

# CMS Draft Analysis Note

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## Search for SM tttt in the same-sign dilepton and multi-lepton final states at $\sqrt{s} = 13$ TeV with the full Run 2 dataset

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### Abstract

This is the AN supporting the Run 2 tttt analysis. It is based on the 2016 tttt AN (AN-17-115) and the 2016 same-sign AN (AN-16-386).

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DRAFT

96 **1** Introduction

The same-sign dilepton final state is often used in searches for new physics due to its ability to stifle all but some rare standard model backgrounds. One of these rare backgrounds is the production of four top quarks,  $pp \rightarrow t\bar{t}t\bar{t}$ , primarily through the two processes depicted in Fig. 1. The lepton multiplicity composition in  $t\bar{t}t\bar{t}$  events is shown in Fig. 2. With four W bosons, the same-sign dilepton and tri-lepton final state, considering electrons and muons (and those from arising from leptonic tau decays), occurs in approximately 11.9% of  $t\bar{t}t\bar{t}$  events.

This Analysis Note describes a  $t\bar{t}t\bar{t}$  analysis using the full Run 2 data, which is based on two analyses of the 2016 data: the SUSY same-sign dilepton analysis and the corresponding  $t\bar{t}t\bar{t}$  analysis. The SUSY same-sign dilepton analysis used  $35.9 \text{ fb}^{-1}$  of 2016 collision data, is documented in [1], and published in [2]. The observed (expected) 95% confidence limit on cross section of four top quark production was found to be  $42 (27^{+13}_{-8}) \text{ fb}$ , and the cross section of the process calculated at NLO is  $9.2^{+2.9}_{-2.4} \text{ fb}$  [3].

The 2016  $t\bar{t}t\bar{t}$  analysis, based on the same dataset as the SUSY search, was developed by modifying the selection of the SUSY analysis to take advantage of the jet and b-jet multiplicity associated with four top quark decays, as well as creating a separate category of events with 3 or more leptons. The resulting 2016  $t\bar{t}t\bar{t}$  analysis is document in [4], and [5]. The observed (expected) 95% confidence limit on cross section of four top quark production was found to be  $41.7 (20.8^{+11.2}_{-6.9})$  fb. The expected significance, based on an NLO cross section for  $t\bar{t}t\bar{t}$  of  $9.2^{+2.9}_{-2.4}$  fb [3], was 1.05 standard deviations, corresponding to a measured signal strength parameter of  $1.0^{+1.2}_{-0.94}$ . The observed significance was 1.56 standard deviations (p-value of 6%), corresponding to a measured signal strength parameter of  $1.83^{+1.42}_{-1.23}$ . The results were also used to constrain the top quark Yukawa coupling of the Higgs boson, based on an LO\*k-factor cross section of  $12.2^{+5.0}_{-4.4}$  fb and its dependence on  $|y_T/y_T^{SM}|$ . The central (upper,lower) value of the theoretical cross section band, combined with our 95% observed limit on  $t\bar{t}t\bar{t}$  production, resulted in a 95% CL limit  $|y_T/y_T^{SM}| < 2.27 (2.03,2.56)$ . Additionally, an extension to the interpretation of heavy (pseudo)scalar bosons in the context of 2HDM in the manner of the SUSY analysis was studied, but it was not included in the  $t\bar{t}t\bar{t}$  publication due to the very limited improvements with respect to the SUSY analysis, as discussed in Section 9 of Ref. [4].

<sup>125</sup> All details of the 2016 analysis can be found in the SUSY AN [1] (Fake and Flip backgrounds)  
<sup>126</sup> and the TOP AN [4] (everything else).

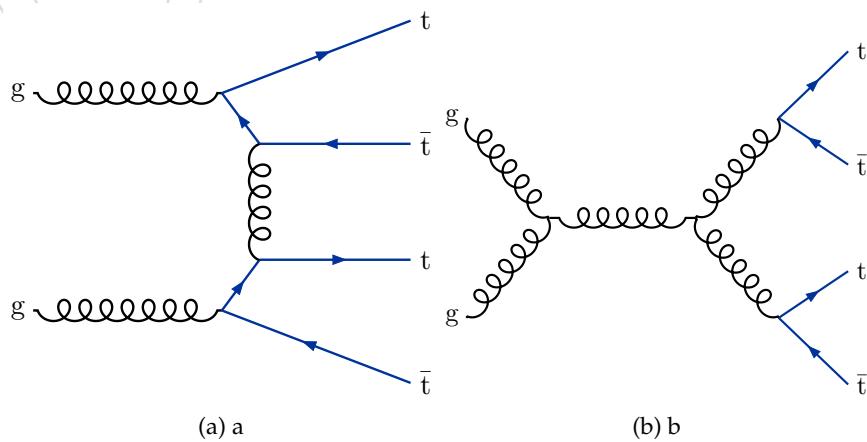


Figure 1: Leading order diagrams for  $t\bar{t}t\bar{t}$  production.

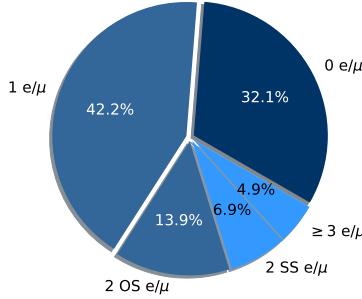


Figure 2: Lepton multiplicity for decay of 4 W bosons.

With respect of the 2016 tt $t\bar{t}$  analysis, the Run2 analysis includes the following changes and improvements:

- Increased luminosity of  $35.922 + 41.53 + 58.83 = 136.3 \text{ fb}^{-1}$
- Update the cut-and-count signal regions from 8 to 18 bins in N(lep)/N(jets)/N(b-jets)
- Note that this was the case for ANv7 and previous versions. Update the correction to ttW+bb and ttZ+bb based on the Data/MC discrepancy in the ttbb/ttjj ratio from  $1.7 \pm 0.6$  (TOP-16-010) to  $1.0 \pm 0.35$ . This is based on the preliminary results of TOP-17-021, which finds a ttbb/ttjj ratio of  $1.0 \pm 0.35$ .
- This is the current prescription for ANv8 and on. Maintain the correction to ttW+bb and ttZ+bb based on the Data/MC discrepancy in the ttbb/ttjj ratio of  $1.7 \pm 0.6$  (TOP-16-010).
- Update t $\bar{t}$ t $\bar{t}$  cross section to the latest and most precise estimate, from 9.2 fb [3] to 11.97 fb [6].
- Reduced normalization uncertainty on t $\bar{t}$ H from 50% to 25% motivated by HIG-17-035.
- Studied two BDTs, one with a long list and one with a short list of variables. See Section 6.2.
- Optimized hyperparameter choice for BDT. See Section 6.3.

Upon looking at the 2017 and 2018 datasets, several more changes were adopted in order to deal with data-taking issues:

- Since the dilepton+HT triggers were not active for the first period of 2017 data taking (RunB, corresponding to 12% of the data), we decided to use the dilepton triggers throughout the 2017 run. The difference resulting from this change is expected to be at the few % level, although we will only know the exact number after all Trigger Scale Factors will be computed.
- Since an increase in non-prompt leptons was observed in our control regions, we tightened the isolation requirements with respect to 2016, as described in Section 4.3.

Issues relating to the missing HEM15/16 sectors starting in Run 319077 of 2018 data-collection were found to impact the fake background only negligibly, and no action was taken to specifically deal with this as the estimation method is data-driven. Additionally, the L1 prefiring issue affecting the efficiency of EGamma objects at high  $|\eta|$  was found to have a small effect (%-level) which can be accounted for by a scale factor applied to simulation. See appendix D for checks,

160 and also see the 2LSS+multi-lepton SUSY analysis presentation at the SUSY Leptonic Meeting  
161 (slides 21-25 for HEM and slides 26-28 for L1 prefiring at <https://indico.cern.ch/event/770079/>).  
162

163 Several changes were considered, but it was decided not to implement them, as motivated  
164 below:

- 165 • Many additional variables were considered for the BDT, including resolved and  
166 boosted top-tagging variables, but they were not found to bring significant ( $> 2\%$ )  
167 gain. See Section 6.2.
- 168 • Separate BDT trainings for different background categories were considered, but  
169 since all backgrounds were found to have a similar BDT shape, it was decided to  
170 continue with a single training based on the sum of all backgrounds (slide 14 from  
Nick at <https://indico.cern.ch/event/709496/>).
- 171 • The LeptonMVA selection (combining ID, isolation and nearby jet information) was  
172 studied, but ultimately not adopted. Due to the different kinematics, it was found  
173 to increase ttV backgrounds by more than the tt $\bar{t}$  signal, resulting in no gain in  
174 significance. See Nick’s presentation at the Top Cross Section meeting (slides 4-8 at  
175 <https://indico.cern.ch/event/709496/>).
- 176 • Jet and b-jet  $p_T$  thresholds were studied in the range of 20-40 GeV. Based on ex-  
177 pected significance, it was decided to stay with 40 GeV threshold for jets, and with  
178 25 GeV threshold for b-jets. See Appendix C.
- 179 • Including events with tau leptons. See Appendix C.
- 180 • Events with a 3rd lepton, with  $10 < p_T < 20$  GeV and not forming a Z mass, were  
181 previously included in the 2-lepton signal regions. We considered removing these  
182 events, or moving them to a separate categories, but found that both treatments  
183 actually worsened (by a few %) the significance, as a non-negligible fraction of tt $\bar{t}$   
184 events falls into this category.

185 An extensive changelog with respect to ANv7 can be found in Appendix E.

186 Changelog with respect to ANv9

- 187 • Incorporation of lepton scale factors for 2017, 2018 (muons are incomplete for 2017,  
188 2018, and electrons are incomplete for 2018)
- 189 • Switched to preliminary Autumn18V3 JECs applied to 2018 data, MC (with no resid-  
190 uals for data)
- 191 • Added Appendix F and Appendix G showing the unblinding of 2016 and 2017  
192 datasets, respectively.
- 193 • For convenience, significance numbers will also be tracked in <https://docs.google.com/spreadsheets/d/140yqrUwEsJtOJ80mDdOun-J8iUrV49aAdUkvfNjpn8k/edit#gid=0>

196 Changelog with respect to ANv10

- 197 • Updated to Autumn18V8 JECs applied to 2018 data, MC (includes residuals for 2018,  
198 split into eras A, B, C, and D)
- 199 • Updated to latest recommended b-tag SFs (V1) for 2018 MC
- 200 • Retrained the BDT with final samples and these updated corrections
- 201 • Added Appendix H which shows the results of unblinding 2018 and the full Run2  
202 dataset, with a retrained BDT

- 203     • Updated documentation to reflect the full unblinding  
 204     • For convenience, significance numbers are tracked in <https://docs.google.com/spreadsheets/d/140yqrUwEsJtOJ8OmDdOun-J8iUrV49aAdUkvfNjpn8k/edit#gid=0>

207 Changelog with respect to ANv12

- 208     • Included final lepton SFs (electrons, muons), finalized 2018 muon scale factors  
 209     • Included 50ipb from new re-reco golden JSON in 2018  
 210     • Updated to latest luminosity recommendation for 2018 ( $59.6 \text{ fb}^{-1}$ ), with a total of  
 211        $137.2 \text{ fb}^{-1}$   
 212     • Switched to final luminosity uncertainty of 2.5% for 2018  
 213     • Included latest 2018 samples for heavy (pseudo)scalar boson interpretation

## 214 2 Samples

### 215 2.1 Collision data

216 The full Run 2 analysis uses an amount of pp collision data corresponding to an integrated  
 217 luminosity of  $35.9 \text{ fb}^{-1}$ ,  $41.5 \text{ fb}^{-1}$ , and  $59.6 \text{ fb}^{-1}$ , for 2016, 2017, and 2018, totaling  $137.2 \text{ fb}^{-1}$ .

218 The following datasets are used:

- 219     • DoubleEG/EGamma: two-electron channel,
- 220     • DoubleMuon: two-muon channel,
- 221     • MuonEG: crossed channel targeting events with one muon and one electron.
- 222     • JetHT: used to recover efficiency for dilepton+ $H_T$  triggers in period 2016H (although  
 223       no additional events are found in this sample).

224 The Reco campaigns used in this analysis are listed in table 1, together with the corresponding  
 225 JSON file. Events present in multiple datasets are properly handled.

Table 1

Year	campaign	JSON file
2016	23Sep2016-v1	Cert_271036-284044_13TeV_23Sep2016ReReco_Collisions16
2017	31March2018-v1	Cert_294927-306462_13TeV_EOY2017ReReco_Collisions17
2018	17Sep2018-v1(2018A,B,C)-PromptReco(2018D)	Cert_314472-325175_13TeV_PromptReco_Collisions18_JSON.txt

226 In order to remove detector noise and unphysical events such beam-halo events, the recom-  
 227 mended JetMET filters [7] have been applied.

### 228 2.2 Monte Carlo Simulation

229 The simulation samples used in this analysis have been produced in the Summer16 campaign  
 230 for comparing to 2016 data, in the RunIIFall17MiniAOD-94X campaign for comparing to 2017  
 231 data, and in the RunIIAutumn18MiniAOD-102X campaign for 2018 data. The samples, almost  
 232 exclusively at NLO, are listed in Table 2.

233 To improve on the MADGRAPH5\_AMC@NLO modeling of the multiplicity of additional jets  
 234 from initial-state radiation (ISR) and final-state radiation (FSR),  $t\bar{t}W$  and  $t\bar{t}Z$  MC events are

235 reweighted based on the number of ISR or FSR jets ( $N_j^{\text{ISR/FSR}}$ ). The reweighting is based on the  
 236 Data/MC ratio in the light-flavor jet multiplicity in dilepton  $t\bar{t}$  events (using MADGRAPH5\_AMC@NLO  
 237 MC), as shown in Figure 3 and described in Ref. [8] for 2016 Data and MC. The method requires  
 238 exactly two b-tagged jets, and assumes that all other jets are ISR or FSR. The reweighting factors  
 239 vary between 0.86 and 0.77 for  $N_j^{\text{ISR/FSR}}$  between 1 and 4 (where dilepton  $t\bar{t}V$  plus 4 ISR/FSR  
 240 jets results in  $N_{\text{jets}}=8$ , corresponding to our highest SR bin). For 2017 MC, new weights were  
 241 derived since the MC uses a new Pythia tune. The corresponding Data/MC ratios are shown  
 242 in Figure 4, where the different plots are based on different number of partons simulated at  
 243 the matrix element. Reweighting factors are applied to  $t\bar{t}W$ ,  $t\bar{t}Z$ , and  $t\bar{t}$  (from inclusive sam-  
 244 ple reweighting). Following the procedure used in the SUSY group, we take one half of the  
 245 deviation from unity as the systematic uncertainty in these reweighting factors.

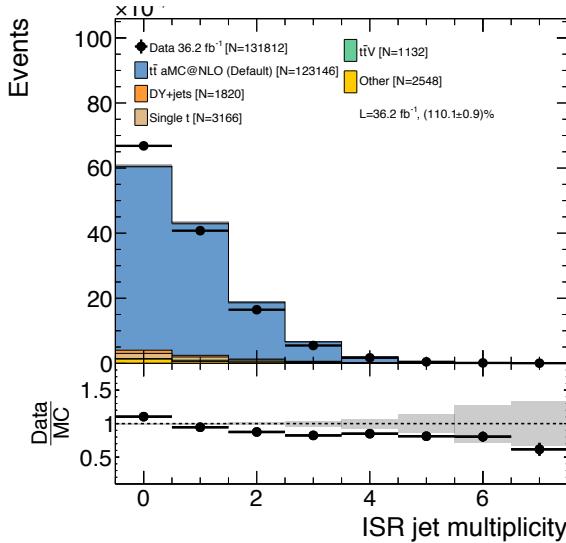


Figure 3: Distribution of the number of light jets (labeled ISR jets in the plot, but referring to ISR and FSR jets) in a dilepton  $t\bar{t}$  sample in 2016 data, compared to 2016 MC.

246 To improve the modeling of the flavor of additional jets, the  $t\bar{t}W$ ,  $t\bar{t}Z$  and  $t\bar{t}H$  simulation is also  
 247 corrected to account for the measured ratio of  $t\bar{t}bb/t\bar{t}jj$  cross sections reported in TOP-16-010.  
 248 More details on these corrections and their uncertainties are provided in Section 9.

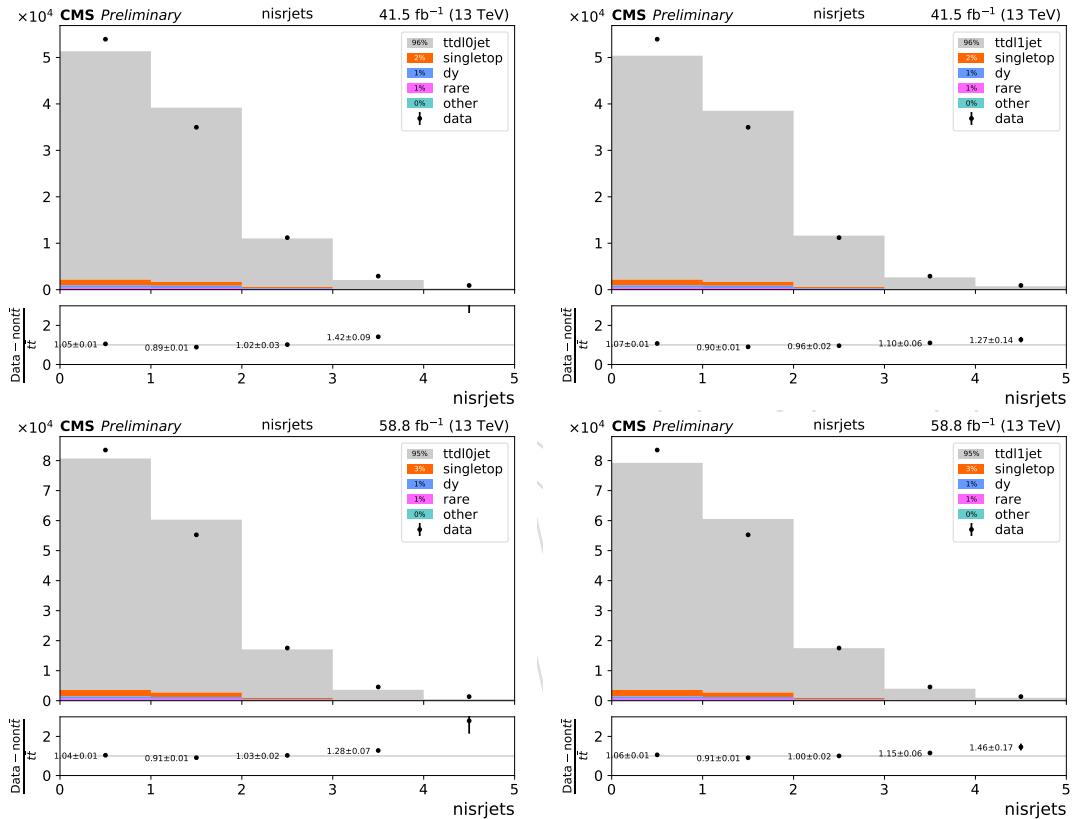


Figure 4: Distribution of the number of light jets in a dilepton  $t\bar{t}$  sample with 0 additional partons (left) or 1 additional parton (right) for 2017 data (top) and 2018 data (bottom) compared to their respective MC samples. In the case of 0 additional partons, the reweighting factor in the last bin is taken instead from the previous bin.

Table 2: Signal and background samples from Summer16 and Fall17 campaigns. For each sample, the name has to be completed with the Summer16 string

RunIISummer16MiniAODv2-PUMoriond17\_80X\_mcRun2\_asymptotic\_2016\_TrancheIV\_v6\_v\* for 2016, and with the Fall17 string RunIIFall17MiniAODv2-PU2017\_12Apr2018\_94X\_mc2017\_realistic for 2017. For 2018, the string is RunIIAutumn18MiniAOD-102X\_upgrade2018. Starting in 2017, the new Pythia tune “TuneCP5” was used instead of the CUETP8M2T4 and CUETP8M1 tunes. The  $t\bar{t}t\bar{t}$  signal samples are highlighted in red. The ISR/FSR variation samples for  $t\bar{t}t\bar{t}$  are not necessary starting in 2017, since the Parton Shower variation weights are already included in the nominal sample. **Remaining missing samples QCD samples are only used as a crosscheck, and the analysis does not depend on them.**

sample name	$\sigma$ (pb)	2016	2017	2018
/DYJetsToLL_M-50.TuneCUETP8M1.13TeV-madgraphMLM-pythia8	6020.85	X	X	X
/DYJetsToLL_M-10to50.TuneCUETP8M1.13TeV-madgraphMLM-pythia8	18610	X	X	X
/WJetsToLNu.TuneCUETP8M1.13TeV-madgraphMLM-pythia8	61334.9	X	X	X
/TT.TuneCUETP8M2T4.13TeV-powheg-pythia8	831.762	X	X	X
/TTJets.TuneCP5.13TeV-amcatnloFXFX-pythia8	831.762	X	X	X
/TTWJetsToLNu.TuneCUETP8M1.13TeV-amcatnloFXFX-madspin-pythia8	0.2043	X	X	X
/TTZToLLNuNu_M-10.TuneCUETP8M1.13TeV-amcatnlo-pythia8	0.2529	X	X	X
/TTZToLL_M-10to10.TuneCUETP8M1.13TeV-madgraphMLM-pythia8	0.0493	X	X	X
/ttHToNonbb_M125.TuneCUETP8M2_ttHanche3.13TeV-powheg-pythia8	0.2710	X	X	X
/fZq_ll_4f_13TeV-amcatnlo-pythia8.TuneCUETP8M1 (ext1)	0.0758	X	X	X
/TTTT_TuneCP5_PSwights.13TeV-amcatnlo-pythia8	0.01197	-	X	X
/TTTT_TuneCUETP8M2T4_13TeV-amcatnlo-isrup-pythia8	0.01197	X	-	X
/TTTT_TuneCUETP8M2T4_13TeV-amcatnlo-fsrup-pythia8	0.01197	X	-	X
/TTTT_TuneCUETP8M2T4_13TeV-amcatnlo-isrdown-pythia8	0.01197	X	-	X
/TTTT_TuneCUETP8M2T4_13TeV-amcatnlo-fsrdown-pythia8	0.01197	X	-	X
/TGJets.TuneCUETP8M1_13TeV_amcatnlo_madspin_pythia8	2.967	X	X	X
/TTGamma_SingleLeptFromT.TuneCUETP8M2T4_13TeV-amcatnlo-pythia8	0.77	X	X	X
/TTGamma_SingleLeptFromTbar.TuneCUETP8M2T4_13TeV-amcatnlo-pythia8	0.769	X	X	X
/TTGamma_Dilept.TuneCUETP8M2T4_13TeV-amcatnlo-pythia8	0.632	X	X	X
/WGToLNuG.TuneCUETP8M1_13TeV-amcatnloFXFX-pythia8	405.271	X	-	X
/WGToLNuG.TuneCP5_13TeV-madgraphMLM-pythia8	405.271	-	X	-
/ZGTo2LG.TuneCUETP8M1_13TeV-amcatnloFXFX-pythia8	123.9	X	X	X
/WZTo3LNu.TuneCUETP8M1_13TeV-powheg-pythia8	4.4297	X	X	X
/ZZTo4L_13TeV_powheg-pythia8	1.256	X	X	X
/ZZZ.TuneCUETP8M1_13TeV-amcatnlo-pythia8	0.01398	X	X	X
/WZZ.TuneCUETP8M1_13TeV-amcatnlo-pythia8	0.05565	X	X	X
/WWZ.TuneCUETP8M1_13TeV-amcatnlo-pythia8	0.1651	X	X	X
/WZG.TuneCUETP8M1_13TeV-amcatnlo-pythia8	0.04123	X	X	X
/WWG.TuneCUETP8M1_13TeV-amcatnlo-pythia8	0.2147	X	X	X
/WWW_4F.TuneCUETP8M1_13TeV-amcatnlo-pythia8	0.2086	X	X	X
/WWTo2LNu.DoubleScattering_13TeV-pythia8	0.16975	X	X	X
/ST_tWIL_5f_LO_13TeV-MadGraph-pythia8	0.01123	X	X	X
/WpWpJJ_EWK-QCD_TuneCUETP8M1_13TeV-madgraph-pythia8	0.03711	X	X	X
/GluGluHToZZTo4L_M125_13TeV_powheg2_JHUGenV6_pythia8	0.01181	X	X	X
/VHToNonbb_M125_13TeV_amcatnloFXFX_madspin_pythia8	0.9561	X	X	X
/TTHH.TuneCUETP8M2T4_13TeV-madgraph-pythia8	0.000757	X	X	X
/TTZH.TuneCUETP8M2T4_13TeV-madgraph-pythia8	0.001535	X	X	X
/TTZZ.TuneCUETP8M2T4_13TeV-madgraph-pythia8	0.001982	X	X	X
/TTWZ.TuneCUETP8M2T4_13TeV-madgraph-pythia8	0.003884	X	X	X
/TTW.TuneCUETP8M2T4_13TeV-madgraph-pythia8	0.000788	X	X	X
/TTJ.TuneCUETP8M2T4_13TeV-madgraph-pythia8	0.000474	X	X	X
/TTWH.TuneCUETP8M2T4_13TeV-madgraph-pythia8	0.001582	X	X	X
/TTWW.TuneCUETP8M2T4_13TeV-madgraph-pythia8	0.01150	X	X	X
/QCD_Pt-20toInf_MuEnrichedPt15.TuneCUETP8M1_13TeV_pythia8	720648000 $\times 0.00042$	X	X	TODO
/QCD_Pt-15to20_MuEnrichedPt5.TuneCUETP8M1_13TeV_pythia8	1273190000 $\times 0.003$	X	X	TODO
/QCD_Pt-20to30_MuEnrichedPt5.TuneCUETP8M1_13TeV_pythia8	558528000 $\times 0.0053$	X	X	TODO
/QCD_Pt-30to50_MuEnrichedPt5.TuneCUETP8M1_13TeV_pythia8	139803000 $\times 0.01182$	X	X	TODO
/QCD_Pt-50to80_MuEnrichedPt5.TuneCUETP8M1_13TeV_pythia8	19222500 $\times 0.02276$	X	X	TODO
/QCD_Pt-80to120_MuEnrichedPt5.TuneCUETP8M1_13TeV_pythia8	2758420 $\times 0.03844$	X	X	TODO
/QCD_Pt-120to170_MuEnrichedPt5.TuneCUETP8M1_13TeV_pythia8	469797 $\times 0.05362$	X	X	TODO
/QCD_Pt-170to300_MuEnrichedPt5.TuneCUETP8M1_13TeV_pythia8	117989 $\times 0.07335$	X	X	TODO
/QCD_Pt-470to600_MuEnrichedPt5.TuneCUETP8M1_13TeV_pythia8	645.528 $\times 0.12242$	X	X	TODO
/QCD_Pt-600to800_MuEnrichedPt5.TuneCUETP8M1_13TeV_pythia8	187.109 $\times 0.13412$	X	X	TODO
/QCD_Pt-1000toInf_MuEnrichedPt5.TuneCUETP8M1_13TeV_pythia8	10.4305 $\times 0.15544$	X	X	TODO
/QCD_Pt-20to30_EMEnriched.TuneCUETP8M1_13TeV_pythia8	557600000 $\times 0.0096$	X	X	TODO
/QCD_Pt-30to50_EMEnriched.TuneCUETP8M1_13TeV_pythia8	136000000 $\times 0.073$	X	X	TODO
/QCD_Pt-50to80_EMEnriched.TuneCUETP8M1_13TeV_pythia8	19800000 $\times 0.146$	X	X	TODO
/QCD_Pt-80to120_EMEnriched.TuneCUETP8M1_13TeV_pythia8	2800000 $\times 0.125$	X	X	TODO
/QCD_Pt-120to170_EMEnriched.TuneCUETP8M1_13TeV_pythia8	477000 $\times 0.132$	X	X	TODO
/QCD_Pt-170to300_EMEnriched.TuneCUETP8M1_13TeV_pythia8	114000 $\times 0.165$	X	X	TODO
/QCD_Pt-300toInf_EMEnriched.TuneCUETP8M1_13TeV_pythia8	9000 $\times 0.15$	X	X	TODO
/QCD_Pt-20to30_bcToE.TuneCUETP8M1_13TeV_pythia8	557627000 $\times 0.00059$	X	X	TODO
/QCD_Pt-30to80_bcToE.TuneCUETP8M1_13TeV_pythia8	159068000 $\times 0.00255$	X	X	TODO
/QCD_Pt-80to170_bcToE.TuneCUETP8M1_13TeV_pythia8	3221000 $\times 0.01183$	X	X	TODO
/QCD_Pt-170to250_bcToE.TuneCUETP8M1_13TeV_pythia8	105771 $\times 0.02492$	X	X	TODO
/QCD_Pt-250toInf_bcToE.TuneCUETP8M1_13TeV_pythia8	21094.1 $\times 0.03375$	X	X	TODO

### 249 3 Triggers

250 The analysis uses two slightly different trigger strategies for 2016 and 2017/2018 data. For 2016  
 251 data, the dilepton +  $H_T$  triggers are used, since they have slightly looser lepton requirements  
 252 (they do not apply isolation), and are therefore slightly more efficient than the pure dilepton  
 253 triggers. For 2017, these triggers were not available in RunB, which represent 12% of the 2017  
 254 dataset, and therefore the pure dilepton triggers are used for the entirety of 2017 and 2018. We  
 255 considered also using the pure dilepton triggers for 2016, to further simplify this treatment, but  
 256 we prefer to avoid unnecessary changes to the published result, wherever possible. The full set  
 257 of triggers used for selecting events in the signal regions are listed in Table 3.

Table 3: Summary of the signal triggers

2016		
$H_{T,off}$	Channel	Trigger Name
$> 300 \text{ GeV}$	$\mu\mu$	HLT_DoubleMu8_Mass8_PFHT300
	$ee$	HLT_DoubleEle8_CaloIdM_TrackIdM_Mass8_PFHT300
	$e\mu$	HLT_Mu8_Ele8_CaloIdM_TrackIdM_Mass8_PFHT300
2017		
$H_{T,off}$	Channel	Trigger Name
all	$\mu\mu$	HLT_Mu17_TrkIsoVVL_Mu8_TrkIsoVVL_DZ_Mass8
	$ee$	HLT_Ele23_Ele12_CaloIdL_TrackIdL_IsoVL_
	$e\mu$	HLT_Mu23_TrkIsoVVL_Ele12_CaloIdL_TrackIdL_IsoVL_DZ HLT_Mu8_TrkIsoVVL_Ele23_CaloIdL_TrackIdL_IsoVL_DZ
2018		
$H_{T,off}$	Channel	Trigger Name
all	$\mu\mu$	HLT_Mu17_TrkIsoVVL_Mu8_TrkIsoVVL_DZ_Mass3p8
	$ee$	HLT_Ele23_Ele12_CaloIdL_TrackIdL_IsoVL_
	$e\mu$	HLT_Mu23_TrkIsoVVL_Ele12_CaloIdL_TrackIdL_IsoVL_DZ HLT_Mu8_TrkIsoVVL_Ele23_CaloIdL_TrackIdL_IsoVL_DZ

258 An additional set of triggers is developed for a control sample selection for fake rate measure-  
 259 ment. They require either a single lepton, or a single lepton and a jet in order to increase a purity  
 260 of the collected data sample. The full list can be found in Table 4. Two sets of the auxiliary trig-  
 261 gers are introduced in order to collect data with different online lepton IDs corresponding to  
 262 the two sets of signal dilepton triggers.

#### 263 3.1 2016 trigger efficiency

264 Computation of trigger efficiencies and scale factors for 2016 data is described in Section 4  
 265 of [1]. Summarizing these results, the dilepton+ $H_T$  trigger efficiency is found to be the product  
 266 of lepton leg efficiencies (93-98% for electrons, 85-95% for muons), and the HT efficiency (98-  
 267 99%). As a result the trigger efficiency for events passing the baseline selection in the 2016 data  
 268 is greater than 95% for ee and  $e\mu$  events, and about 92% for  $\mu\mu$  events.

#### 269 3.2 2017 and 2018 trigger efficiency

270 This section is identical to same-sign SUS analysis (AN-18-280), which provides trigger effi-  
 271 ciency maps. These are currently being applied for all 3 years.

Table 4: Summary of the control triggers ordered by lepton flavor and  $p_T$ .

2016		
$H_{T,off}$	Channel	Trigger Name
$> 300 \text{ GeV}$	$\mu$	HLT_Mu8 HLT_Mu17
	$e$	HLT_Ele8_CaloIdM_TrackIdM_PFJet30 HLT_Ele17_CaloIdM_TrackIdM_PFJet30
2017		
$H_{T,off}$	Channel	Trigger Name
all	$\mu$	HLT_Mu8_TrkIsoVVL HLT_Mu17_TrkIsoVVL
	$e$	HLT_Ele8_CaloIdL_TrackIdL_IsoVL_PFJet30 HLT_Ele23_CaloIdL_TrackIdL_IsoVL_PFJet30
2018		
$H_{T,off}$	Channel	Trigger Name
all	$\mu$	HLT_Mu8_TrkIsoVVL HLT_Mu17_TrkIsoVVL
	$e$	HLT_Ele8_CaloIdL_TrackIdL_IsoVL_PFJet30 HLT_Ele23_CaloIdL_TrackIdL_IsoVL_PFJet30

## 272 4 Object Definition and Selection

273 The object selection is based on the one used in the SUSY 2016 analysis documented in [1]. In  
 274 particular, the same lepton ID is used for same-sign and trilepton events.

275 The somewhat complicated definitions of isolation and lepton IDs were arrived through dedi-  
 276 cated studies within the "SUSY fake lepton working group" before data taking started in 2015.  
 277 The goal was to have a fairly efficient lepton selection, with low fake rate, and with a reliable  
 278 method to estimate nonprompt lepton backgrounds with minimal sample dependence starting  
 279 from "fakeable objects" (FO). The studies are documented in Ref. [9].

### 280 4.1 Electron identification

281 The electrons are reconstructed by associating ECAL-clusters and GSF tracks. Only electrons  
 282 within the tracker and ECAL acceptance are considered:  $|\eta| < 2.5$ .

283 The electron identification is performed using a multivariate (MVA) discriminant built with  
 284 shower-shape variables ( $\sigma_{\eta\eta}$ ,  $\sigma_{\phi\phi}$ , the cluster circularity, widths along  $\eta$  and  $\phi$ ,  $R_9$ , H/E,  
 285  $E_{inES}/E_{raw}$ ), track-cluster matching variables ( $E_{tot}/p_{in}$ ,  $E_{Ele}/p_{out}$ ,  $\Delta\eta_{in}$ ,  $\Delta\eta_{out}$ ,  $\Delta\phi_{in}$ ,  $1/E - 1/p$   
 286 ) and track quality variables ( $\chi^2$  of the KF and GSF tracks, the number of hits used by the  
 287 KF/GSF filters, fbrem). A complete description of the discriminant and training used can be  
 288 found in [10–12]. Three identification working points, summarized in the Table 5, are used in  
 289 this analysis, following the prescriptions of the SUSY lepton scale factors group [13]. Signal  
 290 events are selected using the tight working point while the loose working points are used in  
 291 the estimate of background arising from mistakenly identified or non prompt leptons. Separate  
 292 loose working points are derived for regions where isolated and non isolated dilepton triggers  
 293 are used.

294 Different cuts are used in 2016, 2017, and 2018 based on the Spring16\_GeneralPurpose\_V1, Fall17NoIso,  
 295 and Fall17V2NoIso trainings, respectively. Note that the cuts for 2018 are with respect to  
 296 the "raw" BDT output (i.e., the values are not forced to lie between -1 and 1 by squashing  
 297 the output with a sigmoid function). Raw values can be obtained from squashed output via  
 298 raw =  $\frac{1}{2} \ln \left( \frac{1+\text{squashed}}{1-\text{squashed}} \right)$ . The working points depend of the momentum and pseudorapidity of  
 299 the electrons: for 2016 (resp. 2017/2018) in the region between 15 (resp. 10) and 25 GeV, the  
 300 working point is given by a linear interpolation between the higher and lower value, following  
 301  $\min[A, \max[B, A + C * (p_T - 15(10))]]$ , where  $A$  is the cut at 15(10) GeV,  $B$  is the cut at 25 GeV,  
 302 and  $C = (B - A) / 10(15)$ . Electrons with  $p_T < 10$  GeV are only not used to select signal events,  
 303 so the Tight WP is not assigned for them.

304 As the electron charge can be determined with three different techniques [14], and as the mis-  
 305 measurement of the lepton charge can lead to accept more background after the selection, the  
 306 requirement to get a consistent charge measurement with the three methods is considered in  
 307 some of the electron definitions used in the analysis. To reject electrons originating from photon  
 308 conversion, two variables are considered: the number of missing pixel hits and a conversion  
 309 veto based on the vertex fit probability. Finally, impact parameter variables are also consid-  
 310 ered: impact parameter in the transverse plane  $d_0$ , impact parameter along the z axis  $d_z$ , and  
 311 the impact parameter significance in the detector space SIP<sub>3D</sub>.

### 312 4.2 Muon identification

313 Two working points are considered for the muon identification. The loose working point fol-  
 314 lows the "muon POG Loose ID" described in [15], while the tight working point is given by  
 315 the list of requirements known as the "muon Medium Id", defined in [16]. Only muons within

Table 5: Lower cut on the electron MVA discriminant for the various electron ID used in this analysis.

2016				
pseudorapidity region	momentum [GeV]	loose WP (non iso trigger)	loose WP (iso trigger)	tight WP
$0 <  \eta  < 0.8$	$5 < p_T < 10$	-0.30	-0.46	N/A
	$10 < p_T < 15$	-0.86	-0.48	0.77
	$p_T > 25$	-0.96	-0.85	0.52
$0.8 <  \eta  < 1.479$	$5 < p_T < 10$	-0.36	-0.03	N/A
	$10 < p_T < 15$	-0.85	-0.67	0.56
	$p_T > 25$	-0.96	-0.91	0.11
$1.479 <  \eta  < 2.5$	$5 < p_T < 10$	-0.63	0.06	N/A
	$10 < p_T < 15$	-0.81	-0.49	0.48
	$p_T > 25$	-0.95	-0.83	-0.01

2017				
pseudorapidity region	momentum [GeV]	loose WP (non iso trigger)	loose WP (iso trigger)	tight WP
$0 <  \eta  < 0.8$	$5 < p_T < 10$	-0.135	0.488	N/A
	$10 < p_T < 25$	$-0.930 + \frac{0.043}{15} \times (p_T - 10)$	$-0.788 + \frac{0.148}{15} \times (p_T - 10)$	$0.2 + 0.032 \times (p_T - 10)$
	$p_T > 25$	-0.887	-0.64	0.68
$0.8 <  \eta  < 1.479$	$5 < p_T < 10$	-0.417	-0.045	N/A
	$10 < p_T < 25$	$-0.930 + \frac{0.04}{15} \times (p_T - 10)$	$-0.850 + \frac{0.075}{15} \times (p_T - 10)$	$0.1 + 0.025 \times (p_T - 10)$
	$p_T > 25$	-0.890	-0.775	0.475
$1.479 <  \eta  < 2.5$	$5 < p_T < 10$	-0.470	0.176	N/A
	$10 < p_T < 25$	$-0.942 + \frac{0.032}{15} \times (p_T - 10)$	$-0.810 + \frac{0.077}{15} \times (p_T - 10)$	$-0.1 + 0.028 \times (p_T - 10)$
	$p_T > 25$	-0.910	-0.733	0.320

2018				
pseudorapidity region	momentum [GeV]	loose WP (non iso trigger)	loose WP (iso trigger)	tight WP
$0 <  \eta  < 0.8$	$5 < p_T < 10$	0.053	1.320	N/A
	$10 < p_T < 25$	$-0.106 + 0.062 \times (p_T - 25)$	$1.204 + 0.066 \times (p_T - 25)$	$4.277 + 0.112 \times (p_T - 25)$
	$p_T > 25$	-0.106	1.204	4.277
$0.8 <  \eta  < 1.479$	$5 < p_T < 10$	-0.434	0.192	N/A
	$10 < p_T < 25$	$-0.769 + 0.038 \times (p_T - 25)$	$0.084 + 0.033 \times (p_T - 25)$	$3.152 + 0.060 \times (p_T - 25)$
	$p_T > 25$	-0.769	0.084	3.152
$1.479 <  \eta  < 2.5$	$5 < p_T < 10$	-0.956	0.362	N/A
	$10 < p_T < 25$	$-1.461 + 0.042 \times (p_T - 25)$	$-0.123 + 0.053 \times (p_T - 25)$	$2.359 + 0.087 \times (p_T - 25)$
	$p_T > 25$	-1.461	-0.123	2.359

the muon system acceptance  $|\eta| < 2.4$  are considered. Impact parameter selection is also applied on muons, and use the variables already defined in the previous section. Contrary to the electrons, only one charge can be reconstructed for a muon track. The quality of the charge reconstruction is then given for the muons by a quality criteria on the track reconstruction :  $\delta p_T(\mu)/p_T(\mu) < 0.2$ .

### 4.3 Lepton isolation

The lepton isolation is constructed using three different variables:

- the mini isolation  $I_{\text{mini}}$ . Requiring  $I_{\text{mini}}$  below a given threshold ensures that the lepton is locally isolated, even in boosted topologies. The impact of pileup is mitigated using the so-called effective area correction:

$$I_{\text{mini}} = \frac{\sum_R p_T(h^\pm) - \max(0, \sum_R p_T(h^0) + p_T(\gamma) - \rho A (\frac{R}{0.3})^2)}{p_T(\ell)}. \quad (1)$$

where  $\rho$  is the pileup energy density, where  $\sum_R p_T(h^\pm)$ ,  $\sum_R p_T(h^0)$  and  $\sum_R p_T(\gamma)$  refers to the sum of the transverse momentum of the charged hadrons, neutral hadrons and photons, respectively, within a cone  $R$ , dependent of the lepton  $p_T$ :

$$R = \frac{10}{\min(\max(p_T(\ell), 50), 200)} \quad (2)$$

The effective areas  $\mathcal{A}$  used are listed in Table 6.

- The second variable is the ratio of the lepton  $p_T$  and of the jet matched to the lepton  $p_T^{\text{ratio}} = \frac{p_T(\ell)}{p_T(\text{jet})}$ . This jet is matched geometrically to the lepton, and in most of the case is the jet containing the lepton. If no jet is clustering the lepton, then the closest one is chosen. The use of  $p_T^{\text{ratio}}$  is a simple way to identify leptons in quite boosted topologies, without any jet reclustering. In order to avoid an over-correction on prompt leptons, the application of the jet energy correction is only applied on the hadronic part of the jet, using the following formula  $\text{jet} = \ell + (\text{jet-PU}-\ell) * JEC - PU$ , where  $\ell$  is the lepton,  $PU$  the pileup energy clustered into the jet, and  $JEC$  the jet energy scale correction to be applied. This formula is applied at the Lorentz vector level. This approach is commonly used in the B2G group and is blessed by the JetMET POG.
- The last variable used is the  $p_T^{\text{rel}}$  variable:

$$p_T^{\text{rel}} = \frac{|(\vec{p}(\text{jet}) - \vec{p}(\ell)) \times \vec{p}(\ell)|}{|\vec{p}(\text{jet}) - \vec{p}(\ell)|}. \quad (3)$$

This variable allows to recover leptons from accidental with jets. Similarly to  $p_T^{\text{ratio}}$ , the jet energy scale correction are only applied on the hadronic part of the considered jet.

Using those three variables, a lepton is considered isolated if the following condition is respected:

$$I_{\text{mini}} < I_1 \wedge (p_T^{\text{ratio}} > I_2 \vee p_T^{\text{rel}} > I_3) \quad (4)$$

The  $I_i, i = 1, 2, 3$  values depends of the flavor of the lepton: as the probability to misidentify a jet is higher for the electrons, tighter isolation values are used. The loose lepton isolation is significantly relaxed, as well as an extra definition (fakeable) used for the fake lepton background estimation. The different values are summarized in the Table 7. The logic beyond that isolation is a relaxing of the local isolation, compensated by the requirement that the lepton carries the major part of the energy of the corresponding jet, or if not, if the lepton is considered as accidentally overlapping with a jet.

More details and figures about the lepton isolation can be found in the previous version of the analysis described in [17].

#### 4.4 Trigger emulation selection

Due to the differences between online and offline reconstruction and selection, some leptons can be selected offline but not online. Even though we simulate this effect by applying trigger selection in the simulation, we still choose to reduce online/offline differences by applying a set of offline requirements to ensure that electrons will pass the trigger. Effectively this increases the efficiency of our trigger selection, reducing the role of trigger Data/MC scale factors. The requirements are summarized in Table 8.

Table 6: Effective areas for muons and electrons.

2016			
Muons		Electrons	
$ \eta $ range	$\mathcal{A}(\mu)$ neutral	$ \eta $ range	$\mathcal{A}(e)$ neutral
0.0 – 0.8	0.0735	0.0 – 1.0	0.1752
0.8 – 1.3	0.0619	1.0 – 1.479	0.1862
1.3 – 2.0	0.0465	1.479 – 2.0	0.1411
2.0 – 2.2	0.0433	2.0 – 2.2	0.1534
2.2 – 2.5	0.0577	2.2 – 2.3 2.3 – 2.4 2.4 – 2.5	0.1903 0.2243 0.2687

2017/2018			
Muons		Electrons	
$ \eta $ range	$\mathcal{A}(\mu)$ neutral	$ \eta $ range	$\mathcal{A}(e)$ neutral
0.0 – 0.8	0.0566	0.0 – 1.0	0.1440
0.8 – 1.3	0.0562	1.0 – 1.479	0.1562
1.3 – 2.0	0.0363	1.479 – 2.0	0.1032
2.0 – 2.2	0.0119	2.0 – 2.2	0.0859
2.2 – 2.5	0.0064	2.2 – 2.3 2.3 – 2.4 2.4 – 2.5	0.1116 0.1321 0.1654

Table 7: Isolation working points

2016			
isolation value	loose WP ( $e/\mu$ )	$\mu$ (Medium) WP	e (Tight) WP
$I_1$	0.4	0.16	0.12
$I_2$	0	0.76	0.80
$I_3$	0	7.2	7.2

2017/2018			
isolation value	loose WP ( $e/\mu$ )	$\mu$ (Medium) WP	e (Tight) WP
$I_1$	0.4	0.11	0.07
$I_2$	0	0.74	0.78
$I_3$	0	6.8	8.0

## 4.5 Lepton definitions

The loose, tight and fakeable lepton definition, based the different identification and isolation definitions described in the previous sections, are summarized in the Table 9.

## 4.6 Jets

Jets are reconstructed from particle flow candidates, clustered with the anti- $k_T$  algorithm and with a cone size of  $\Delta R < 0.4$ . Only jets above a transverse momentum  $p_T > 40$  GeV and within the tracker acceptance  $|\eta| < 2.4$  are considered. To reject noise and mis-measured jets, the selected jets have to fulfill several identification criteria.

For the 2016 data analysis the following criteria, corresponding to the loose JetID selection, are applied:

- neutral hadronic energy fraction  $< 0.99$

Table 8: Electron selection criteria used for the emulation of the high level trigger selection .

variable	barrel, $ \eta  < 1.4442$	endcaps, $1.566 <  \eta  < 2.4$
Identification criteria		
$ \Delta\eta In  <$	0.01	0.01
$ \Delta\phi In  <$	0.04	0.08
$\sigma I\eta I\eta <$	0.011	0.031
$H/E <$	0.08	0.08
$ 1/E - 1/p  <$	0.01	0.01
Isolation criteria (Not used with Dilepton+ $H_T$ triggers, since those triggers do not apply isolation.)		
relative ecal PFCluster isolation ( $dR=0.3$ )		0.45
relative hcal PFCluster isolation ( $dR=0.3$ )		0.25
relative track $p_T$ isolation		0.2

variable	muons			electrons		
	loose	fakable	tight	loose	fakable	tight
identification	loose ID	medium ID	medium ID	loose WP	loose WP	tight WP
isolation	loose WP	loose WP	$\mu$ WP	loose WP	loose WP	e WP
HLT emulation	-	-	-	-	-/yes	-/yes
$d_0$ (cm)	0.05	0.05	0.05	0.05	0.05	0.05
$d_z$ (cm)	0.1	0.1	0.1	0.1	0.1	0.1
SIP <sub>3D</sub>	-	< 4	< 4	-	< 4	< 4
missing inner hits	-	-	-	$\leq 1$	= 0	= 0
conversion veto	-	-	-	yes	yes	yes
tight charge	-	yes	yes	-	yes	yes

Table 9: Summary of the lepton selection.

- neutral electromagnetic energy fraction  $< 0.99$
- number of constituents  $> 1$
- charged hadron fraction  $> 0$
- charged multiplicity  $> 0$
- charged EM fraction  $< 0.99$

For the 2017 and 2018 data analysis we followed the JetMet recommendation to switch to the Tight JetID selection defined as the following set of conditions:

- neutral hadronic energy fraction  $< 0.9$
- neutral electromagnetic energy fraction  $< 0.9$
- number of constituents  $> 1$
- charged hadron fraction  $> 0$
- charged multiplicity  $> 0$

The most recently recommended set of jet energy corrections (JEC) is applied:

- 2016: Summer16\_23Sep2016 (corresponding Global Tag: 80X\_dataRun2\_2016SeptRepro\_v7)
- 2017: Fall17\_17Nov2017B\_V32
- 2018: Autumn18\_RunX\_V8

To avoid double counting due to jets matched geometrically with a lepton, the jet the closest matched to a fakeable or tight lepton within  $\Delta R < 0.4$  is not considered as a jet in the event.

From those selected jets, the key variable  $H_T$  is defined by

$$H_T = \sum_{jets} p_T.$$

### 381 4.7 Jets originating from b quarks

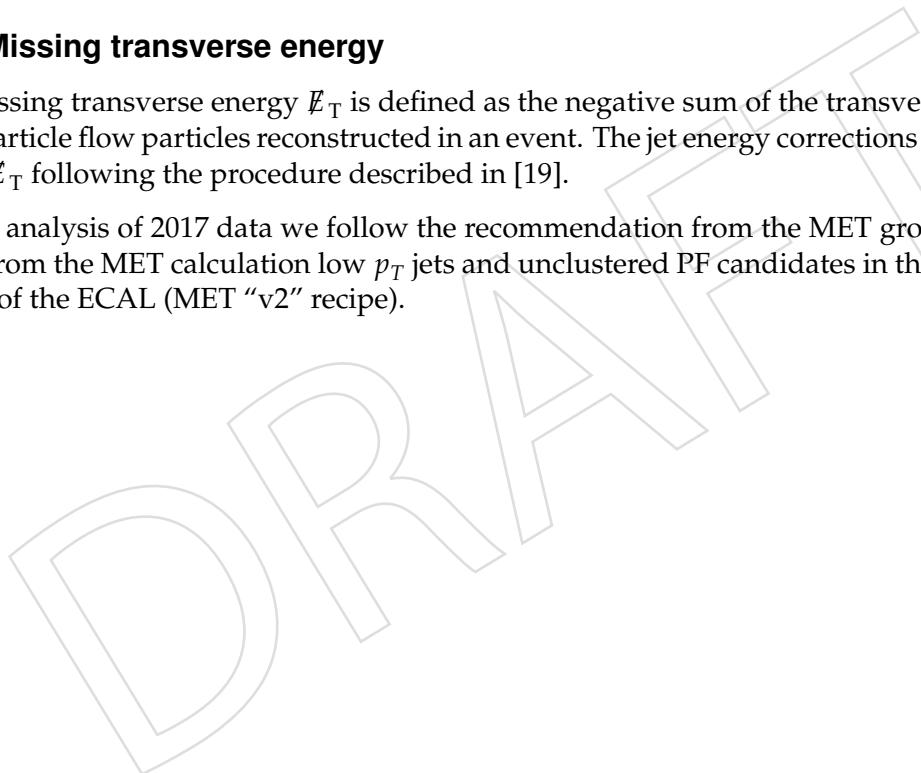
382 Jets defined in the previous section - with the only difference of a lower  $p_T$  threshold down  
383 to 25 GeV - can be promoted as jets originating from b quark. The “deepCSV” discrimina-  
384 tor is used for all years. The discriminator value for each jet is obtained by summing the  
385 “ $x:probb$ ” and “ $x:probbb$ ” discriminators. In 2016,  $x = \text{deepFlavourJetTags}$ , in 2017,  $x =$   
386  $\text{pfDeepCSVJetTags}$ , and in 2018,  $x = \text{pfDeepCSVJetTags}$  [18]. The medium working point  
387 requires the discriminant value to be greater than 0.6324 in 2016, 0.4941 in 2017, and 0.4184 in  
388 2018.

389 When the jet multiplicity and  $H_T$  are computed, b-jets with  $p_T > 40$  GeV are counted along  
390 with standard jets.

### 391 4.8 Missing transverse energy

392 The missing transverse energy  $\cancel{E}_T$  is defined as the negative sum of the transverse momentum  
393 of all particle flow particles reconstructed in an event. The jet energy corrections are propagated  
394 to the  $\cancel{E}_T$  following the procedure described in [19].

395 For the analysis of 2017 data we follow the recommendation from the MET group [20] and ex-  
396 clude from the MET calculation low  $p_T$  jets and unclustered PF candidates in the most forward  
397 region of the ECAL (MET “v2” recipe).



## 398 5 Signal Extraction Strategy: cut-based

### 399 5.1 Baseline selection

400 With respect to Ref. [1], the baseline region was modified for the 2016 analysis to take advantage  
 401 of the four top kinematics. The baseline region is designed to reject the majority of background  
 402 (while preserving signal) before splitting kinematic phase space to form the signal regions.

403 With respect to the 2016 data analysis, the definition of individual signal regions has been  
 404 modified (5.3), while the control regions (5.2) and the overall baseline selection (below) have  
 405 remained the same.

406 Selected events must contain at least two tight (as defined in Table 9) same-sign leptons. To  
 407 reject low-mass resonances, events with a pair of same-sign electrons with an invariant mass  
 408  $m_{\ell\ell} < 12 \text{ GeV}$  are rejected. Events with three leptons of the same charge are also rejected.

409 In total, the baseline selection consists of the following:

- 410 • tight same-sign dileptons,
- 411 •  $N_{\text{jets}} \geq 2$ ,
- 412 •  $N_{\text{b jets}} \geq 2$ ,
- 413 •  $\cancel{E}_T > 50 \text{ GeV}$ ,
- 414 •  $H_T > 300 \text{ GeV}$ ,
- 415 •  $p_{\text{T,lep1}} \geq 25 \text{ GeV}, p_{\text{T,lep2}} \geq 20 \text{ GeV}$ .

416 Events containing a third loose (as defined in Table 9) lepton that makes a DY candidate (opposite charge, same flavor pair with  $|m_{\ell\ell} - m_Z| < 15 \text{ GeV}$  or  $m_{\ell\ell} < 12 \text{ GeV}$ ) with one of the two  
 417 SS dileptons are rejected from the signal regions, and a subset of them is assigned to the  $t\bar{t}Z$   
 418 control region mentioned below. The transverse momentum of the third lepton is required to  
 419 be larger than 5(7) GeV for muons (electrons).

421 If a Z candidate is not found, a tight third lepton present in the event with  $p_{\text{T}} > 20 \text{ GeV}$  contributes  
 422 to the  $N_{\text{leps}}$  count.

### 423 5.2 Control regions selection

424 Two control regions have been introduced to simultaneously constrain two dominant SM backgrounds:  $t\bar{t}W$  and  $t\bar{t}Z$ .

426 The control region for  $t\bar{t}Z$  (“CRZ”) consists of events passing the baseline selection with three  
 427 leptons (where the third lepton has  $p_{\text{T}} > 20 \text{ GeV}$  and passes the tight selection), with a Z boson  
 428 candidate as described above.

429 The control region for  $t\bar{t}W$  (“CRW”) consists of events with a same-sign lepton pair passing the  
 430 tight lepton requirement, fewer than 6 jets, and exactly 2 b-tagged jets.

### 431 5.3 Signal regions selection

432 Signal regions are formed by subdividing the baseline region by number of leptons ( $N_{\text{leps}}$ ),  
 433 number of b-jets ( $N_{\text{b jets}}$ ), and number of jets ( $N_{\text{jets}}$ ).

434 The 2016+2017+2018 analysis signal regions and CRW/CRZ are tabulated in Table 10, and  
 435 shown graphically in Figure 5, while the old 2016 signal regions are shown in Table 11 for  
 436 reference.

Table 10: Definition of the 14 SRs and two CRs, CRW and CRZ, for the full Run2 analysis.

$N_{\text{leps}}$	$N_{\text{b jets}}$	$N_{\text{jets}}$	Region
2	2	$\leq 5$	CRW
		6	SR1
		7	SR2
		$\geq 8$	SR3
	3	5	SR4
		6	SR5
		7	SR6
		$\geq 8$	SR7
	$\geq 4$	$\geq 5$	SR8
	$\geq 3$	5	SR9
		6	SR10
		$\geq 7$	SR11
		4	SR12
		$\geq 5$	SR13
	$\geq 6$	$\geq 6$	SR14
inverted Z-veto			CRZ

Table 11: SR definitions and CRW definition for the 2016 analysis.

$N_{\text{leps}}$	$N_{\text{b jets}}$	$N_{\text{jets}}$	Region
2	2	$\leq 5$	CRW
		6	SR1
		7	SR2
		$\geq 8$	SR3
	3	5, 6	SR4
		$\geq 7$	SR5
		$\geq 4$	$\geq 5$
	$\geq 3$	2	$\geq 5$
		$\geq 3$	$\geq 4$
inverted Z-veto			CRZ

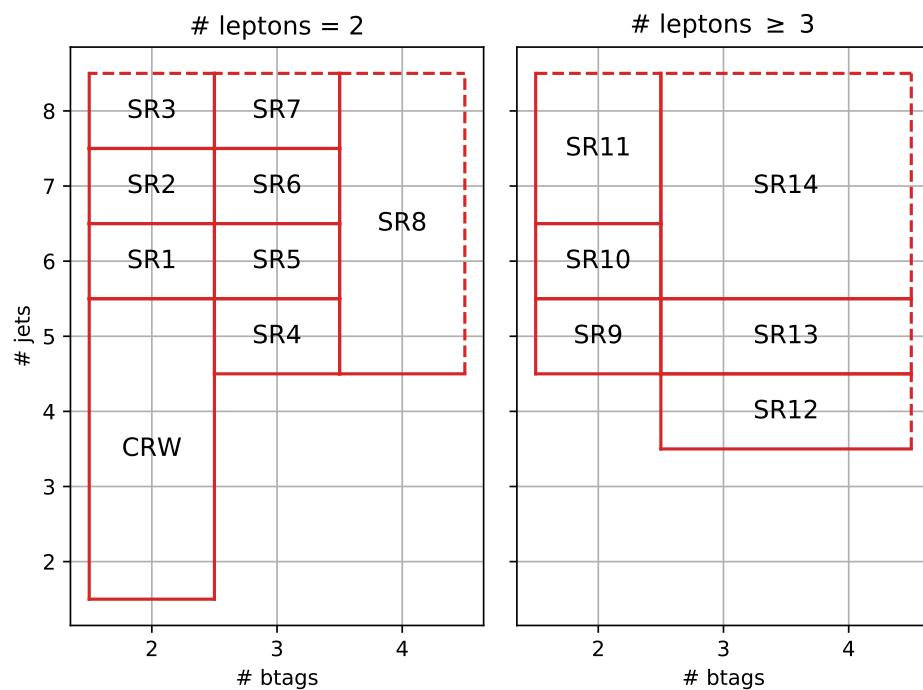


Figure 5: Signal regions and CRZ (15 regions) with CRZ not displayed. Dotted edges indicate no explicit boundary.

## 437 6 Signal Extraction Strategy: BDT

### 438 6.1 Intro

439 A BDT approach was studied for the 2016 analysis, indicating that a 10-20% improvement  
440 in the expected significance could be reached with respect to the cut-based approach, in an  
441 optimistic scenario without additional systematics on the BDT input variables. For the 2016  
442 analysis, the cut-based approach was chosen, resulting in a faster publication and a slightly  
443 lower performance. For the full Run 2 analysis, we have revisited the BDT, with an extended  
444 study of the input variables and hyperparameters, as discussed below. Some of the studies  
445 below are based on  $75 \text{ fb}^{-1}$  of integrated luminosity, as they were performed with a 2016+2017  
446 analysis in mind, but their conclusions also apply to the full Run 2 analysis.

447 A BDT was explored using a well-known python package *xgboost* for a gradient boosted decision  
448 tree. Both ROOT-based TMVA, and python-based *xgboost* were explored, with the latter  
449 being used in the end with appropriate binning since it shows slightly higher performance.  
450 The sections will be structured as follows. A list of variables used as inputs to the BDT will be  
451 motivated and described, followed by the hyperparameters used in the training of the decision  
452 tree. The end-result comparison with the nominal cut-based analysis will be presented, and  
453 studies contributing to this result (e.g., training sets/parameters, etc) will follow.

### 454 6.2 Selection and variables

455 BDT training is susceptible to low statistics, so in order to increase statistics, a relaxed baseline  
456 selection for training was created, consisting of

- 457 • Lepton 1  $p_T > 15$
- 458 • Lepton 2  $p_T > 15$
- 459 •  $H_T > 250$
- 460 • MET  $> 30$
- 461 • Njets  $\geq 2$
- 462 • Nbtags  $\geq 1$

463 Signal was defined to be from four top MC and background was taken as all of the SM back-  
464 grounds previously described and taken in this analysis. Variations of this are discussed in the  
465 last section of miscellaneous studies.

466 Over 40 variables were considered. Ranked in approximate descending order of discriminative  
467 power, as reported by TMVA, they are

- 468 • Njets
- 469 • Nbtags
- 470 • Nleps
- 471 •  $m_T(\ell_2)$
- 472 •  $m(\ell_1, \ell_2)$
- 473 • MET
- 474 •  $m_T(\ell_1)$
- 475 •  $H_T^b$ :  $H_T$  made from b jets
- 476 •  $m(j_1, j_2)$ : invariant mass of leading two jets
- 477 •  $m(\ell_1, j_2)$

- 478 •  $\Delta\phi(j_1, j_2)$
- 479 •  $\Delta\phi(\ell_1, j_1)$
- 480 •  $p_T(j_i)$  for  $i = 1 - 8$
- 481 •  $q_1$ : sign of the same-sign lepton pair
- 482 •  $\Delta\eta(\ell_1, \ell_2)$
- 483 •  $H_T$
- 484 •  $H_T^{\text{ratio}}$ : Ratio of  $H_T$  of first four leading jets to rest
- 485 •  $m(\ell_1, j_1)$
- 486 • Nlooseb: number of btags passing loose threshold
- 487 • Ntightb: number of btags passing tight threshold
- 488 •  $\max(m(j)/p_T)$ : ratio of jet mass to momentum to discriminate merged jets
- 489 • Wcands: number of jet pairs with invariant mass within 30GeV of the W mass

490 Other more general variables like  $m(\ell_i, j_j)$  were also considered but showed little to no discrimination power, where i and j encompass the first two leading objects. The same generalization applied to other indexed variables. Note that this ranking takes into account the correlation between variables, explaining why  $H_T$  is not as highly ranked, since discriminative power is first taken from Njets. Past approximately 22 variables, no extra AUC (area under curve) was gained, so it was determined to proceed with this smaller set of variables for simplicity. The AUC metric is the area under the signal versus background efficiency Receiver Operating Characteristic (ROC) curve, bounded by 0 and 1, where 1 is equivalent to perfect discrimination between signal and background, and 0.5 represents discrimination no better than random guessing.

500 The 19 approximately most performant variables were then selected to continue optimization.  
501 They are

- 502 • (a) Nbtags
- 503 • (b) Njets
- 504 • (c) Nlooseb
- 505 • (d) MET
- 506 • (e) Ntightb
- 507 • (f)  $p_T(\ell_2)$
- 508 • (g)  $m(\ell_1, j_1)$
- 509 • (h)  $p_T(j_1)$
- 510 • (i)  $p_T(j_7)$
- 511 • (j)  $\Delta\phi(\ell_1, \ell_2)$
- 512 • (k)  $p_T(j_6)$
- 513 • (l)  $\max(m(j)/p_T(j))$
- 514 • (m) Nleps
- 515 • (n)  $p_T(\ell_1)$
- 516 • (o)  $\Delta\eta(\ell_1, \ell_2)$
- 517 • (p)  $p_T(j_8)$
- 518 • (q)  $H_T^b$

- 519     • (r)  $p_T(\ell_3)$   
520     • (s)  $q_1$

521 Kinematic distributions for these input variables are shown in Figure 6.

### 522 6.3 Hyperparameters

523 Hyper-parameter tuning was performed in order to maximize discrimination ( $s/\sqrt{s+b}$  and  
524 AUC) while avoiding over-training from limited statistics. The selected set of TMVA hyper-  
525 parameters is given by the strings

- 526     • NTrees=500  
527     • nEventsMin=150  
528     • MaxDepth=5  
529     • BoostType=AdaBoost  
530     • AdaBoostBeta=0.25  
531     • SeparationType=GiniIndex  
532     • nCuts=20  
533     • PruneMethod=NoPruning

534 In total, the complete set of combinations considered in the hyperparameter scan were

- 535     • NTrees=200,500,1000  
536     • nEventsMin=50,150,300  
537     • MaxDepth=4,5,6  
538     • BoostType=AdaBoost,Bagging,Grad  
539     • AdaBoostBeta=0.1,0.25,0.8  
540     • SeparationType=GiniIndex,CrossEntropy  
541     • nCuts=5,20,100

542 A similar hyperparameter scan for xgboost yielded the parameters

- 543     • n\_estimators = 500  
544     • eta = 0.07  
545     • max\_depth = 5  
546     • subsample = 0.6  
547     • alpha = 8.0  
548     • gamma = 2.0  
549     • lambda = 1.0  
550     • min\_child\_weight = 1.0  
551     • colsample\_bytree = 1.0

552 Note that “n\_estimators” represents the number of trees, and “eta” is the learning rate. In  
553 particular, for a given learning rate, the number of trees and the depth of each tree are the most  
554 impactful hyper-parameters.

555 Using these optimal hyperparameters, Figure 7 compares ROC curves and  $s/\sqrt{s+b}$  curves for  
556 TMVA and xgboost, showing that xgboost yields a significant performance than TMVA.

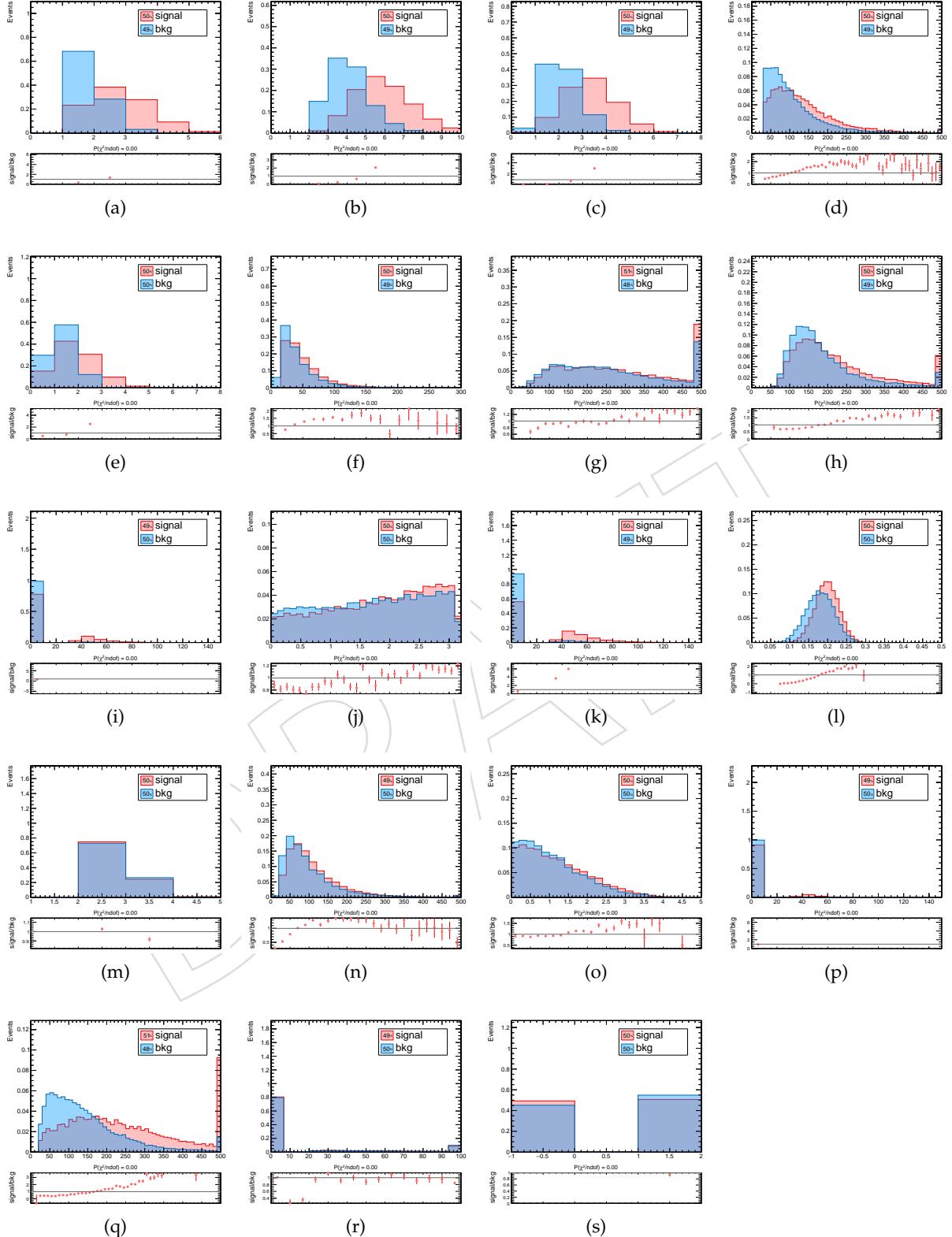


Figure 6: Distributions of kinematic inputs to the BDT. Integrals have been normalized so only shapes are relevant here.

557 Figure 8 shows that xgboost discriminator shapes for train and test sets separately for back-  
 558 ground and signal. There appears to be no signs of overtraining, given the above hyperparam-  
 559 eters.

560 The raw xgboost output discriminant using these 19 variables is shown in Figure 9, lumping  
 561 rare MC processes into "Rares", and fakes/flips into "Others".

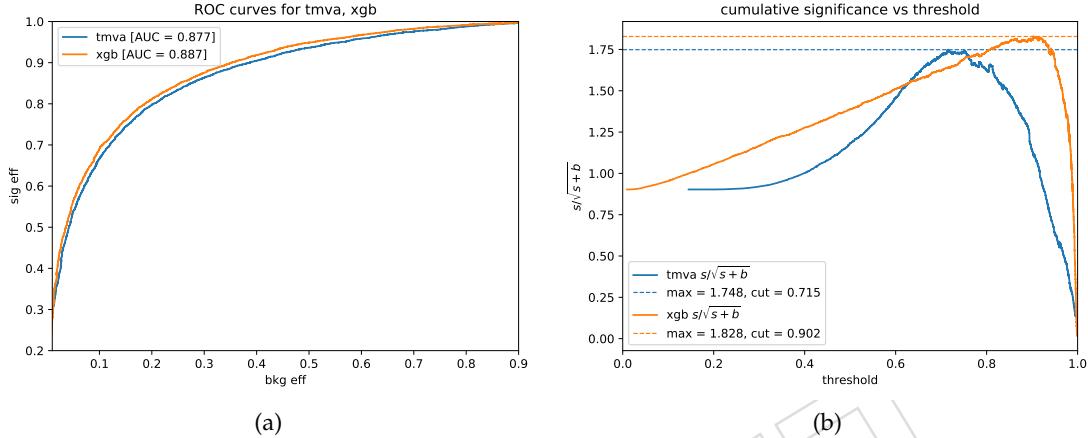


Figure 7: Left: ROC curves for TMVA and xgboost with AUC metric displayed in the legend. Right: Cumulative significance vs threshold value on the discriminant for TMVA and xgboost. The maximum of this curve is marked by a dotted line.

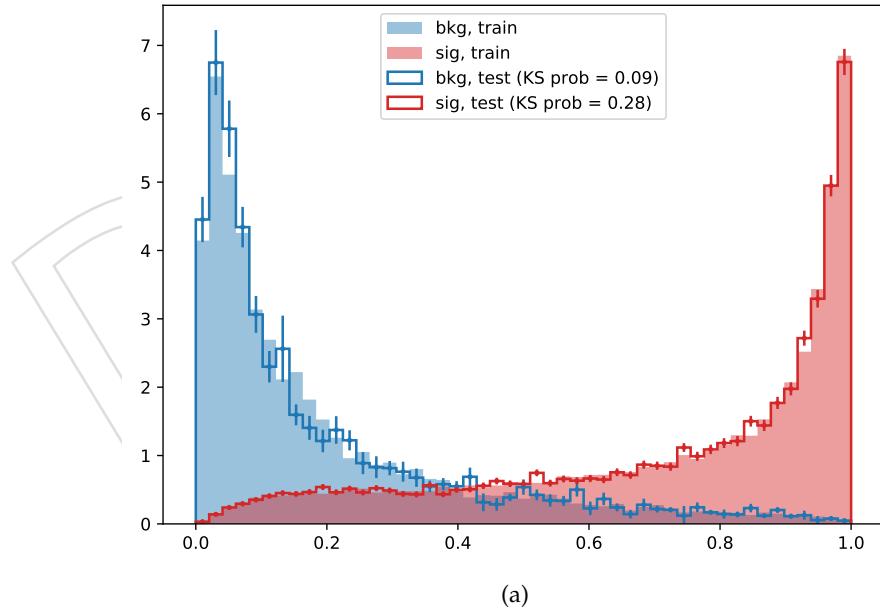


Figure 8: Raw xgboost discriminator shape for train and test sets separately for background and signal.

#### 562 6.4 Comparison with cut-based

563 For 18 bins in total, there are 17 bins along the discriminator output from 0 to 1 corresponding  
 564 to the bin edges of 0.0000, 0.0362, 0.0659, 0.1055, 0.1573, 0.2190, 0.2905,  
 565 0.3704, 0.4741, 0.6054, 0.7260, 0.8357, 0.9076, 0.9506, 0.9749, 0.9884,

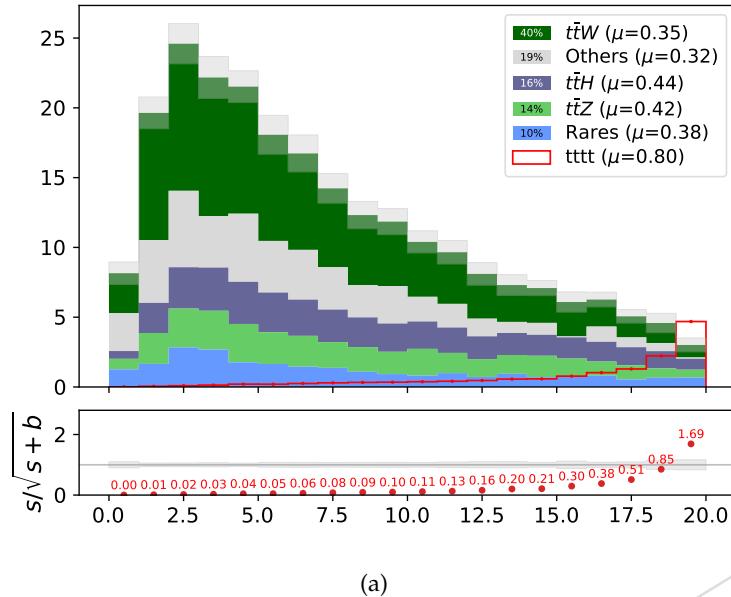


Figure 9: Raw xgboost BDT discriminator shape with a panel showing the cumulative  $s/\sqrt{s+b}$  with MC scaled to  $75 \text{ fb}^{-1}$ .

566    0.9956, 1.0000 . The initial bin contains events matching the CRZ selection to help constrain  
 567    the  $t\bar{t}Z$  background. Increasing the binning beyond this does not result in a gain in significance, and also gives issues with sample statistics.

569    Note that the binning was chosen to create a shape different from the raw xgboost discriminator  
 570    output. The resulting shape is similar to the output of an AdaBoost algorithm from TMVA, and  
 571    has the effect of distributing signal events normally contained within 1-2 bins into a few more,  
 572    allowing for a better constraint in the fit.

573    In order to compare properly with the cut-based, the region in which the BDT discriminant is  
 574    calculated and applied corresponds to the nominal baseline selection. The 18-bin BDT shown  
 575    in Figure 10 yields an expected significance of  $2.68\sigma$ , with the 16-bin cut-based result yielding  
 576    an expected significance of  $2.46\sigma$ , which is a gain slightly larger than 8%.

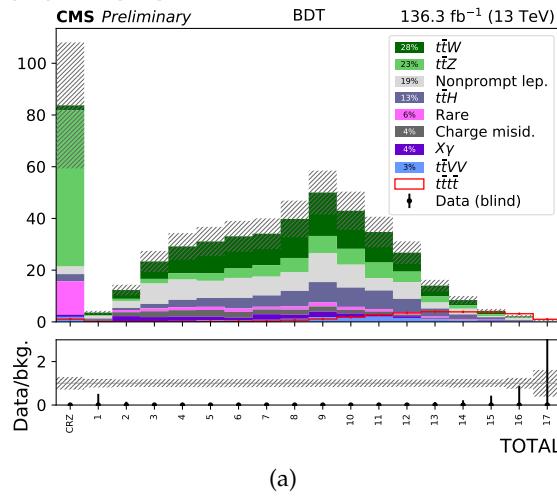


Figure 10: BDT discriminant in discrete signal bins, including CRZ in the first bin.

577 **6.5 Misc studies, all numbers and text is from 2016**

578 **6.5.1 Definition of background**

579 Training is performed on the full mix of backgrounds. Backgrounds modeled with simulation  
 580 are normalized to a common luminosity, i.e. weighted events are used. For comparison, we  
 581 also trained separate BDTs for each of the ttV backgrounds, but saw almost no gain in expected  
 582 sensitivity. Note that signal-background have similar separation for each of the individual  
 583 background processes, as shown in Figure 11.

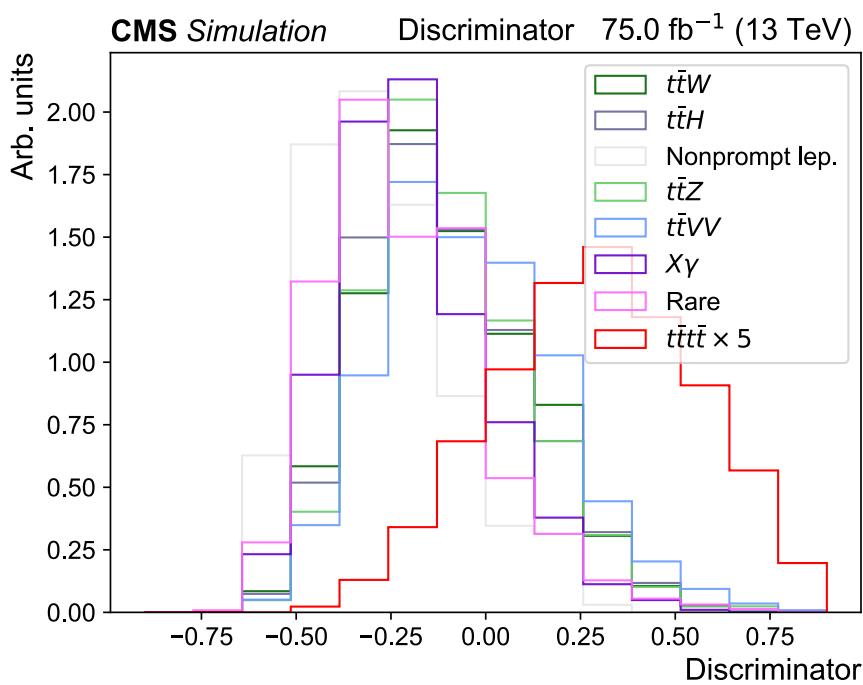


Figure 11: Unstacked BDT discriminator shape for signal and background process

## 584 7 Validation of kinematic variables

585 In this section we describe the control-regions to validate with data the kinematic variables we  
 586 use in defining our cut-based signal regions:  $H_T$ ,  $\cancel{E}_T$ ,  $N_{\text{jets}}$ ,  $N_{\text{b jets}}$ . In addition, the variables  
 587 used for the BDT are also studied:  $H_T^b$ , Nlooseb, Ntightb,  $\Delta\phi(\ell_1, \ell_2)$ ,  $\Delta\eta(\ell_1, \ell_2)$ ,  $p_T(\ell_1)$ ,  $p_T(\ell_2)$ ,  
 588  $p_T(\ell_3)$ ,  $q_1$ ,  $m(\ell_1, j_1)$ ,  $\max(m(j)/p_T(j))$  and the  $p_T$  for jets 1,6,7, and 8.

589 The control regions considered are opposite-sign dilepton (Section 7.1), tight/loose same-sign  
 590 dilepton (Section 7.2). CRW and CRZ, which are closest to the signal region, will be presented  
 591 in Section 10.

### 592 7.1 Opposite-sign dilepton events

593 In this control region, the same requirements on  $H_T$ ,  $\cancel{E}_T$  and  $N_{\text{jets}}$  as in our inclusive *baseline*  
 594 selection are applied (see Section 5), but we require two opposite sign tight leptons and we  
 595 remove the  $Z/\gamma^*$  veto. This control region coincides with the application region we use for the  
 596 data-driven method to estimate the *Charge flips* background (see Section 8).

597 Distributions are shown in Figs. 12,13,14 (2017, 2018, 2016+2017+2018) for the main variables,  
 598 and Figs. 15,16,17 for additional variables used in the BDT. The overall agreement is very good,  
 599 but there are some discrepancies in the  $N_{\text{b jets}}$  and BDT distributions that we will monitor as  
 600 we include additional corrections to the MC. The ISR/FSR corrections described in Section 2,  
 601 corresponding to Figure 4, is applied to the  $t\bar{t}$  sample here. Note that quantities like the  $p_T$   
 602 for jet 7,8 correspond to high (5+) ISR/FSR jet multiplicities in OS  $t\bar{t}$ . These are anyway not  
 603 directly probed in the signal regions, consisting mainly of  $t\bar{t}W$  and  $t\bar{t}Z$  for which the matrix  
 604 element provides more partons.

### 605 7.2 Same-sign tight+fail dilepton events

606 In this control region, the same requirements on  $H_T$ ,  $\cancel{E}_T$  and  $N_{\text{jets}}$  as in our inclusive *baseline*  
 607 selection are applied (see Section 5), but we require one tight lepton and one same-sign lepton  
 608 failing the tight requirement. The control region is enriched in events with one fake lepton.  
 609 It corresponds to the application region for the *fake-rate* method, the data-driven estimate of  
 610 backgrounds with fake-leptons (see Section 8). The only difference is that in the fake-rate ap-  
 611 plication we also allow events where the two leptons fail the tight requirements.

612 Distributions are shown in Figs. 18,19,20 (2017, 2018, 2016+2017+2018) for the main variables,  
 613 and Figs. 21,22,23 for additional variables used in the BDT. There is an overall underestimate,  
 614 consistent with what was seen in 2016. The underestimate is flavor-independent, and does not  
 615 show large trends in the main kinematics. Such an underestimate in loose-not-tight leptons, if  
 616 paired with a data FakeRate equal or larger to the MC one, can indicate that the Nonprompt  
 617 lepton background is underestimated by a pure MC prediction. Since this background is pre-  
 618 dicted from data, that is not an issue with the analysis strategy.

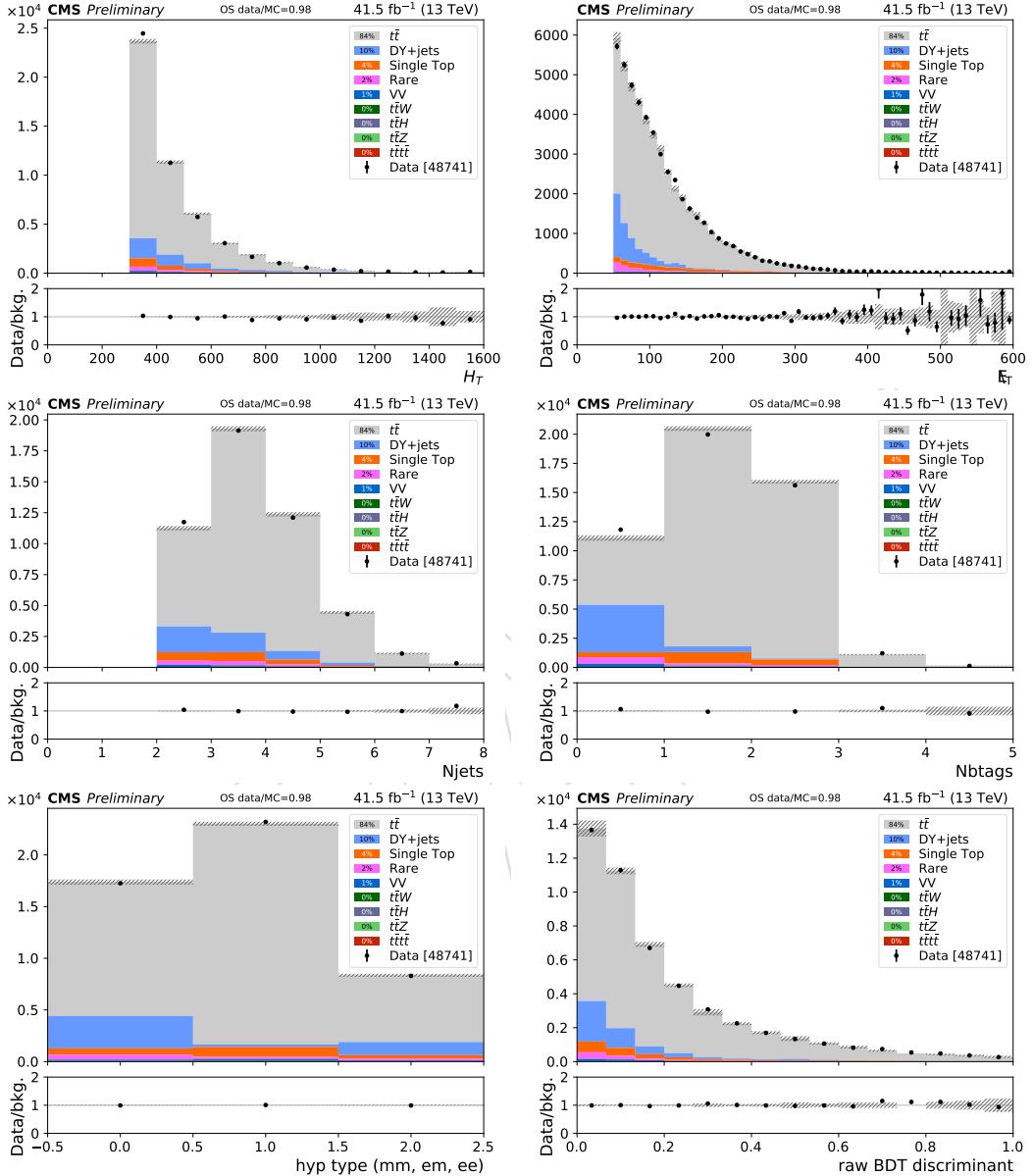


Figure 12: Data to simulation comparisons for 2017. From left top to right bottom the  $H_T$ ,  $\cancel{E}_T$ ,  $N_{\text{jets}}$ ,  $N_{\text{btags}}$ , lepton flavor and raw BDT discriminant distributions are shown for the opposite-sign dilepton baseline region. Shaded band shows MC stat uncertainty, except in the case of the discriminator distribution, where it has been added in quadrature with scale, btag, JEC, and JER variations

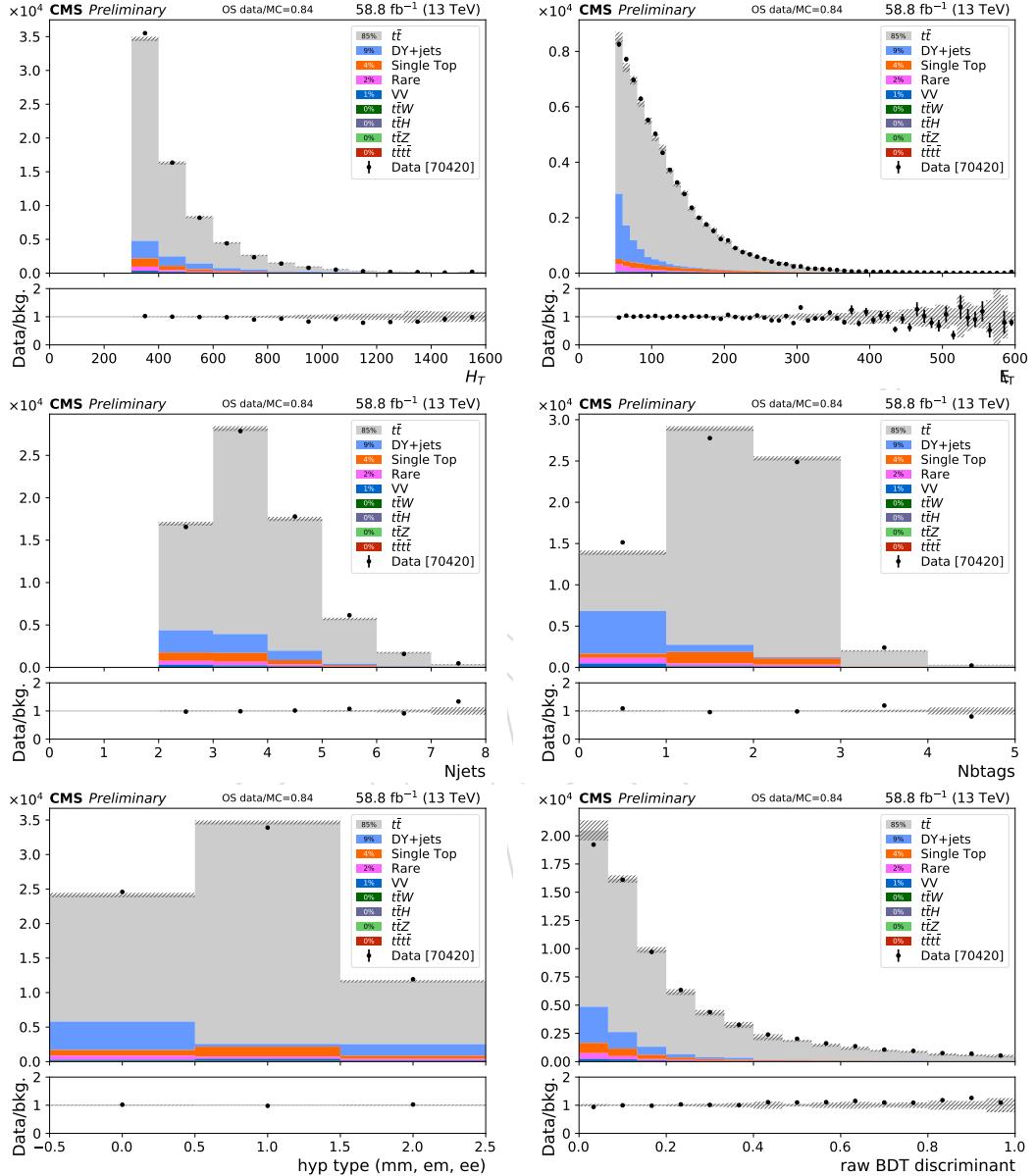


Figure 13: Data to simulation comparisons for 2018. From left top to right bottom the  $H_T$ ,  $\cancel{E}_T$ ,  $N_{\text{jets}}$ ,  $N_{\text{btags}}$ , lepton flavor and raw BDT discriminant distributions are shown for the opposite-sign dilepton baseline region. Shaded band shows MC stat uncertainty, except in the case of the discriminator distribution, where it has been added in quadrature with scale, btag, JEC, and JER variations.

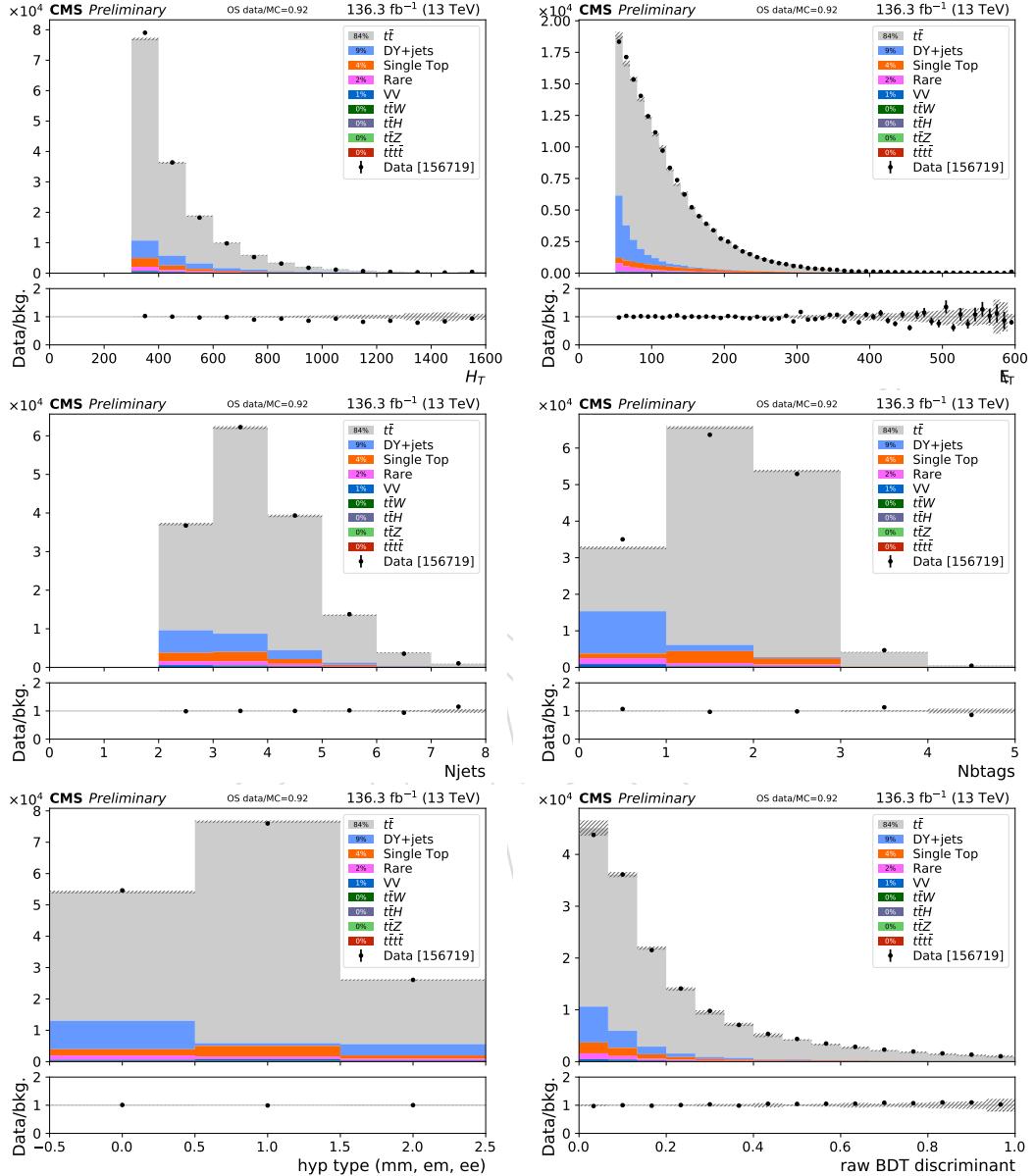


Figure 14: Data to simulation comparisons for 2016+2017+2018. From left top to right bottom the  $H_T$ ,  $\not{E}_T$ ,  $N_{\text{jets}}$ ,  $N_{\text{btags}}$ , lepton flavor and raw BDT discriminant distributions are shown for the opposite-sign dilepton baseline region. Shaded band shows MC stat uncertainty, except in the case of the discriminator distribution, where it has been added in quadrature with scale, btag, JEC, and JER variations

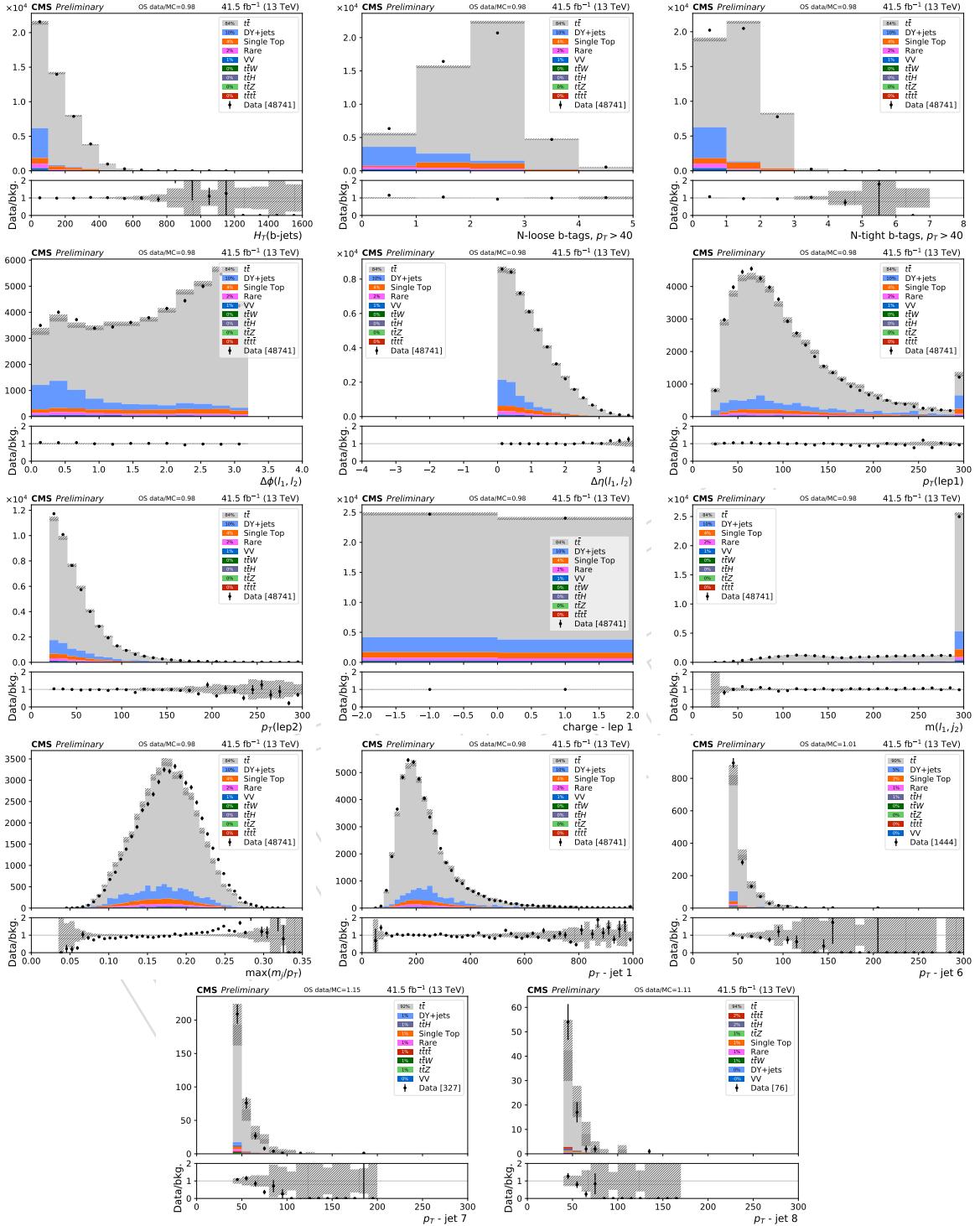


Figure 15: Data to simulation comparisons for 2017, for the additional variables used by the BDT. From left top to right bottom,  $H_T^b$ , Nlooseb, Ntightb,  $\Delta\phi(\ell_1, \ell_2)$ ,  $\Delta\eta(\ell_1, \ell_2)$ ,  $p_T(\ell_1)$ ,  $p_T(\ell_2)$ ,  $p_T(\ell_3)$ ,  $q_1$ ,  $m(\ell_1, j_1)$ ,  $\max(m(j)/p_T(j))$  and the  $p_T$  for jets 1,6,7, and 8; shown for the opposite-sign dilepton baseline region.

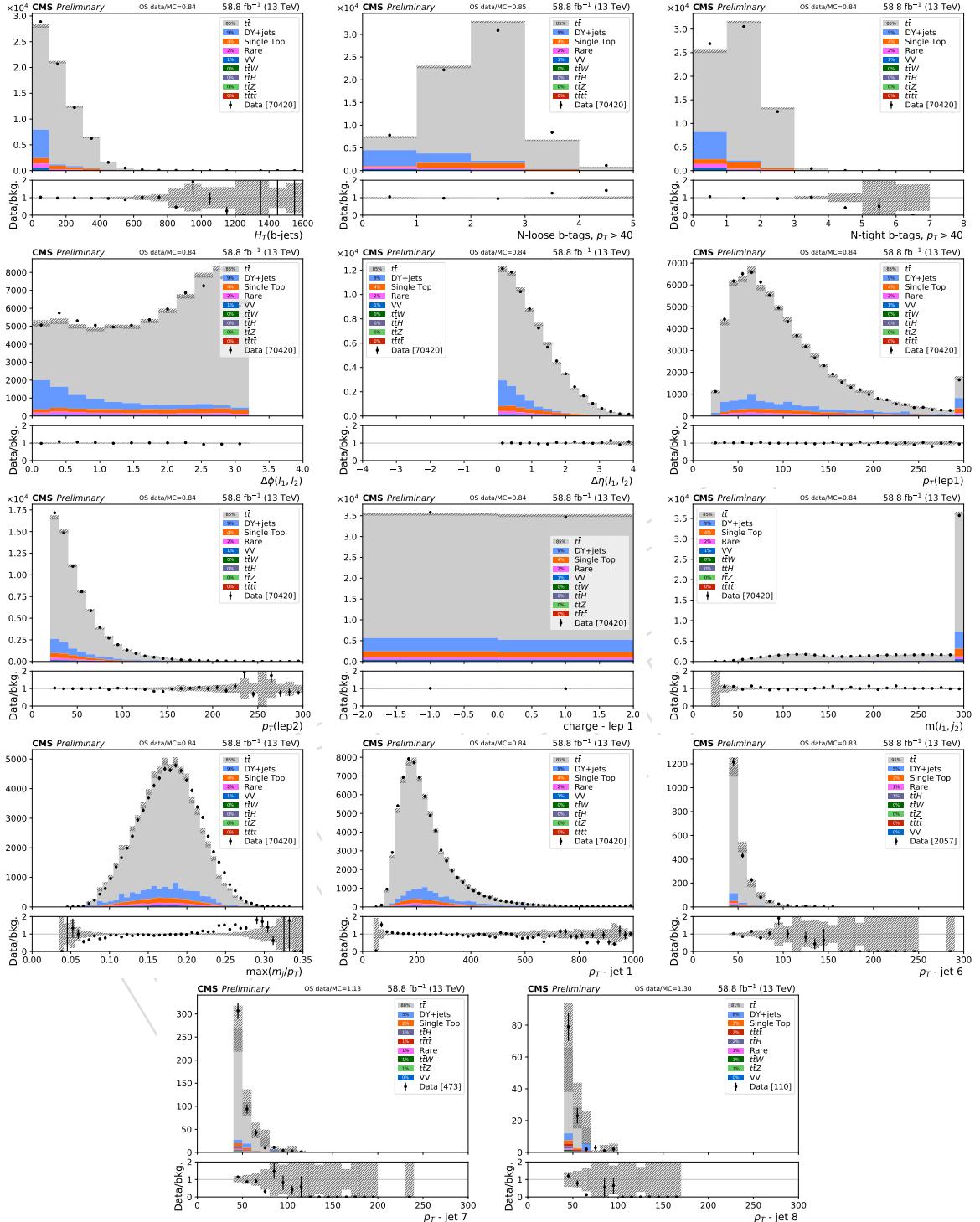


Figure 16: Data to simulation comparisons for 2018, for the additional variables used by the BDT. From left top to right bottom,  $H_T^b$ , Nlooseb, Ntightb,  $\Delta\phi(\ell_1, \ell_2)$ ,  $\Delta\eta(\ell_1, \ell_2)$ ,  $p_T(\ell_1)$ ,  $p_T(\ell_2)$ ,  $p_T(\ell_3)$ ,  $q_1$ ,  $m(\ell_1, j_1)$ ,  $\max(m(j)/p_T(j))$  and the  $p_T$  for jets 1,6,7, and 8; shown for the opposite-sign dilepton baseline region.

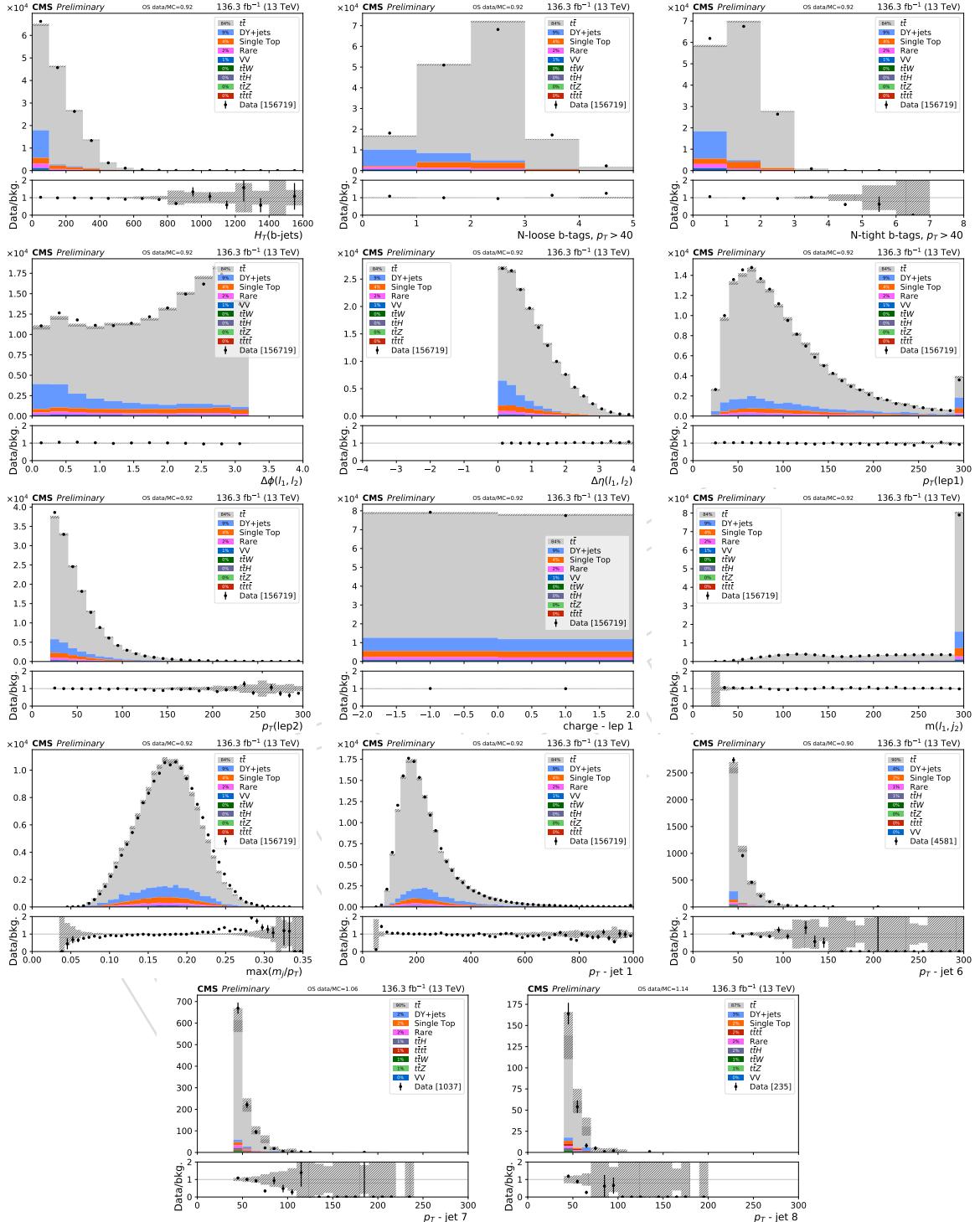


Figure 17: Data to simulation comparisons for 2016+2017+2018, for the additional variables used by the BDT. From left top to right bottom,  $H_T^b$ , Nlooseb, Ntightb,  $\Delta\phi(\ell_1, \ell_2)$ ,  $\Delta\eta(\ell_1, \ell_2)$ ,  $p_T(\ell_1)$ ,  $p_T(\ell_2)$ ,  $p_T(\ell_3)$ ,  $q_1$ ,  $m(\ell_1, j_1)$ ,  $\max(m(j)/p_T(j))$  and the  $p_T$  for jets 1,6,7, and 8; shown for the opposite-sign dilepton baseline region.

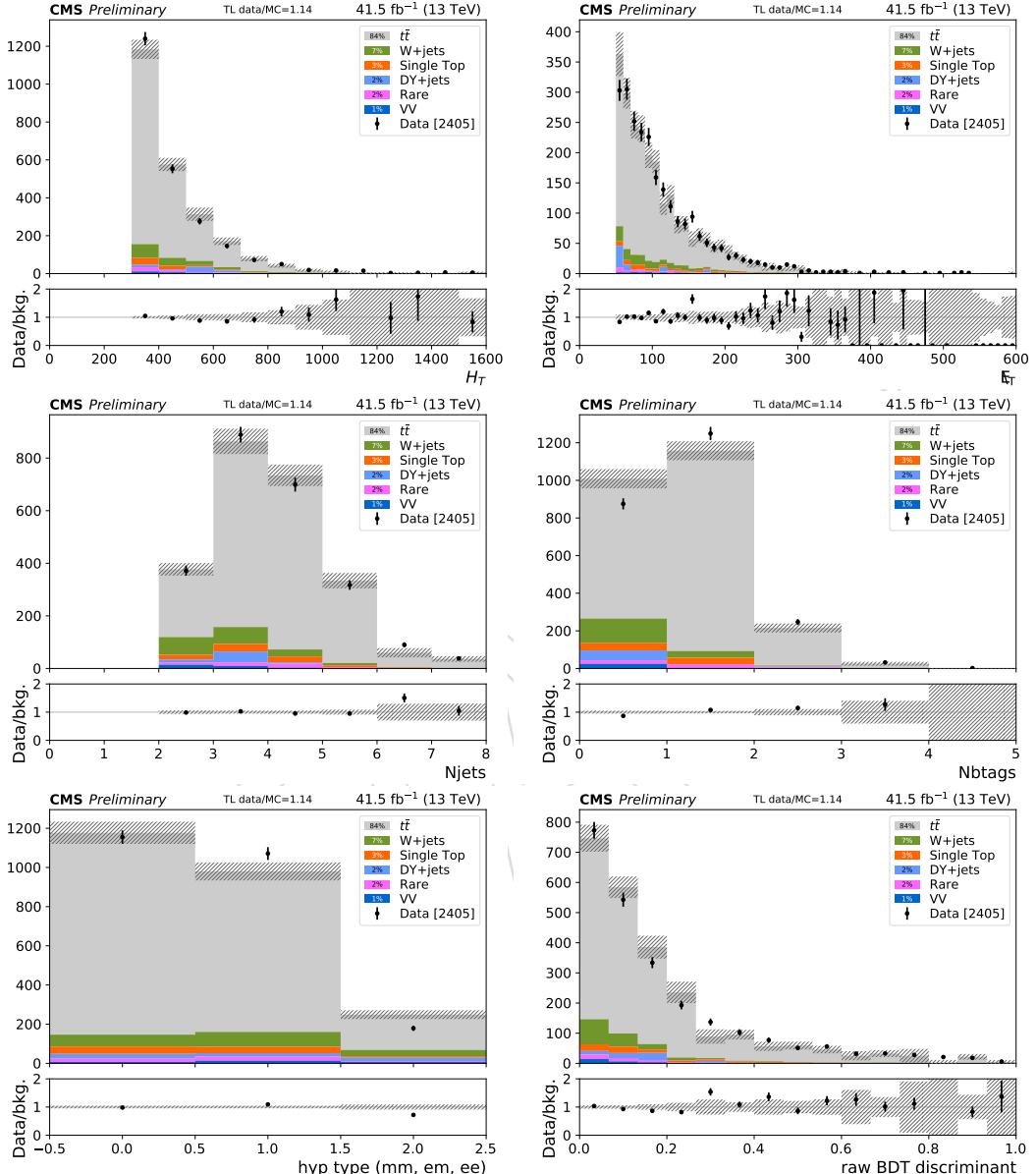


Figure 18: Data to simulation comparisons for 2017. From left top to right bottom the  $H_T$ ,  $\cancel{E}_T$ ,  $N_{\text{jets}}$ ,  $N_{\text{bjets}}$ , lepton flavor and raw BDT discriminant distributions are shown for the same-sign tight+fail dilepton baseline region. Shaded band shows MC stat uncertainty, except in the case of the discriminator distribution, where it has been added in quadrature with scale, btag, JEC, and JER variations

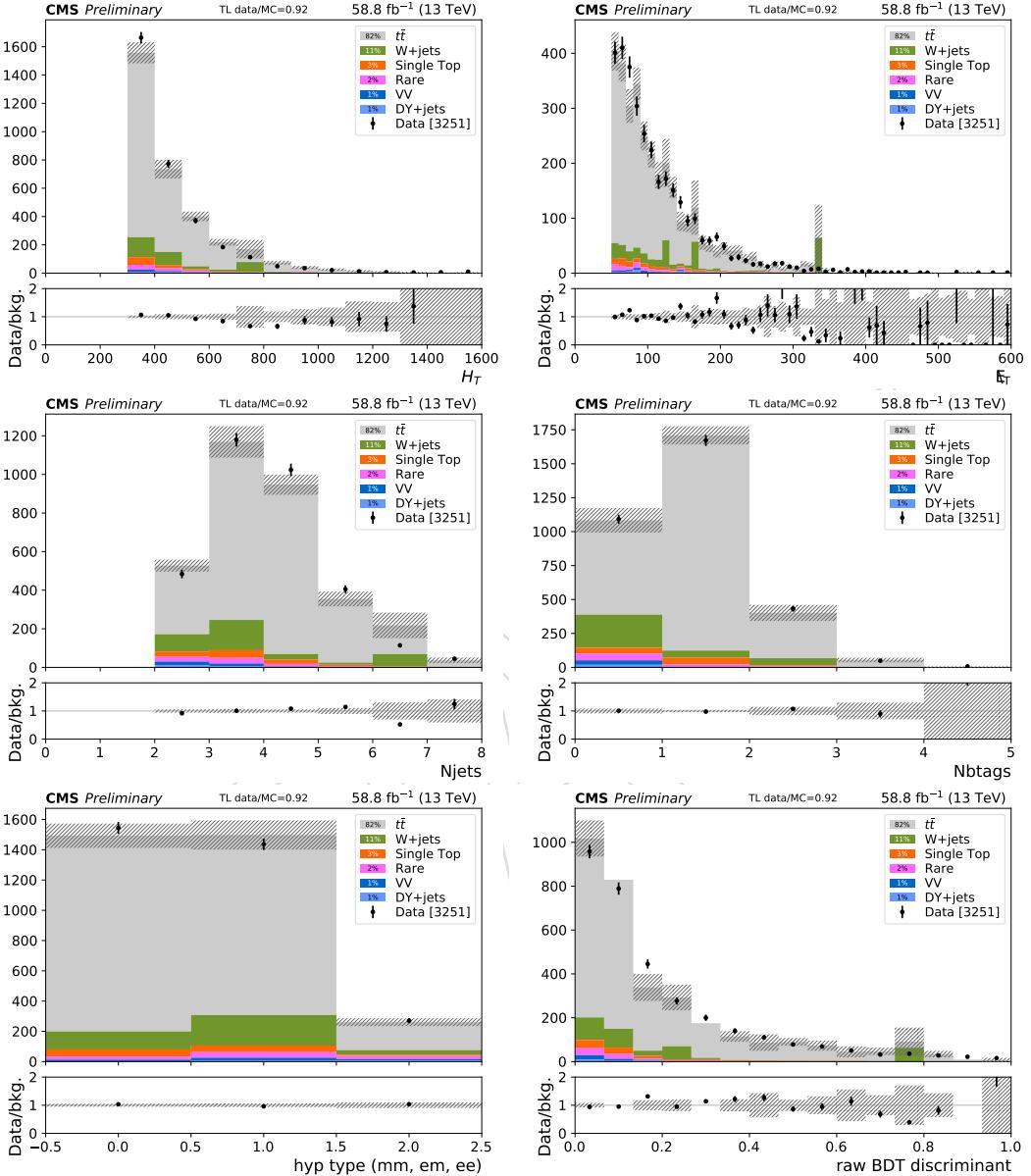


Figure 19: Data to simulation comparisons for 2018. From left top to right bottom the  $H_T$ ,  $\cancel{E}_T$ ,  $N_{\text{jets}}$ ,  $N_{\text{bjets}}$ , lepton flavor and raw BDT discriminant distributions are shown for the same-sign tight+fail dilepton baseline region. Shaded band shows MC stat uncertainty, except in the case of the discriminator distribution, where it has been added in quadrature with scale, btag, JEC, and JER variations

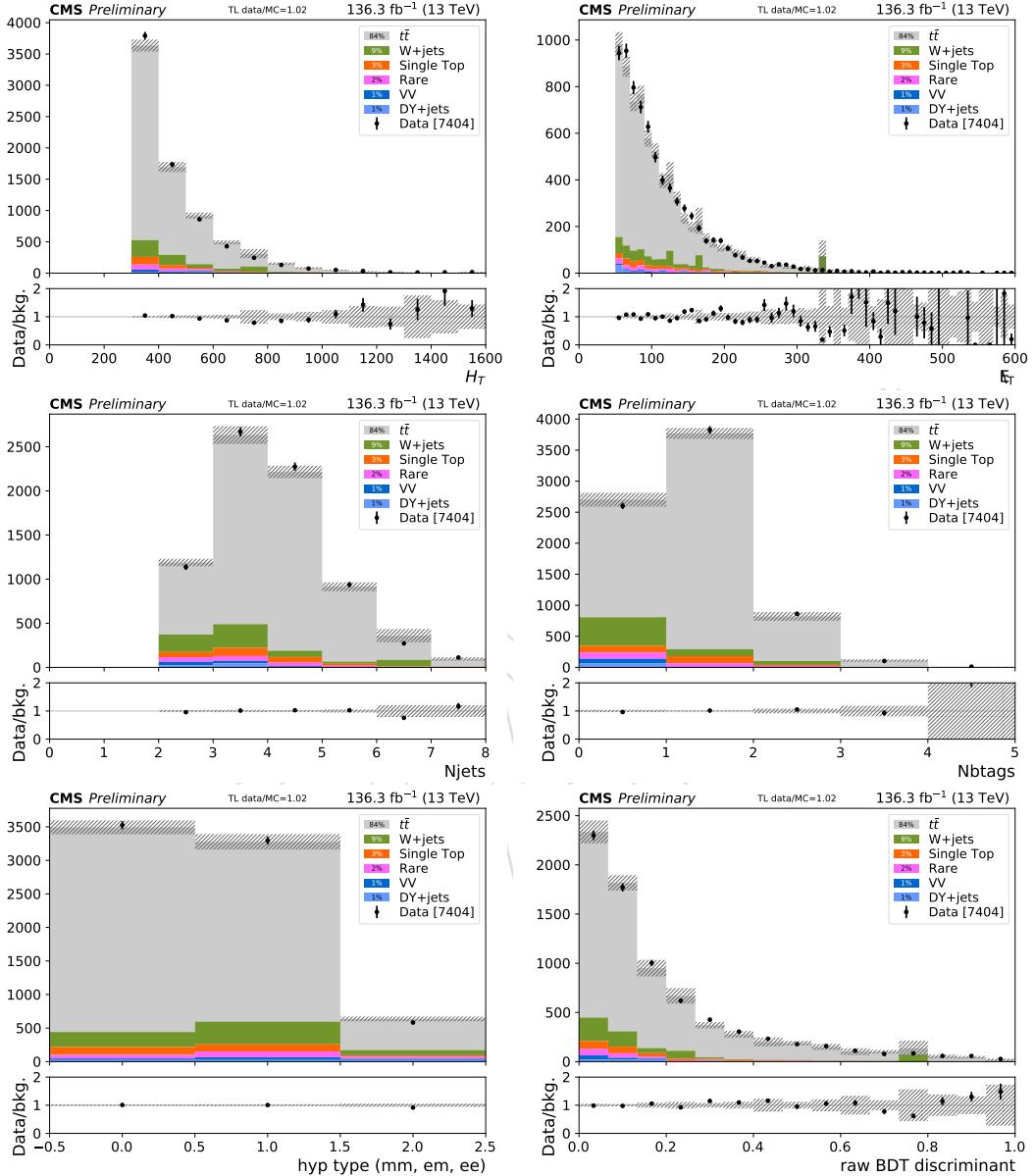


Figure 20: Data to simulation comparisons for 2016+2017+2018. From left top to right bottom the  $H_T$ ,  $\cancel{E}_T$ ,  $N_{\text{jets}}$ ,  $N_{\text{btags}}$ , lepton flavor and raw BDT discriminant distributions are shown for the same-sign tight+fail dilepton baseline region. Shaded band shows MC stat uncertainty, except in the case of the discriminator distribution, where it has been added in quadrature with scale, btag, JEC, and JER variations

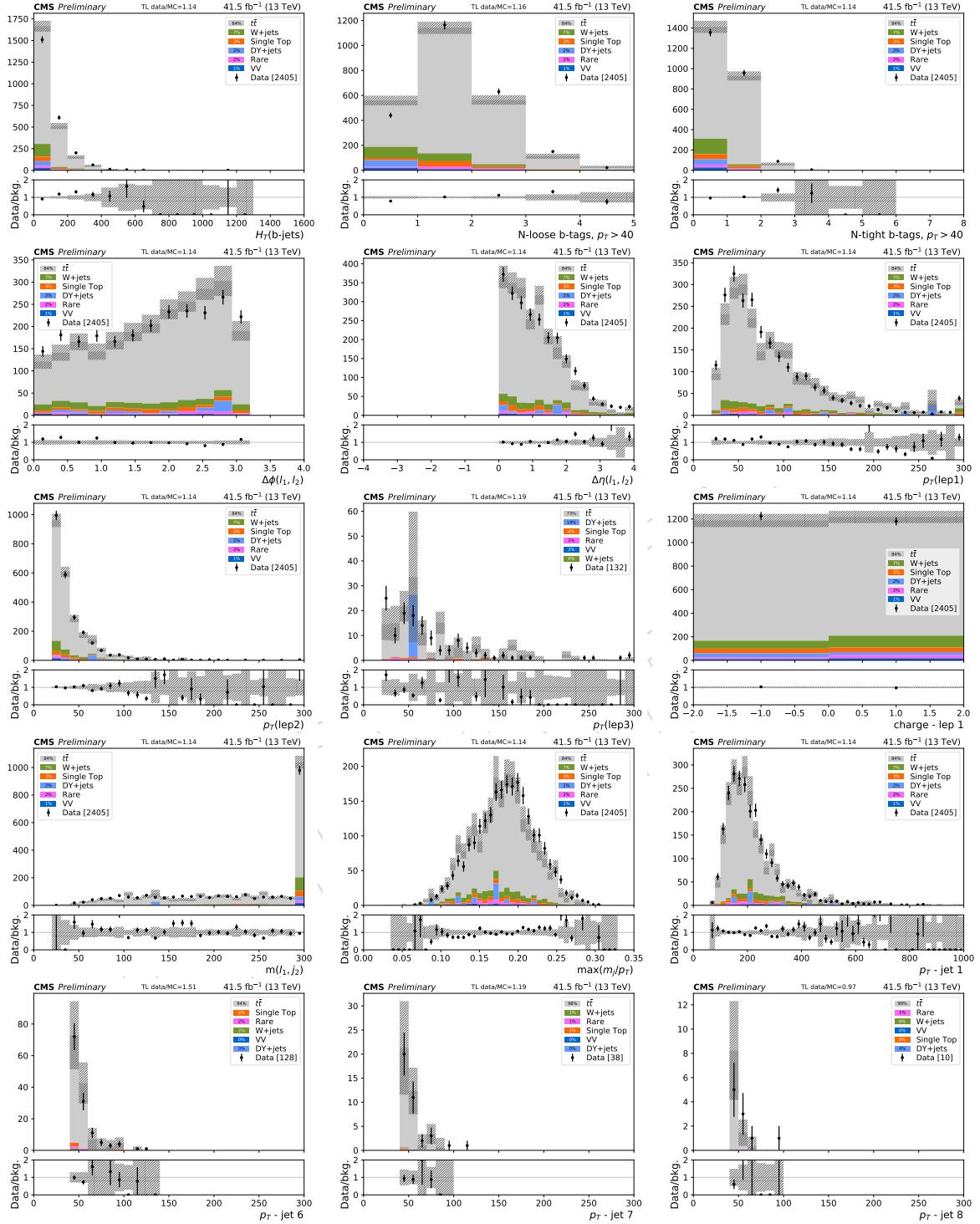


Figure 21: Data to simulation comparisons for 2017, for the additional variables used by the BDT. From left top to right bottom,  $H_T^b$ ,  $N_{\text{loose}b}$ ,  $N_{\text{tight}b}$ ,  $\Delta\phi(\ell_1, \ell_2)$ ,  $\Delta\eta(\ell_1, \ell_2)$ ,  $p_T(\ell_1)$ ,  $p_T(\ell_2)$ ,  $p_T(\ell_3)$ ,  $q_1$ ,  $m(\ell_1, j_1)$ ,  $\max(m(j)/p_T(j))$  and the  $p_T$  for jets 1,6,7, and 8; shown for the same-sign tight+fail dilepton baseline region.

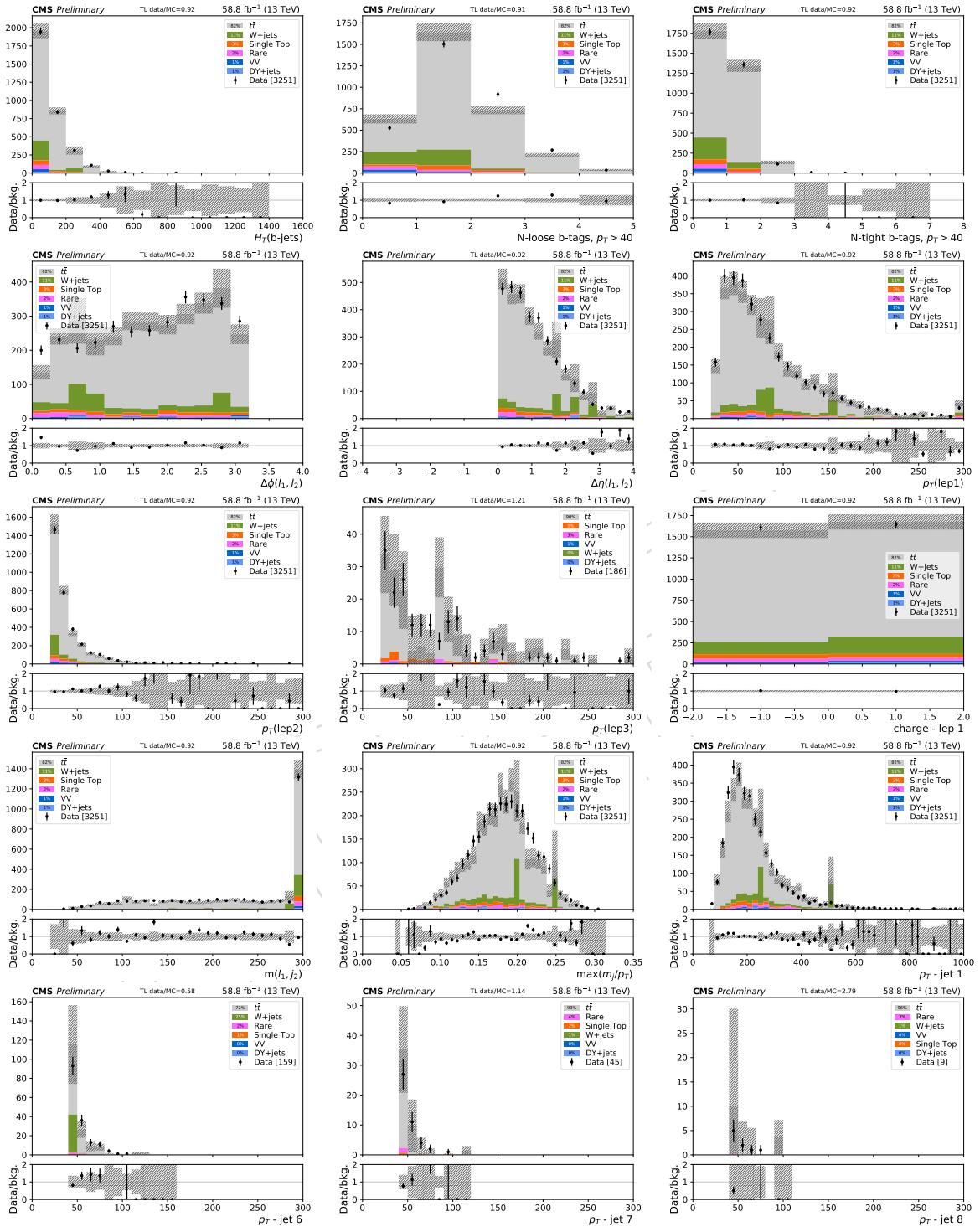


Figure 22: Data to simulation comparisons for 2018, for the additional variables used by the BDT. From left top to right bottom,  $H_T^b$ , Nlooseb, Ntightb,  $\Delta\phi(\ell_1, \ell_2)$ ,  $\Delta\eta(\ell_1, \ell_2)$ ,  $p_T(\ell_1)$ ,  $p_T(\ell_2)$ ,  $p_T(\ell_3)$ ,  $q_1$ ,  $m(\ell_1, j_1)$ ,  $\max(m(j)/p_T(j))$  and the  $p_T$  for jets 1,6,7, and 8; shown for the same-sign tight+fail dilepton baseline region.

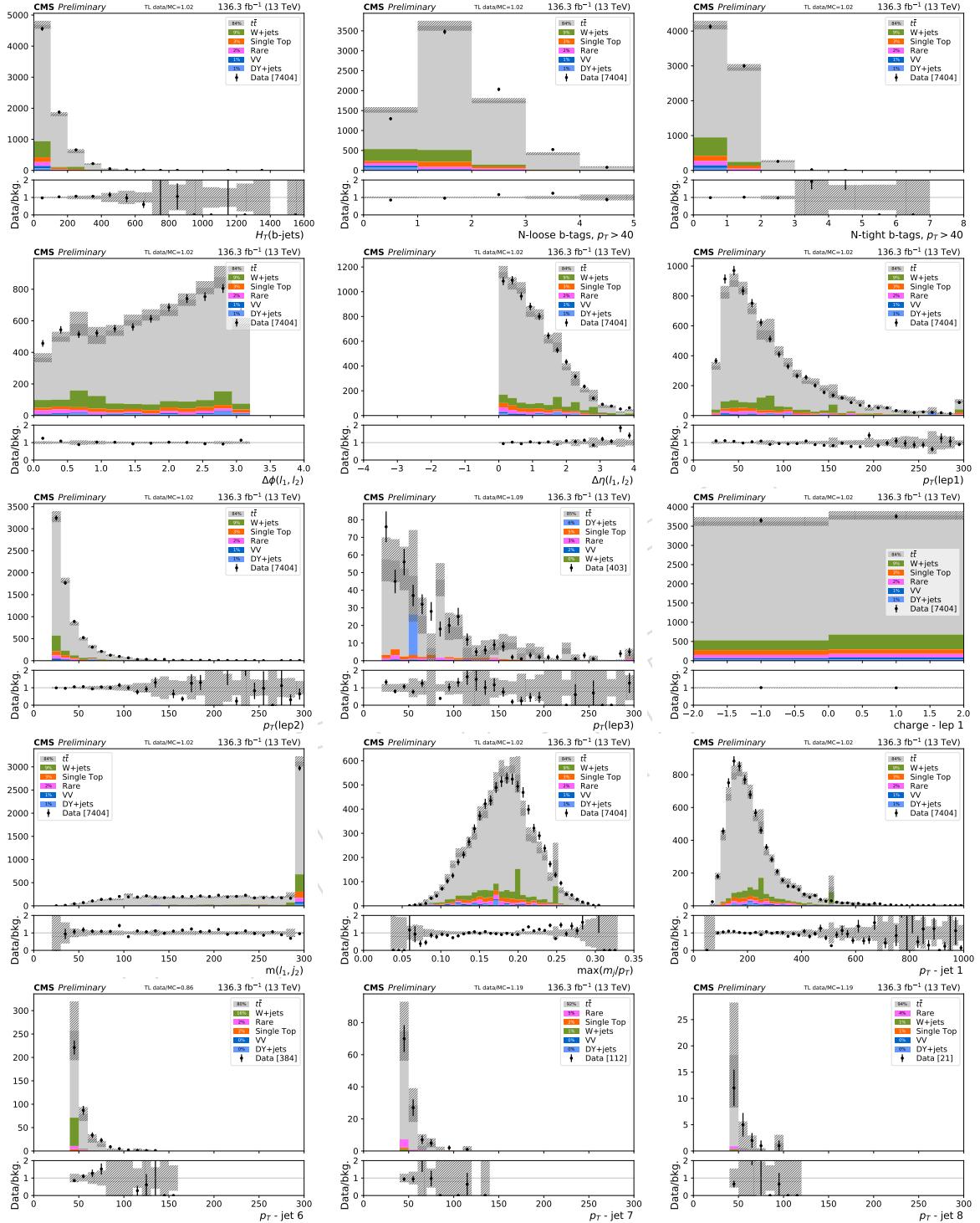


Figure 23: Data to simulation comparisons for 2016+2017+2018, for the additional variables used by the BDT. From left top to right bottom,  $H_T^b$ ,  $N_{\text{loose} b}$ ,  $N_{\text{tight} b}$ ,  $\Delta\phi(\ell_1, \ell_2)$ ,  $\Delta\eta(\ell_1, \ell_2)$ ,  $p_T(\ell_1)$ ,  $p_T(\ell_2)$ ,  $p_T(\ell_3)$ ,  $q_1$ ,  $m(\ell_1, j_1)$ ,  $\max(m(j)/p_T(j))$  and the  $p_T$  for jets 1,6,7, and 8; shown for the same-sign tight+fail dilepton baseline region.

### 619 7.3 Fake-enriched validation region in data

620 In this control region, the same requirements on  $H_T$ ,  $\cancel{E}_T$  and  $N_{\text{jets}}$  as in our inclusive *baseline* se-  
 621 lection are applied (see Section 5), except we relax the  $H_T$  requirement and require  $N_b \text{ jets} = 1$ .  
 622 This region has a significant non-prompt component and allows us to check the overall clo-  
 623 sure of the method in data. In the plots, the fake background is data-driven. Distributions are  
 624 shown in Figs. 24. The overall data/MC normalization factor in this region is 1.06. If fakes are  
 625 entirely responsible for this discrepancy, and given that fakes constitute half of the background,  
 626 this represents a 12% normalization increase of fakes, well within the 30% normalization un-  
 627 certainty taken on this process.

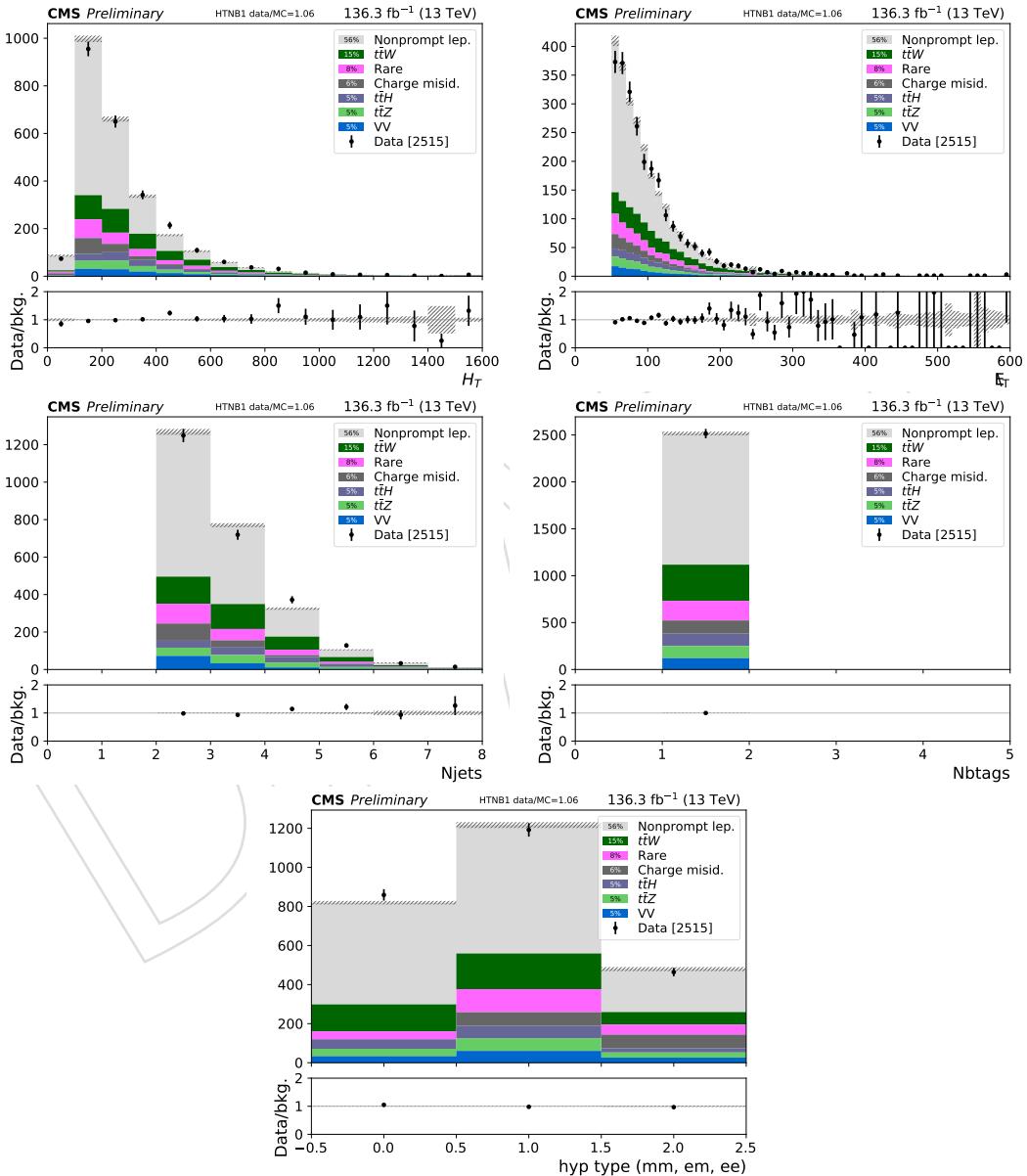


Figure 24: Data to simulation comparisons for 2016+2017+2018. From left top to right bottom the  $H_T$ ,  $\cancel{E}_T$ ,  $N_{\text{jets}}$ ,  $N_b \text{ jets}$ , lepton flavor and raw BDT discriminant distributions are shown for the same-sign dilepton region with  $N_b \text{ jets} = 1$ ,  $N_{\text{jets}} \geq 2$ ,  $E_T^{\text{miss}} \geq 50$ . Shaded band shows MC stat uncertainty.

## 628 8 Background Estimations

629 Backgrounds for the same-sign dilepton final state can be divided in three categories:

- 630 • **Fake leptons:** “Non-Prompt” or “Fake” leptons are leptons from heavy-flavour de-  
631 cays, misidentified hadrons, muons from light-meson decays in flight, or electrons  
632 from unidentified photon conversions. Depending on the signal regions, this back-  
633 ground is dominated by  $t\bar{t}$  or  $W+jets$  processes.
- 634 • **Charge flips:** Charge misidentification, i.e. events with opposite-sign isolated lep-  
635 tons where the charge of one of the leptons is misidentified because of severe bremsstrahlung  
636 in the tracker material. This background, is relevant only for electrons and is negli-  
637 gible for muons.
- 638 • **Rare SM processes:** Rare SM processes yield same-sign leptons, mostly from  $ttW$   
639 and  $ttZ$ . We also include the contribution from the SM Higgs boson produced in as-  
640 sociation with a vector boson or a pair of top quarks in this category of background.  
641 With the exception of  $ttZ$  and  $ttW$ , rares are estimated from simulation.

### 642 8.1 Fake leptons

643 The fake lepton prediction is determined as described in the SUSY AN [1] using the SUSY  
644 same-sign analysis baseline selection. The description is also included below, for convenience.

645 The uncertainties related to this estimate are discussed in Section 9.3. Note that in the case  
646 of 0 events in the application region as an input to the extrapolation into the signal regions,  
647 the prompt-nonprompt events in simulation are multiplied by the data fake rate to obtain a  
648 prediction with associated error. This error is then set as the statistical upper limit on the data-  
649 driven prediction of 0.

#### 650 8.1.1 The fake rate method

651 Background from fake leptons is estimated with the “fake rate” method. The number of events  
652 in the sample with at least one lepton that passes a loose selection but fails the full set of tight  
653 identification and isolation requirements (application region) is weighted using the “tight-to-  
654 loose” ratio, i.e. the probability that a loosely identified non-prompt lepton also passes the  
655 full set of requirements. This probability is measured as a function of lepton  $p_T$  and  $\eta$ , as  
656 well as event kinematics, in a control sample of QCD multijet events that are enriched in non-  
657 prompt leptons (measurement region). Such region is triggered by the auxiliary triggers de-  
658 fined in Table 4 and requires only one denominator lepton in the event, one recoiling jet with  
659  $\Delta R(jet, lep) > 1.0$  and low MET and MT to suppress the contribution from  $W$  and  $Z$ . The main  
660 systematic effects are the non-universality of the “tight-to-loose” ratio, particularly due to the  
661 dependency from the mother parton  $p_T$  and the flavor composition of the sample, and the  
662 prompt contribution in the measurement region.

663 In the SUSY same-sign 13 TeV analyses, developments were deployed [9][17] in the fake rate  
664 estimation in order to reduce the dependency to the mother parton  $p_T$  by using a new proxy of  
665 the lepton  $p_T$ :

$$\begin{aligned}
 & \text{if } p_T^{rel} > I_3 : \quad p_T \rightarrow p_T \cdot (1 + \max(0, I_m - I_1)) \\
 & \text{else : } \quad p_T \rightarrow \max(p_T, p_T(jet) \cdot I_2)
 \end{aligned} \tag{5}$$

666 In addition to the corrected lepton  $p_T$  definition, another development has been introduced  
667 since [17]: for the electrons, the flavor dependency of the tight-to-loose ratio is reduced by

668 extrapolating on both isolation and lepton MVA ID. These improvement are included in the  
 669 2016 and in the current (2016+2017)  $t\bar{t}t\bar{t}$  analyses.

### 670 8.1.2 Fake Rate measurement

671 We derive different versions of the fake rate for muons and electrons, collected with triggers  
 672 with an isolation requirement. The numerator and denominator selections are defined as in  
 673 section 4.5.

674 Events for fake rate measurement need to pass the following requirements:

- 675 • pass a specific auxiliary trigger, described in Table 4
- 676 • only one denominator lepton (FO)
- 677 • at least one jet with  $\Delta R(\text{jet}, \text{FO}) > 1$
- 678 •  $E_T^{\text{miss}} < 20 \text{ GeV}$ ,  $M_T < 20 \text{ GeV}$ .

679 The requirements on  $E_T^{\text{miss}}$  and  $M_T$  are intended to suppress the contribution from prompt  
 680 leptons in the measurement region. Such contribution is subtracted from the fake rate using  
 681 DY, WJets, and  $t\bar{t}$  Monte Carlo samples. For the 2016 data, these samples were normalized in  
 682 the control region defined with  $E_T^{\text{miss}} > 20 \text{ GeV}$ ,  $70 < M_T < 120 \text{ GeV}$  and a tight lepton. For  
 683 the 2017 data, due to the increased PU contribution, an improved normalization technique is  
 684 used, based on template fits of the full  $M_T$  distribution. The  $M_T$  distribution in data with  
 685  $E_T^{\text{miss}} > 30 \text{ GeV}$  and lepton  $p_T > 20 \text{ GeV}$  is fitted with the sum of two templates derived  
 686 from MC, one for QCD and one for the sum of the electroweak processes (DY, WJets and  $t\bar{t}$ ).  
 687 The normalization of the electroweak processes is extracted from the fit and half the difference  
 688 between the normalization and unity is taken as an uncertainty. As the shape of the non-prompt  
 689 component may not be well-modeled, we repeat the fit replacing the QCD MC template with  
 690 a data-driven template extracted from events failing the isolation cut. The difference between  
 691 the two fits is taken as an additional uncertainty in the normalization, added in quadrature  
 692 with the uncertainty obtained from the fit above.

693 The nominal and alternative fits for electrons and muons are shown in Figure 25 (2017) and  
 694 Figure 26 (2018), and the resulting normalization corrections for the electroweak samples are  
 695 shown in Table 12. Statistical uncertainties on the measurement are assumed to be negligible.  
 696 Results with  $35.9 \text{ fb}^{-1}$  of 2016 data are shown in [1], while the 2017 and 2018 results with  
 697  $41.5 \text{ fb}^{-1}$  and  $59.6 \text{ fb}^{-1}$  of data are shown in Figs. 27-29, 28-30 below.

	template	isolated		non-isolated	
		$e$	$\mu$	$e$	$\mu$
2017	MC	1.215	1.222	1.208	1.202
	data	1.277	1.195	1.298	1.178
2018	MC	1.200	1.283	1.202	1.288
	data	1.268	1.252	1.297	1.250

Table 12: Normalization scale factors for electroweak samples derived with two different  $M_T$  templates for QCD: MC and data (the data template refers to the inverted isolation region).

### 698 8.1.3 Fake Rate closure in MC: $t\bar{t}$ and W+jets: 2017/2018 samples

699 Using these definitions, we tested the closure of the method in the baseline and signal regions  
 700 and, inclusively in lepton  $p_T$ , for the most relevant kinematic distributions in 31 and 32. The  
 701 level of closure obtained is typically at 30% or better, similar to the 2016 one. The closure in the  
 702 electron channel showed a potential trend at high  $p_T$ , with deviations up to 60%.

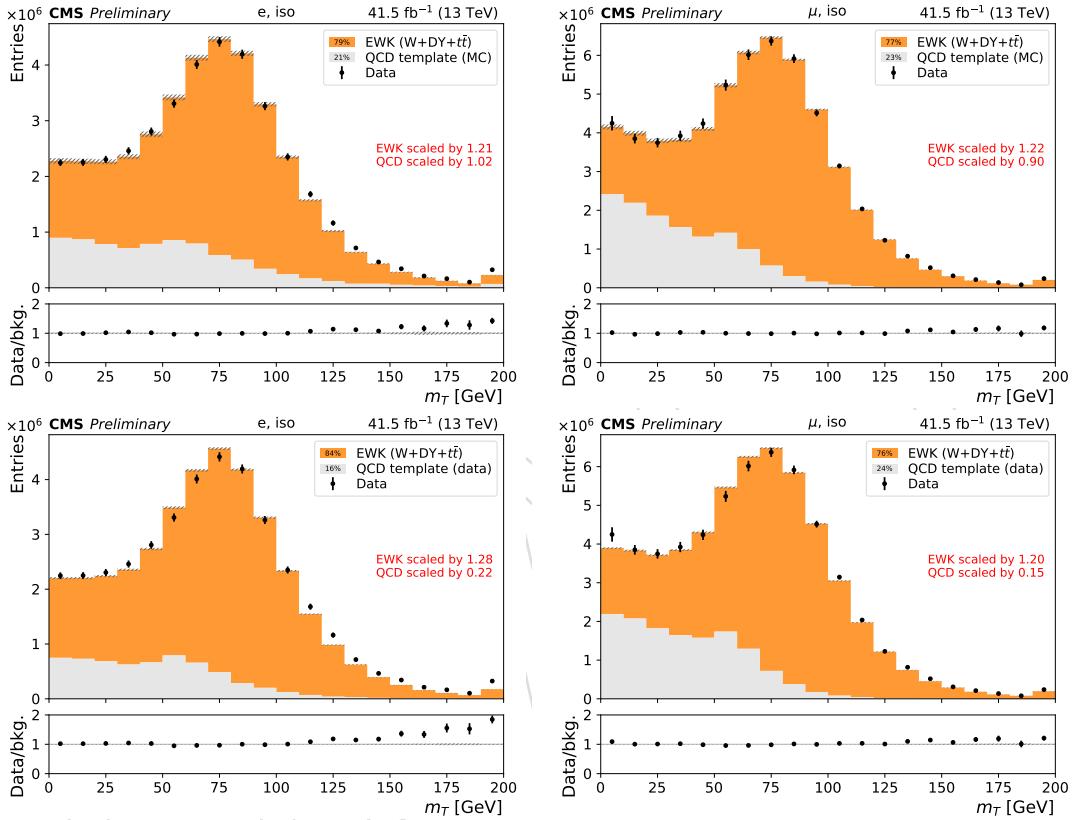


Figure 25: Isolated triggers, 2017 Data and MC: Fitted  $M_T$  distribution used to derive the normalization of electroweak samples (DY, WJets,  $t\bar{t}$ ) in the fake rate measurement region. Electrons are shown on the left, muons on the right. From top to bottom, the results from the nominal selection ( $E_T^{\text{miss}} > 30 \text{ GeV}$  and lepton  $p_T > 20 \text{ GeV}$ ) with the QCD MC template and alternative data QCD template are shown.

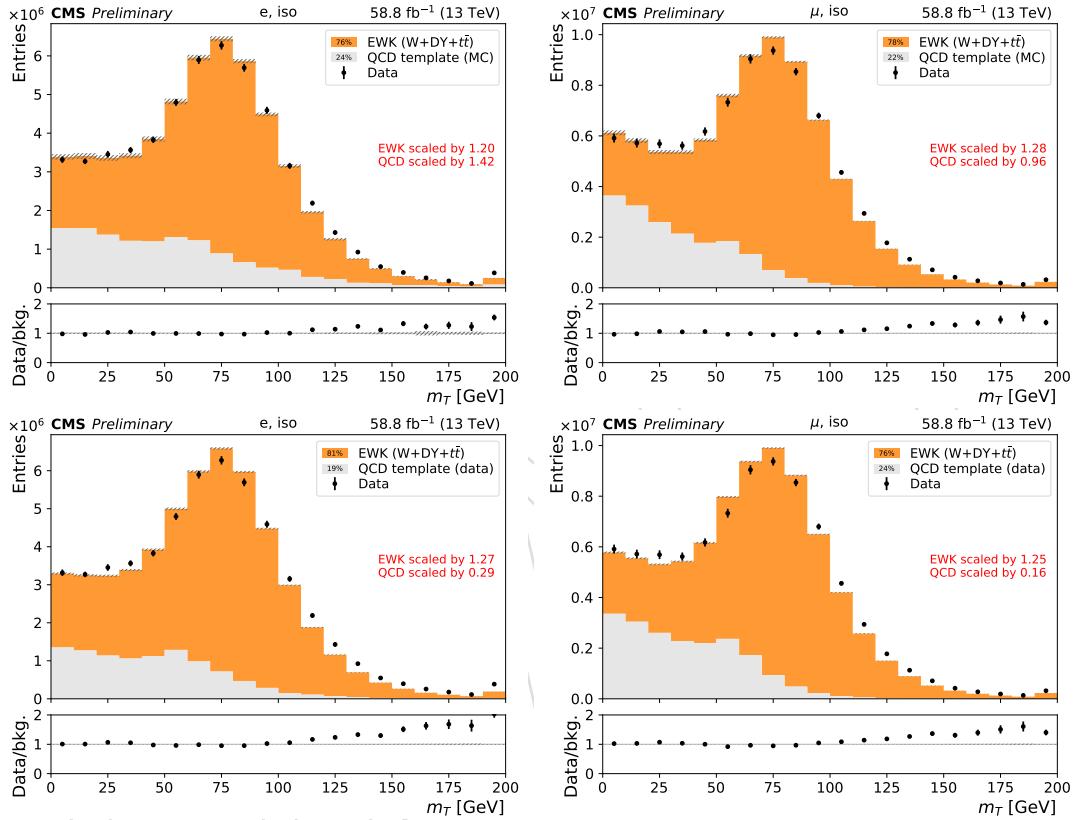


Figure 26: Isolated triggers, 2018 Data and MC : Fitted  $M_T$  distribution used to derive the normalization of electroweak samples (DY, WJets,  $t\bar{t}$ ) in the fake rate measurement region. Electrons are shown on the left, muons on the right. From top to bottom, the results from the nominal selection ( $E_T^{\text{miss}} > 30$  GeV and lepton  $p_T > 20$  GeV) with the QCD MC template and alternative data QCD template are shown.

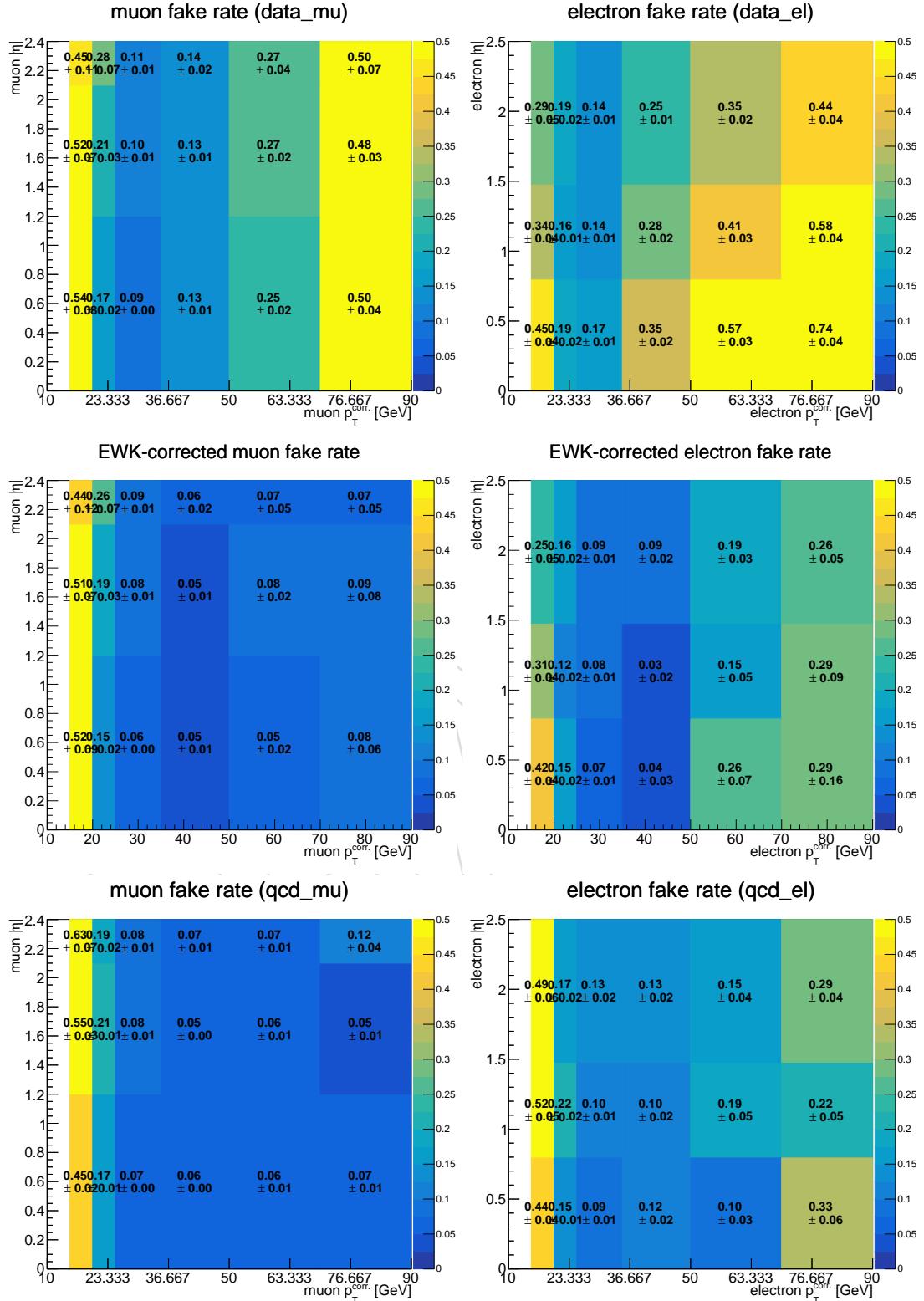


Figure 27: Isolated triggers: Fake rate for muons (left) and electrons (right) for: 2017 data uncorrected (top), 2017 data corrected for EWK contribution (middle) and 2017 QCD MC (bottom). Uncertainties are only statistical.

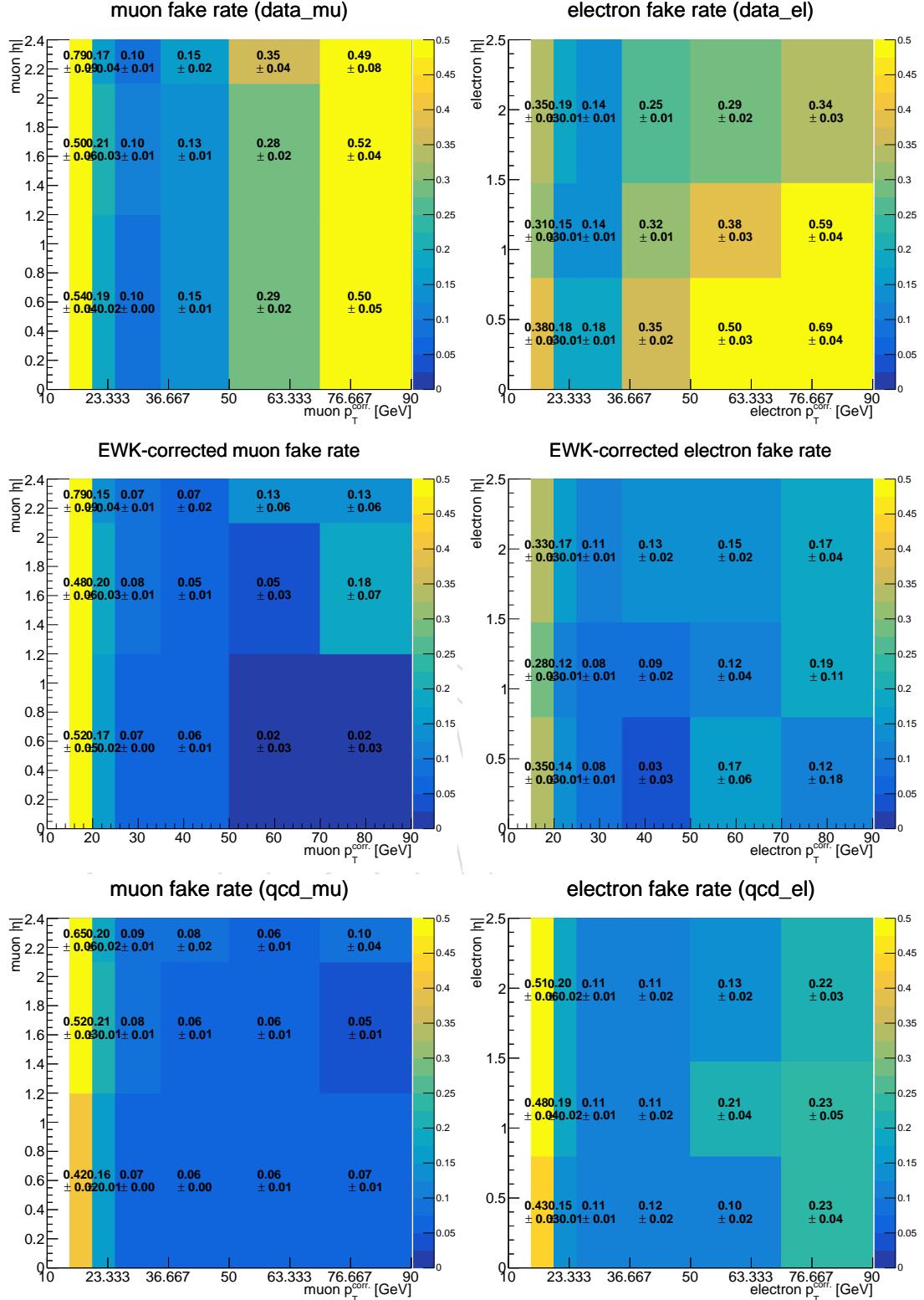


Figure 28: Isolated triggers: Fake rate for muons (left) and electrons (right) for: 2018 data uncorrected (top), 2018 data corrected for EWK contribution (middle) and [2017](#) QCD MC (bottom). Uncertainties are only statistical.

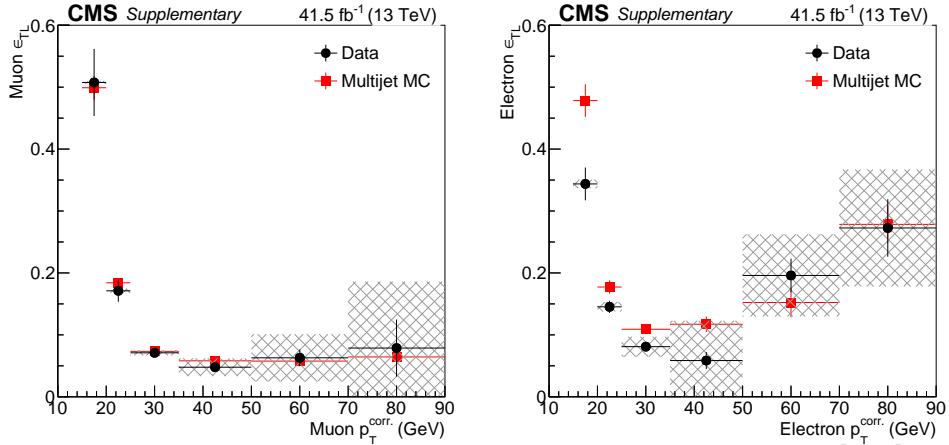


Figure 29: EWK-corrected data fake rate projected vs  $p_T$  for 2017 data (black) and 2017 QCD MC (red), for muons (left) and electrons (right). The shaded band in the projection is the systematic uncertainty related to the EWK contamination.

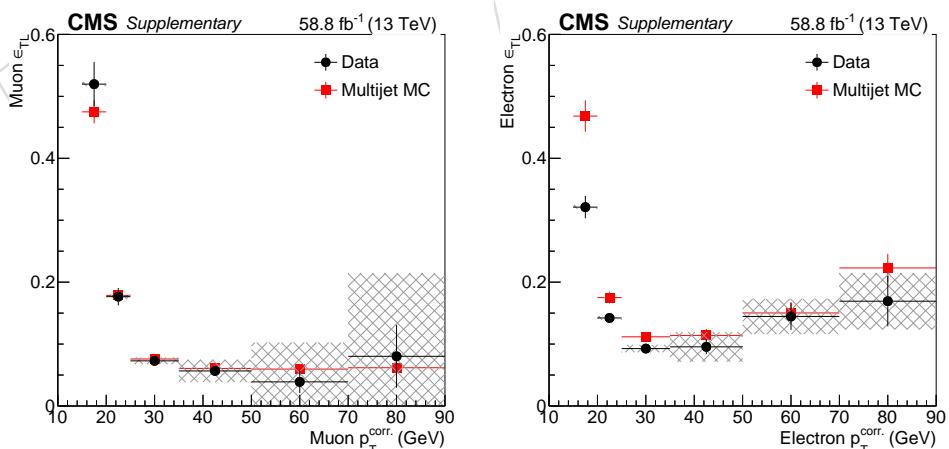


Figure 30: EWK-corrected data fake rate projected vs  $p_T$  for 2018 data (black) and 2017 QCD MC (red), for muons (left) and electrons (right). The shaded band in the projection is the systematic uncertainty related to the EWK contamination.

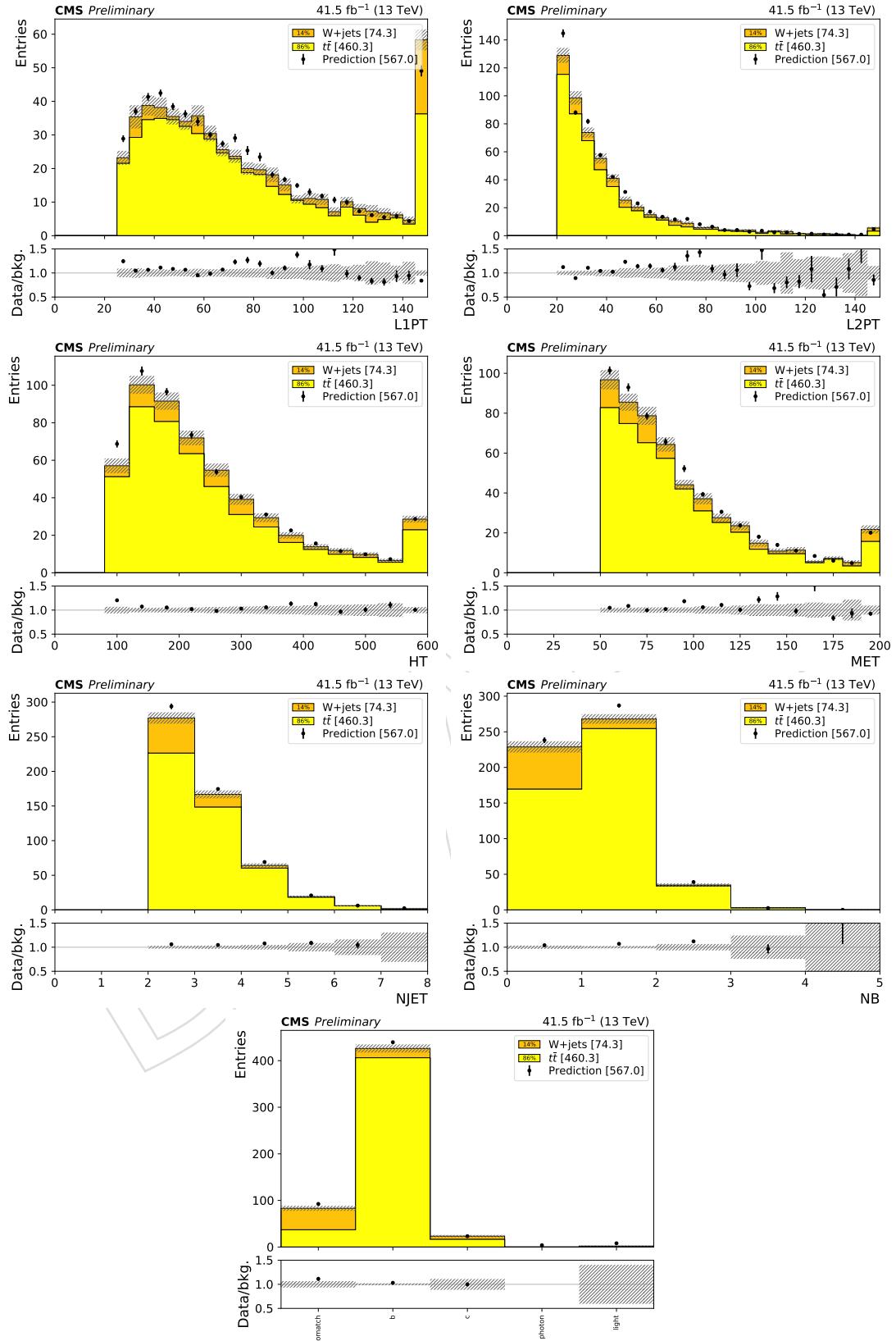


Figure 31: 2017 MC: Electron+muon fake rate closure for QCD measurement in MC soup.

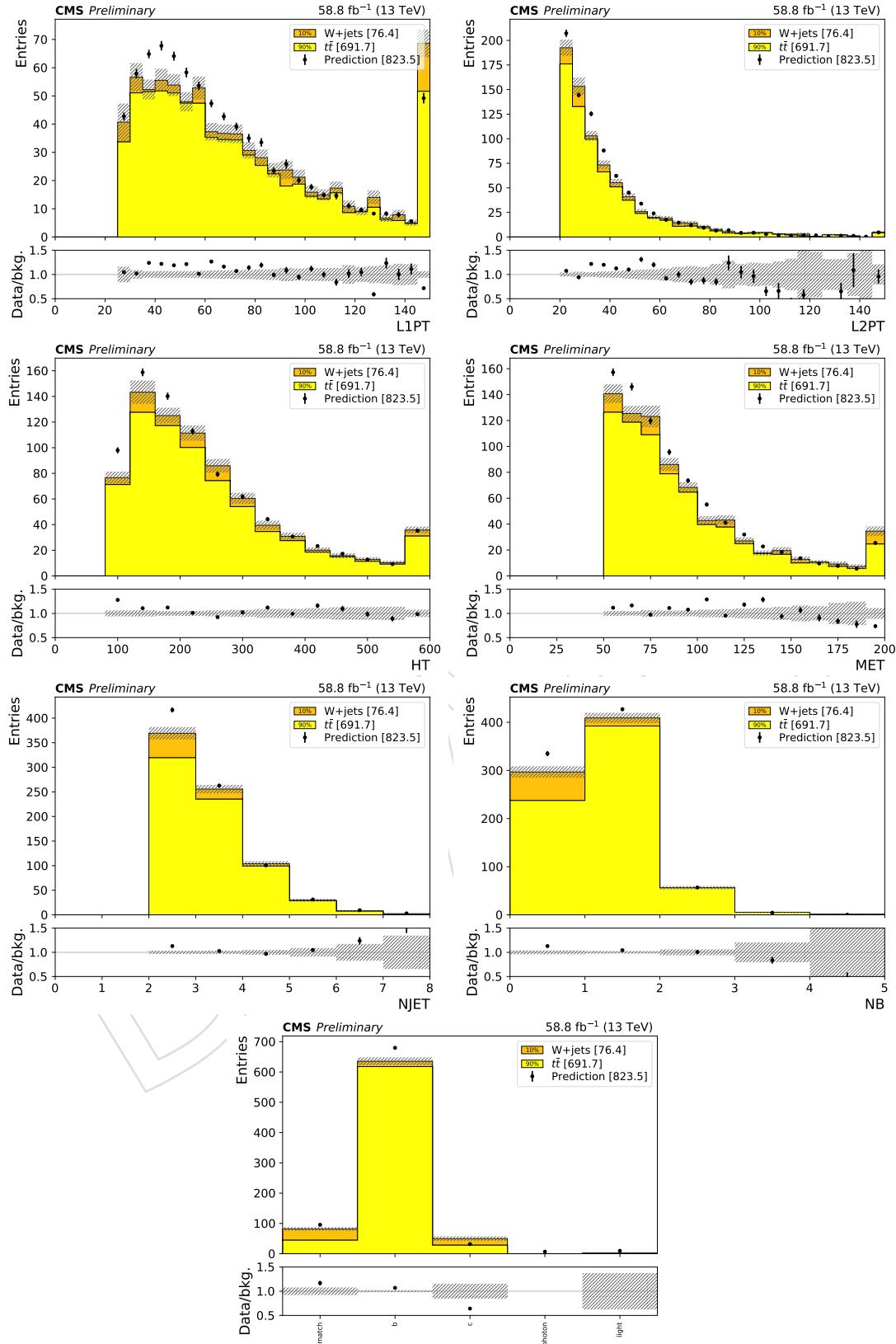


Figure 32: 2018 MC: Electron+muon fake rate closure for QCD measurement in MC soup. **QCD**  
**Fake rate to be updated. 2018 bcToE QCD MC is not available yet**

703 **8.2 Charge misidentification**

704 The charge flip prediction is determined as described in the SUSY AN [1] using the SUSY same-  
 705 sign analysis baseline selection. The description is also included below, for convenience. The  
 706 uncertainties related to this estimate are discussed in Section 9.3.

707 The background due to charge flips is estimated by selecting opposite-sign ee or e $\mu$  events passing  
 708 the full kinematic selection and then weighting them by the  $p_T$  and  $\eta$ -dependent probability  
 709 of electron charge mismeasurement

710 This probability, shown in Fig. 33, is obtained from a soup of TTbar and DY simulation and is  
 711 then validated with a control data sample of same-sign Z $\rightarrow$ ee events, using a  $E_T^{\text{miss}} < 50 \text{ GeV}$   
 712 requirement to be orthogonal to the signal region. The level of agreement in this control region  
 713 is used to gauge the associated systematic uncertainty and to derive a correction to the MC-  
 714 based rate estimation. In the 2016 data, we find good agreement between prediction and data  
 715 in the control region [1]. In the 2017 and 2018 data, the MC Flip Rate is significantly lower  
 716 than the 2016 one due to the upgraded pixel detector. However, the prediction in the same-  
 717 sign Z $\rightarrow$ ee region is found to only be about 30% lower than the observed number of events in  
 718 this region, as shown in Figure 35. Consequently, the 2017 and 2018 charge misidentification  
 719 prediction are scaled by nearly 50%, as seen in Figure 34. The yearly scaling is given in Table 13.  
 720 Since we do not find significant trends in the lepton kinematics, so we do not consider  $p_T$  and  
 721  $\eta$ -dependent corrections. In addition to the statistical uncertainties, we apply a 20% systematic  
 722 uncertainty on this background prediction for all years. In MC, the flip rate for muons is found  
 723 to be  $O(10^{-6})$  and is therefore neglected.

year	obs/pred
2016	1.01
2017	1.44
2018	1.41

Table 13: Ratio of observed flip rate in data to the flip rate in simulation. These are the multiplicative correction to the MC-based charge flip probabilities.

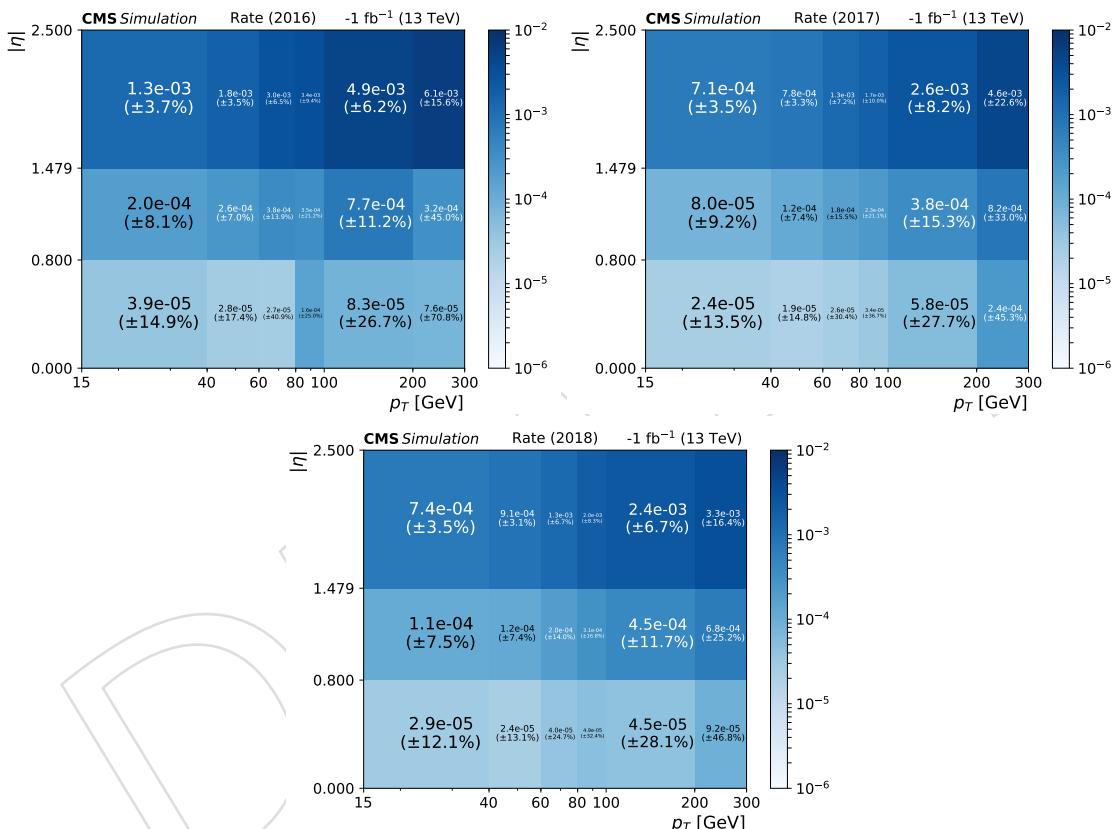


Figure 33: Electron charge flip rate for 2016, 2017, and 2018.

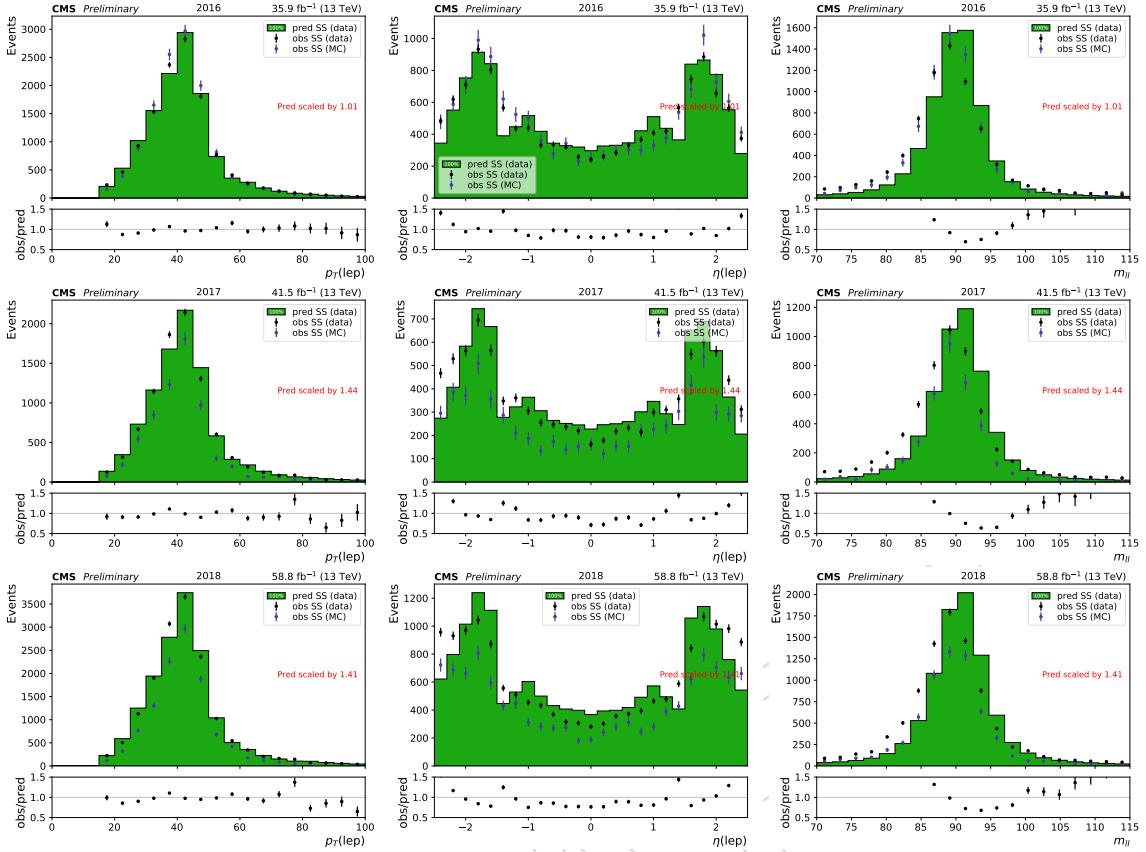


Figure 34: Predicted and observed lepton  $p_T$  (left) and  $\eta$  (middle) and  $m_{\ell\ell}$  (right) in a same-sign  $Z \rightarrow ee$  peak for years 2016, 2017, and 2018 from top to bottom. The prediction is normalized to the observed data.

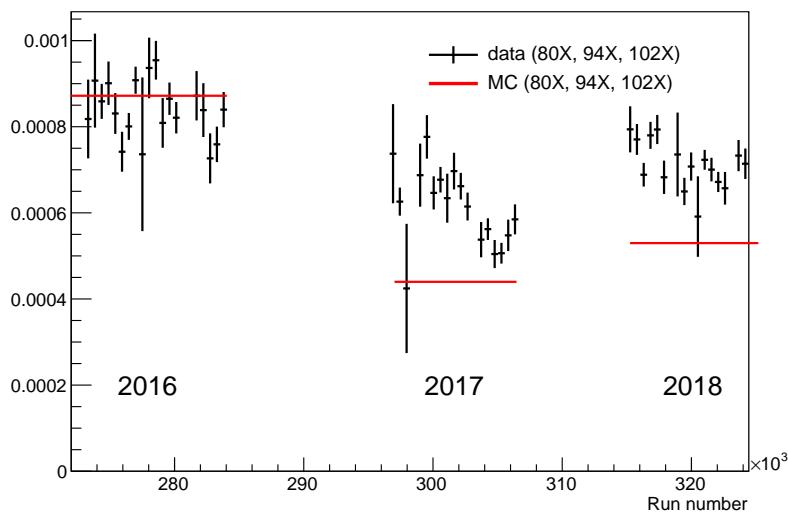


Figure 35: Electron charge flip rate in simulation and data as a function of time

### 724 8.3 Rare SM processes

725 Rare SM processes that result in the production of same-sign leptons are estimated using MC  
 726 samples. The processes with the largest contributions, ttW, ttZ and ttH, are treated individually  
 727 and assigned separate uncertainties. Processes with smaller contributions, including diboson  
 728 (WZ, ZZ) and triboson (WWW, WWZ, WZZ, ZZZ), Higgs (HZZ, VH), same-sign WW from  
 729 both single (qqWW) and double-parton scattering (DPS WW), rare top channels (tZq), and  
 730 ttt+W,tt+j are grouped in a single category termed "Rare SM" in all plots. Similarly, processes  
 731 where one of the leptons is an electron from an unidentified photon conversion are grouped  
 732 in a category called "X+ $\gamma$ " in all plots: these include  $W\gamma$ ,  $Z\gamma$ ,  $t\bar{t}\gamma$  and  $t\gamma$ . Both the "Rare SM"  
 733 and "X+ $\gamma$ " categories are assigned a large ( $\pm 50\%$ ) theoretical normalization uncertainty. The  
 734 breakdown of the individual processes of the "Rare SM" category is shown in Figure 36. The  
 735 theory uncertainties considered are described in more detail in Sec. 8.4.

736 In addition to the theoretical uncertainties, all samples are assigned uncertainties based on  
 737 reconstruction, as summarized in Section 9, due to JES, b-tagging, lepton and trigger scale  
 738 factors uncertainties and luminosity. Finally, the statistical uncertainty of the MC samples is  
 739 also taken into account.

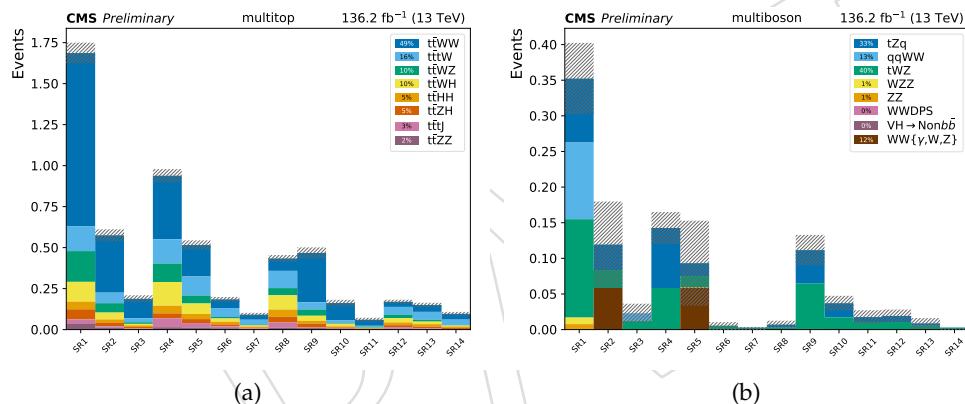


Figure 36: Relative composition of multi-top (left) and multi-boson (right) rare backgrounds in the signal regions for all MC.

## 740 8.4 Theoretical uncertainties on TTW, TTZ and TTH

741 For the largest SM backgrounds, we carefully assign uncertainties on the overall normalization  
 742 and on the shape across signal regions.

### 743 8.4.1 Normalization

744 For  $t\bar{t}W$  and  $t\bar{t}Z$ , a 40% normalization uncertainty is applied, but the results are not sensitive to  
 745 this initial choice, because the respective control regions included in the maximum likelihood  
 746 fit, discussed in more detail in Sec. 11, are used to further constrain the normalization of these  
 747 backgrounds.

748 For  $t\bar{t}H$ , the normalization uncertainty used in the 2016 analysis was 50%, to cover the 1.5  
 749 signal strength observed in the results of HIG-17-004. For the full Run2 analysis, we rely on the  
 750 updated measurement of HIG-17-035, which finds a signal strength of  $1.26^{+0.31}_{-0.26}$ . We therefore  
 751 apply a 25% normalization uncertainty.

### 752 8.4.2 Shape

753 To evaluate the theory uncertainties on the shape, we explore scale and PDF variations, as well  
 754 as differences between LO and NLO samples, where both are available. These studies were  
 755 performed with 2016 MC.

756 The PDF shape uncertainties for these samples are generally smaller than the MC statistical  
 757 uncertainties on the background yield in each region, so no additional shape uncertainty is  
 758 assigned to account for them.

759 The NLO scale uncertainties for these samples are fairly stable at around 15% for  $t\bar{t}W$ ,  $t\bar{t}Z$ , and  
 760  $t\bar{t}H$ , as shown in Figures 38, 37, 39. With the exception of  $t\bar{t}H$ , these figures also show the LO  
 761 scale uncertainty, which covers well the difference between LO and NLO. The LO agreement  
 762 with the NLO shapes, within LO scale uncertainty, gives confidence in the NLO scale uncer-  
 763 tainties as they are evaluated using the same variations of renormalization and factorization  
 764 scales.

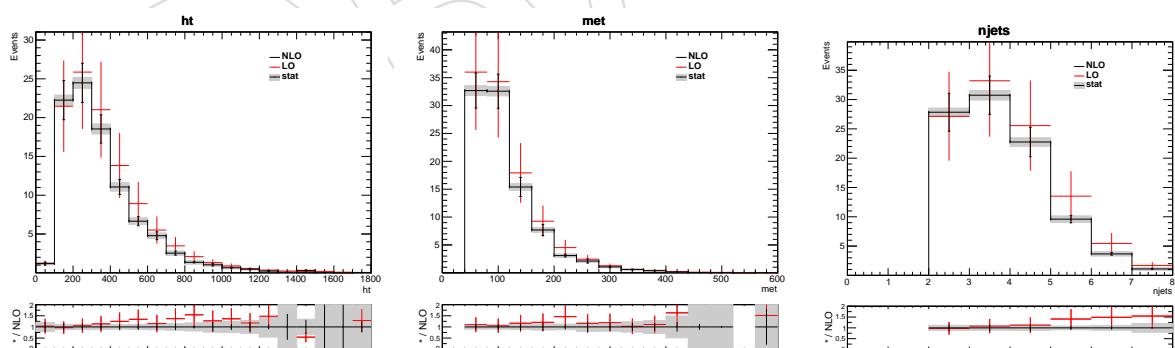


Figure 37: Comparison of LO and NLO kinematics for  $ttZ$  with scale uncertainties in the base-  
 line region and statistical uncertainties on NLO. The error bars on the NLO (black) are the NLO  
 uncertainties, while the error bars on the LO (red) are the LO uncertainties. The gray band is  
 the stat uncertainty on the NLO

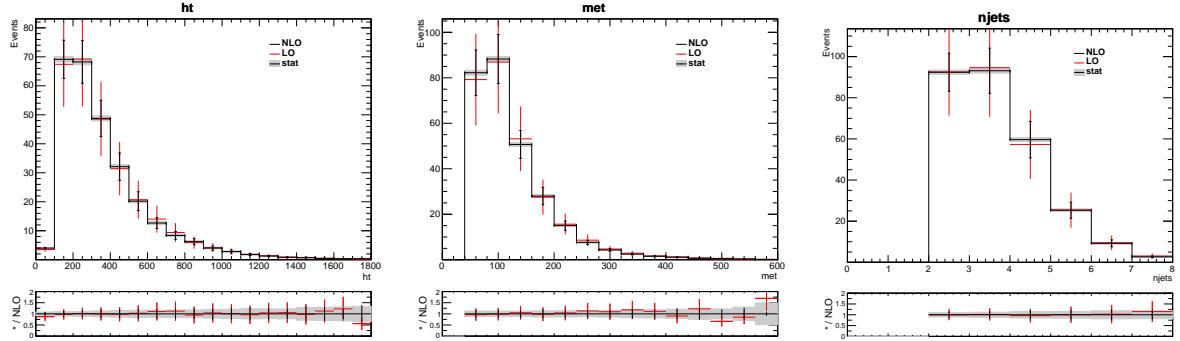


Figure 38: Comparison of LO and NLO kinematics for  $t\bar{t}W$  with scale uncertainties in the baseline region and statistical uncertainties on NLO. The error bars on the NLO (black) are the NLO uncertainties, while the error bars on the LO (red) are the LO uncertainties. The gray band is the stat uncertainty on the NLO

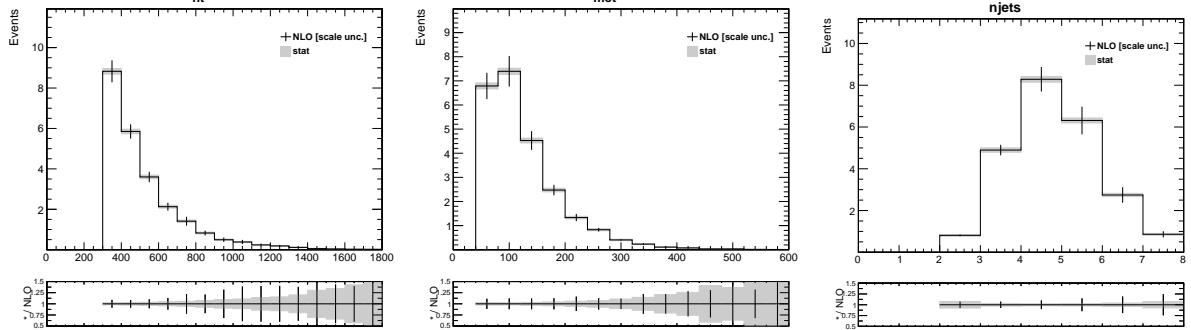


Figure 39:  $t\bar{t}H$  NLO with NLO scale uncertainties in the baseline region. The gray band is the stat uncertainty on the NLO

## 765 9 Systematic Uncertainties

766 With respect to the 2016 SUSY same-sign analysis [1], the 2016 tt $t\bar{t}$  analysis only changed uncer-  
 767 tainties on  $t\bar{t}Z$ ,  $t\bar{t}H$ , and  $t\bar{t}W$ , including also uncertainty for  $t\bar{t}W+bb$ ,  $t\bar{t}Z+bb$ , JER, and adding  
 768 ISR/FSR variations from dedicated samples. With respect to the 2016 tt $t\bar{t}$  analysis [4], the cross-  
 769 section for  $t\bar{t}H$  has a 25% uncertainty taken as a systematic. Other uncertainties, which have  
 770 changed as a result of new measurements, are marked in blue if the change is also applied  
 771 to 2016 data. All uncertainties marked in black have remained unchanged since 2016. These  
 772 changes are reflected in Table 14.

### 773 9.1 Correlation model for 2016+2017+2018

774 In order to combine the 2016+2017+2018 datasets, we have to make assumptions about the  
 775 correlation of uncertainties in the different years. Since the analysis is statistically dominated,  
 776 these correlations are expected to have a very small influence. The current model we have  
 777 chosen makes the following assumptions,

- 778 • Statistical uncertainties (FO CR stat, OS CR stat, MC stat): uncorrelated
- 779 • Luminosity, JEC, JER, b-tag, PU, lepton eff., HLT, eff.: uncorrelated (derived from

source	magnitude	effect on yield
luminosity	2.5, 2.3, 2.5%	2.5, 2.3, 2.5% (2016-2018)
JES	1 – 8%	1 – 15%
JER	1 – 10%	1 – 10%
$t\bar{t}W$ , $t\bar{t}Z$ ISR/FSR	< 15%	1 – 8%
$t\bar{t}W/Z+bb$	< 30%	< 15%
b-tag efficiency	$\sim 1$ – 20%	1 – 15%
pileup	5%	0-5%
lepton efficiencies	2 – 5% (leg)	2 – 10%
HLT efficiencies	1 – 5% (leg)	2 – 7%
FO CR stat.	1 – 100%	1 – 100% (fake bkg. only)
FR extrapolation	30-60%	30-60% (fake bkg. only)
EWK subtraction in FR	25-50% (FR)	1 – 30% (fake bkg. only)
OS CR stat.	4 – 100%	4 – 100% (charge misId. only)
charge misId.	20%	20% (charge misId. only)
$t\bar{t}W$ norm.	40%	40 % ( $t\bar{t}W$ only)
$t\bar{t}Z$ norm.	40%	40 % ( $t\bar{t}Z$ only)
$t\bar{t}H$ norm.	25%	25% ( $t\bar{t}H$ only)
MonteCarlo stat.	1 – 25%	1 – 25%
QCD scales and PDFs	$\times 0.5 / \times 2$	10 – 20% ( $t\bar{t}W, t\bar{t}Z, t\bar{t}H$ )
other bkgs.	50%	50% (Rare and $X + \gamma$ )

Table 14: Summary of the sources of uncertainties, their magnitude and their effects. The second column indicates the magnitude of the yield variation. Reported uncertainties are representative for the most relevant signal regions.

- 780 different datasets)
- 781 • ISR/FSR: uncorrelated (derived from different datasets and with respect to different  
782 MC samples)
- 783 • ttW/Z+bb, ttH normalization: correlated (based on theory predictions and CMS  
784 measurements)
- 785 • ttW and ttZ normalizations: correlated (based on theory predictions and consistent  
786 Control Regions)
- 787 • Rare and X+ $\gamma$  normalizations: correlated (based on theory predictions)
- 788 • QCD scale and PDF uncertainties: correlated (based on theory predictions)
- 789 • EWK subtraction in FR: uncorrelated (based on differently prescaled triggers in data)
- 790 • FR extrapolation: correlated (method is the same, closure in MC shows the same  
791 features)
- 792 • Charge misid.: uncorrelated (agreement in validation region is quite different, prob-  
793 ably due to different Pixel detector conditions)

794 We ran limits under three correlation assumptions with 18 cut-based bins. With respect to  
795 the nominal correlation model, fully uncorrelating (correlating) nuisances between the 3 years  
796 decreases (increases) the expected significance by 2%.

## 797 9.2 Experimental sources of uncertainties

798 One of the main experimental sources of uncertainty is the knowledge of the jet energy scale  
799 (JES), affecting all the simulated backgrounds and considered signals. The 13 TeV uncertainties  
800 vary the jet energy scale by 1–8%, depending of the transverse momentum and pseudorapidity  
801 of the jet. The impact of these uncertainties is assessed by shifting the jet energy correction  
802 factors for each jet up and down by  $\pm 1\sigma$  before the calculation of all kinematic quantities. The  
803 variations are correlated among the different signal regions as bin-by-bin migration is allowed.  
804 Variations are asymmetric in nature and are used as asymmetric nuisance parameters in the  
805 result interpretation. The JES uncertainties are propagated to the missing transverse energy  
806 and all jet-related variables (number of jets,  $H_T$ , number of b-jets) used in this analysis. As  
807 some of the simulation samples are statistically limited, the size of the JES uncertainties can  
808 reach high values in several regions not well populated by several processes. As usually those  
809 variations impact mostly background samples that do not contribute much to those regions,  
810 we consider large variations as they are obtained out of the box. Most populated signal regions  
811 shows yield variations of 8% when the jet energy scale is varied by one standard deviation.

812 The uncertainty related to the knowledge of the jet energy resolution (JER) can be considered  
813 as well for the simulated backgrounds and the signals. The effect of smearing jets on the accep-  
814 tance of events is assessed in the same manner as the JES uncertainty.

815 A similar approach is used for the uncertainties associated to the corrections for the b-tagging  
816 efficiencies for light and bottom flavor jets, which are parametrized as a function of  $p_T$ ,  $\eta$ .  
817 The variation of the scale factor is at maximum of the order of 1–20% per jet, and lead to an  
818 overall effect on yield included between 1 and 15% depending on the signal region and on the  
819 topology of the events included in those signal regions. If considering only highly populated  
820 signal regions to get an overview of the main effects on the background yields, the bulk of the  
821 ttW and ttH yield varies by  $\sim 8\%$  and the ttZ yield by  $\sim 6\%$ .

822 When applying b-tagging efficiency scale factors, it is possible to apply period-dependent scale  
823 factors, rather than an average scale factor for the entire 2016 dataset. This is relevant for signal

regions where several mistags from light jets are expected, since the mistag SFs for light jets are time-dependent. When checking the flavor composition of our main background  $t\bar{t}W$ , as in Figure 40, we find that none of our  $N_b$  bins are dominated by events with multiple light jets, so the period-dependent scale factors are not necessarily warranted. As expected, when we do apply them, we find negligible differences in our background predictions with respect to the nominal scale factors. For 2017, we apply period-dependent b-tagging SFs.

Events from processes such  $t\bar{t}W$  and  $t\bar{t}Z$ , which have two b quarks from the decay of the two top quarks, primarily enter regions requiring three or more b-tagged jets by either the mis-tag of a charm quark from the hadronic decay of a W boson or from the production of additional heavy flavor, primarily from an ISR or FSR gluon which splits to a  $b\bar{b}$  pair.

In 2016, events in simulation containing additional HF quarks not from top or W decays were scaled up to account for  $SF_{\sigma(t\bar{t}bb)/\sigma(t\bar{t}jj)} \sim 1.7 \pm 0.6$  as measured by TOP-16-010. In the high Nntag signal region bins where this effect is dominant, this resulted in a systematic uncertainty up to 20%. In 2017 and 2018, we apply the same correction as 2016 onto  $t\bar{t}$ ,  $t\bar{t}W$ ,  $t\bar{t}Z$ , and  $t\bar{t}H$ , as it is observed to bring agreement in the high statistics opposite-sign  $t\bar{t}$  control region.

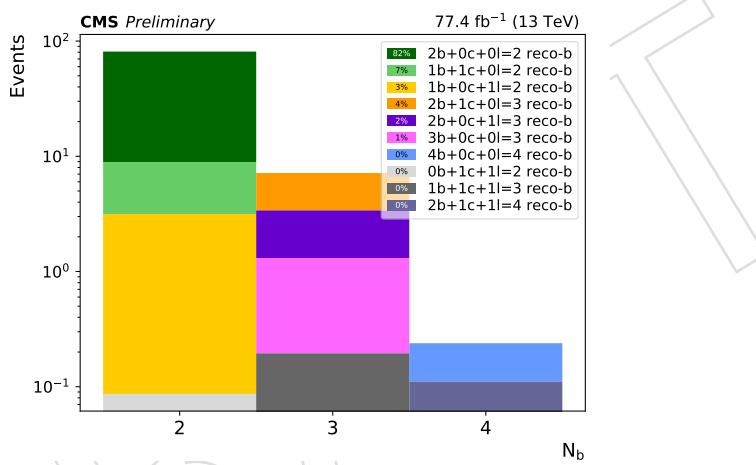


Figure 40: Flavor composition breakdown of the  $t\bar{t}W$  sample, in bins of number of b-tagged jets for all MC

- Trigger efficiency scale factors obtained from data are applied to correct the simulation.
- Lepton efficiency scale factors [13], accounting for differences between the Data and MC for reconstruction and identification of electrons and muons, are applied to all MC events. They result in uncertainties of approximately 2-3% for muons and 3-5% for electrons for 2016.
- The simulation is reweighted to match the expected data distribution in the number of collision per events; the uncertainty on the minimum bias cross-section is propagated to the final yields with an effect at the level of 4% or less across the three years.
- The  $t\bar{t}W$  and  $t\bar{t}Z$  simulation is reweighted to match the number of additional ISR/FSR jets observed in data, as described in Section 2. An uncertainty equivalent to 50% of the difference between 1 and the reweighting factor is applied to this procedure. In 2016, since the reweighing factor can be as large as 0.77, this results in an uncertainty of 15% on the highest  $N_{\text{jets}}$  signal regions. In 2017 and 2018, the reweighing factor is as large as 1.4, which results in a similar uncertainty to 2016 on the highest  $N_{\text{jets}}$  signal regions.

### 852 **9.3 Uncertainties from data-driven background estimations**

853 The following uncertainties are defined for the data-driven background estimations, based on  
 854 the statistics of the control region and the extrapolation from control to signal region.

855 For the nonprompt lepton and charge misidentified lepton backgrounds, the statistical uncer-  
 856 tainty from the control sample varies greatly depending on the signal region considered. Both  
 857 the flip rate and the fake rate are smaller than 1, so the statistical uncertainty is always smaller  
 858 (or much smaller, for charge misidentification) than the size of the predicted background. In  
 859 addition to the statistical uncertainty, the nonprompt lepton background is assigned an overall  
 860 normalization uncertainty of 30%, based on a comparison of the non-closure of two alterna-  
 861 tive methods (nominal fake-rate and in-situ fake-rate). This uncertainty is increased to 60%  
 862 for electrons with  $p_T > 50 \text{ GeV}$ , to account for the trends observed at high  $p_T$  in the electron  
 863 closure for both methods. On top of the uncertainties mentioned above, the nonprompt lepton  
 864 background prediction also includes an uncertainty related to the electroweak contamination  
 865 subtraction in the region where the fake rate is computed. To evaluate this uncertainty the fake-  
 866 rate is computed based on varying the size of the electroweak contamination, and the effect is  
 867 propagated through the whole analysis. The overall effect on the nonprompt lepton back-  
 868 ground yield lies between 1 and 50% depending of the signal region considered. The charge  
 869 misidentified lepton background is assigned a systematic uncertainty of 20%.

870 For  $t\bar{t}W$  and  $t\bar{t}Z$ , the respective control regions are used only to set the overall normalization,  
 871 while the shape is taken from simulation with the experimental uncertainties defined above  
 872 and the theoretical uncertainties described in section 8.4. The two control regions are included  
 873 in the fit with a 40% normalization nuisance parameter, and the fit constrains this parameter to  
 874 approximately 20% given the statistics of the control region, where the corresponding post-fit  
 875 constraint was 30% in the 2016 analysis.

### 876 **9.4 Systematic uncertainties from statistical sources**

877 The statistical precision given by the Monte Carlo samples has to be taken into account for  
 878 all rare SM processes. Within the framework of the HiggsCombine tool, we use the Barlow-  
 879 Beeston [21] method of handling both MC statistics and data-driven background statistics (for  
 880 the nonprompt and misidentified-charge backgrounds) via the “autoMCStats” parameter in  
 881 the datacard.

### 882 **9.5 Signal uncertainties**

883 The  $t\bar{t}t\bar{t}$  signal is assigned uncertainties based on all the effects described above. The uncer-  
 884 tainty values are summarized in Table 15. For the experimental uncertainties we report the  
 885 range of their effect across SRs, while for the theoretical uncertainties we separate the corre-  
 886 lated (acceptance\*efficiency,  $\mathcal{A}\epsilon$ ) and uncorrelated (shape) effects. The  $\mathcal{A}\epsilon$  of our SR selection,  
 887 including branching ratio, is  $1.5^{+0.02}_{-0.03}\%$ , where the uncertainties represent the QCD scale varia-  
 888 tions. The effect of QCD scale variations on the shape of  $t\bar{t}t\bar{t}$  within the set of SRs is shown in  
 889 Figure 41: is as large as 10%. PDF uncertainties on acceptance and shape are both negligible as  
 890 they are smaller than 1%

891 Uncertainties for ISR and FSR variations during sample generation for  $t\bar{t}t\bar{t}$  are assessed by  
 892 considering 4 additional samples and are shown in Figure 42. The variation in acceptance  
 893 (shape) is 1% (<8%) for ISR and 6% (<10%) for FSR.

894 For the SM  $t\bar{t}t\bar{t}$  measurement, as well as the BSM interpretations related to off-shell effects (low  
 895 mass particles, yukawa coupling,  $\hat{H}$ ), no uncertainty is assigned to the cross-section itself, since

we are measuring this process. For the BSM interpretations related to on-shell effects (heavy scalar and pseudoscalar in 2HDM and dark matter models), the SM  $t\bar{t}t\bar{t}$  process is treated as a background and assigned its SM cross section and uncertainty  $12.0^{+2.2}_{-2.5}$  fb [6].

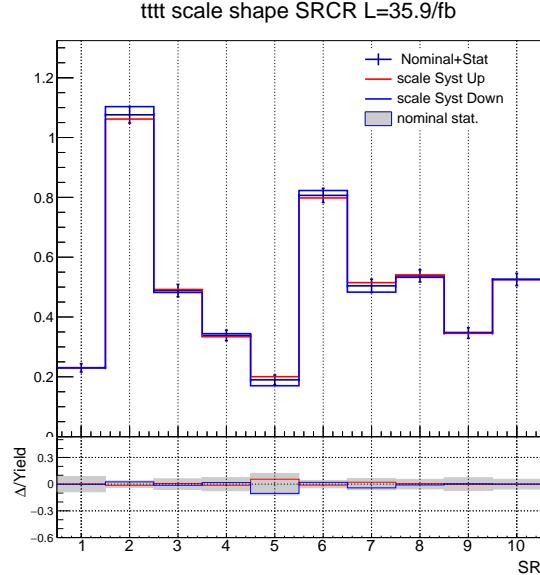


Figure 41: QCD scale uncertainties on the shape of the  $t\bar{t}t\bar{t}$  signal.

source	magnitude	effect on yield
jet ES	1 – 8%	1 – 15%
JER	1 – 10%	1 – 10%
b-tag efficiency	$\sim 1$ – 20%	1 – 15%
pileup	5%	1 – 4%
ISR	-	<8%
FSR	-	<10%
lepton efficiencies	2 – 5% (leg)	2 – 10%
HLT efficiencies	1 – 5% (leg)	2 – 7%
MonteCarlo stat.	1 – 100%	< 8%
QCD scales (acceptance)	$\times 0.5 / \times 2$	2%
QCD scales (shape)	$\times 0.5 / \times 2$	1-10%
PDFs (acceptance)	Envelope	$\leq 1\%$
PDFs (shape)		$\leq 1\%$
$\alpha_S$ (acceptance)		$\leq 6\%$
$\alpha_S$ (shape)		$\leq 1\%$

Table 15: Summary of the sources of uncertainties for  $t\bar{t}t\bar{t}$  signal.

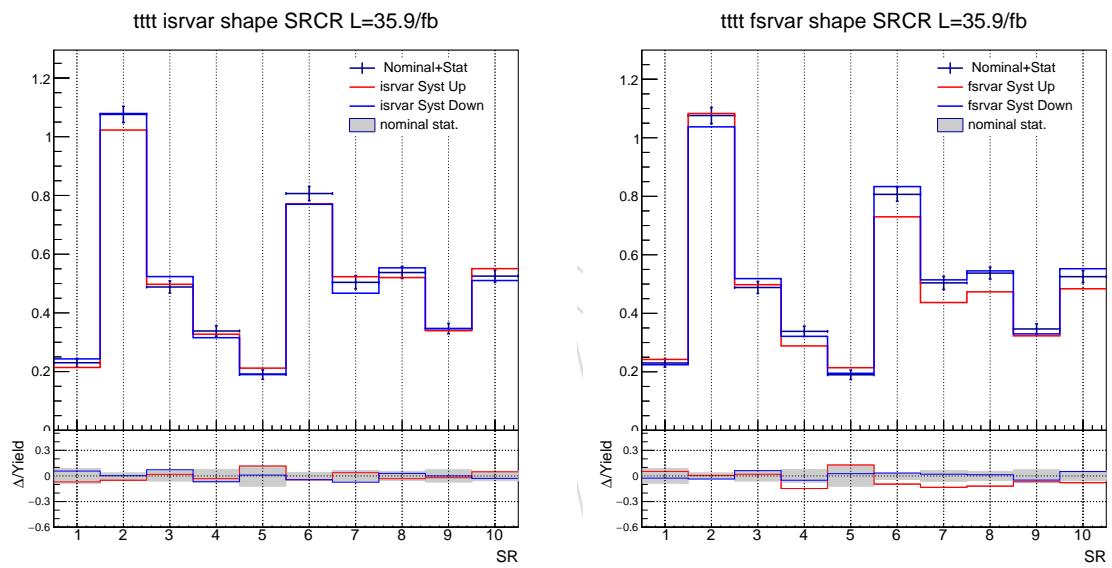


Figure 42: Effect of ISR (left) and FSR (right) variations on the  $t\bar{t}t\bar{t}$  signal.

## 899 9.6 Summary of theoretical uncertainties on MC-based processes

900 Table 16 shows a summary of the cross-section, scale, and pdf uncertainties on the MC-based  
 901 processes, including signal, indicating whether the nuisance contains an acceptance and/or  
 902 shape component. Associated with each row, there are 3 nuisances, parameterized as one log-  
 903 normal, and two shape nuisances. Some notes about individual processes follows.

- 904 • Signal  $t\bar{t}t\bar{t}$  has no log-normal nuisance associated with cross-section uncertainty be-  
 905 cause it is being measured.
- 906 • Both  $t\bar{t}W$  and  $t\bar{t}Z$  do not have an acceptance component in the scale/pdf variations  
 907 because the 40% log-normal is taken to generally cover cross-section  $\times$  acceptance as  
 908 they are constrained by dedicated control regions.
- 909 •  $t\bar{t}H$  has a 25% cross-section log-normal to reflect the result of HIG-17-035 as men-  
 910 tioned in Section 8.4.
- 911 • For X+Gamma and ttVV, the largest cross-section uncertainty on any constituent  
 912 subprocess is taken on the whole process. Thus, x+gamma obtains an 11% nor-  
 913 malization uncertainty driven by  $t\bar{t} + \gamma$  (+9.9%, -11.2%), and ttVV obtains a 11%  
 914 normalization uncertainty from ttWW (+8.6% -11.3%).
- 915 • For rares, we obtain uncertainties directly from MC samples by calculating the summed  
 916 yield of rares for nominal variation, scale up, scale down, pdf up, and pdf down.  
 917 Adding scale and pdf variations in quadrature yields a 20% variation for 2016 and  
 918 2017 samples. We proceed with a log-normal of 20% on the sum of these processes.

DRAFT

lnN		scale		pdf	
		acceptance	shape	acceptance	shape
$t\bar{t}t\bar{t}$		x	x	x	x
$t\bar{t}W$	1.40		x		x
$t\bar{t}Z$	1.40		x		x
$t\bar{t}H$	1.25	x	x	x	x
X+ $\gamma$	1.11	x	x	x	x
$t\bar{t}VV$	1.11	x	x	x	x
Rares	1.20	x	x	x	x

Table 16: Summary of the theory uncertainties

---

## 919 10 Kinematics

920 Prefit kinematic distributions of the 2017 and 2018 data events passing the baseline selection  
 921 and falling into the  $t\bar{t}W$  control region are shown in Figs. 43-45. Similarly, distributions of  
 922 events falling into the  $t\bar{t}Z$  control region are shown in Figs. 47-49. The background prediction  
 923 for the sum of signal regions is shown in Figs. 51-53.

process	SF (2016)	SF (2017)	SF (2018)	SF (Run2)
ttz	1.58	1.005	1.207	1.258
ttw	1.347	1.35	1.156	1.299
tth	1.087	1.089	1.045	1.088
ttt	1.175	0.845	1.451	1.053
fakes	1.064	1.163	1.081	1.125
xg	1.06	1.035	1.015	1.014
rares	1.055	1.017	1.023	1.017
ttvv	1.028	1.018	1.02	1.011
flips	1.016	1.007	0.999	1.001

Table 17: Postfit/prefit scale factors with the BDT analysis. These are the ratio of the normalizations before and after the fit, and they are impacted by all nuisance parameters affecting each background. For this reason, they are not necessarily expected to be consistent between years, as are the individual nuisance parameters themselves (shown in Appendix A).

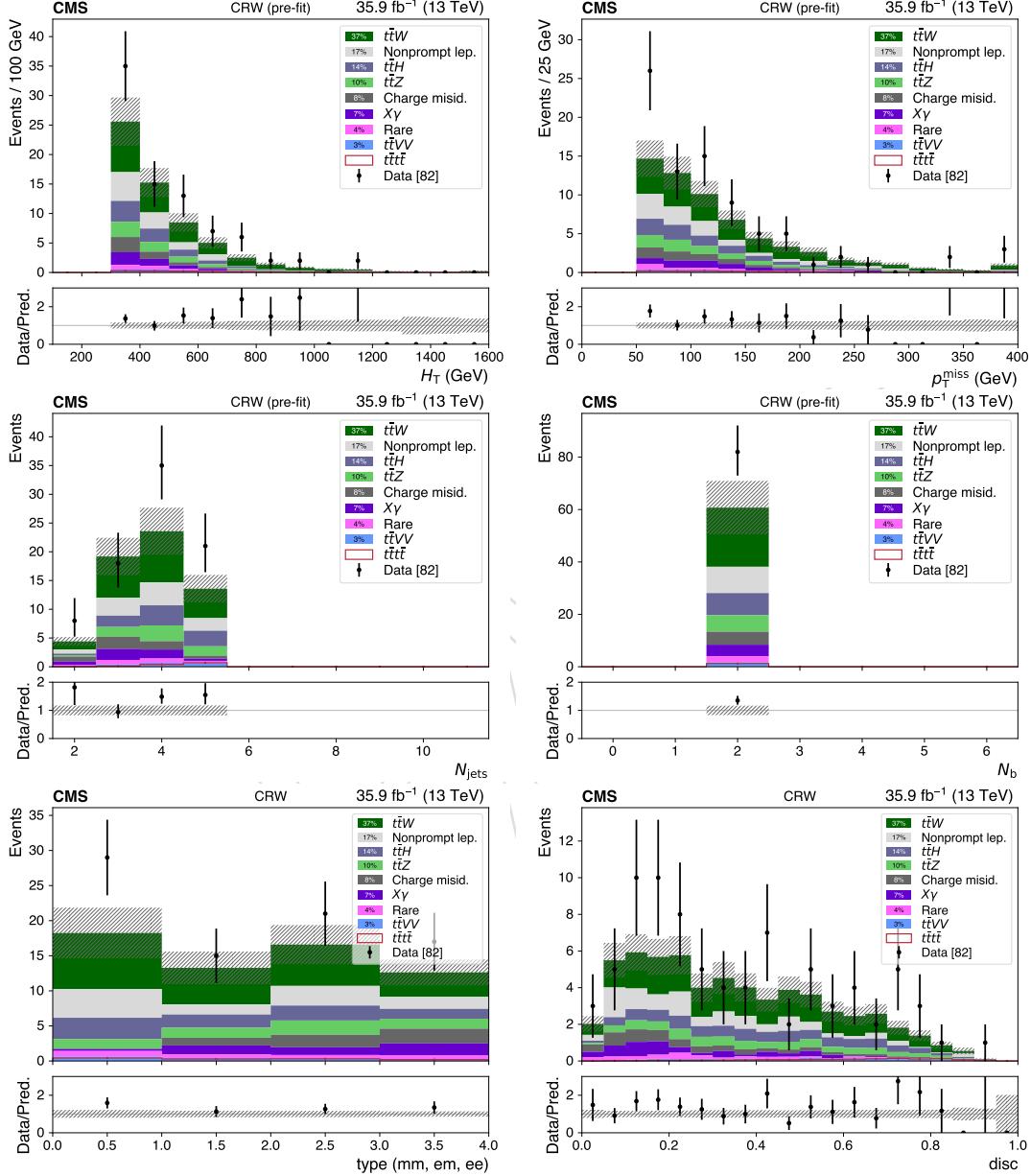


Figure 43: 2016 data and predictions: Prefit distributions of the main analysis variables in the  $t\bar{t}W$  control region:  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_b$  jets, lepton flavor, and raw BDT discriminant where the last bin includes the overflow.

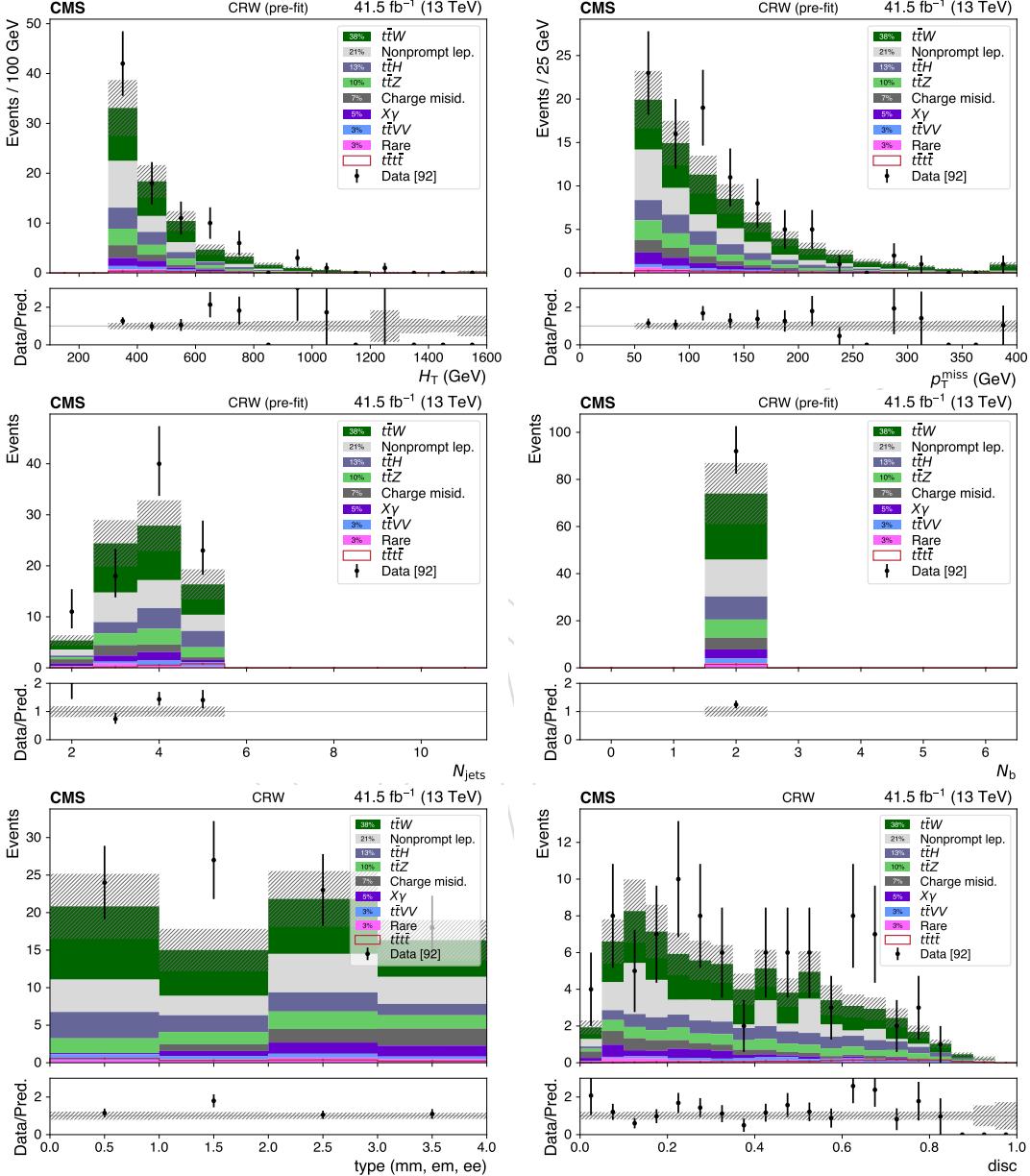


Figure 44: 2017 data and predictions: Prefit distributions of the main analysis variables in the  $t\bar{t}W$  control region:  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_b$  jets, lepton flavor, and raw BDT discriminant where the last bin includes the overflow.

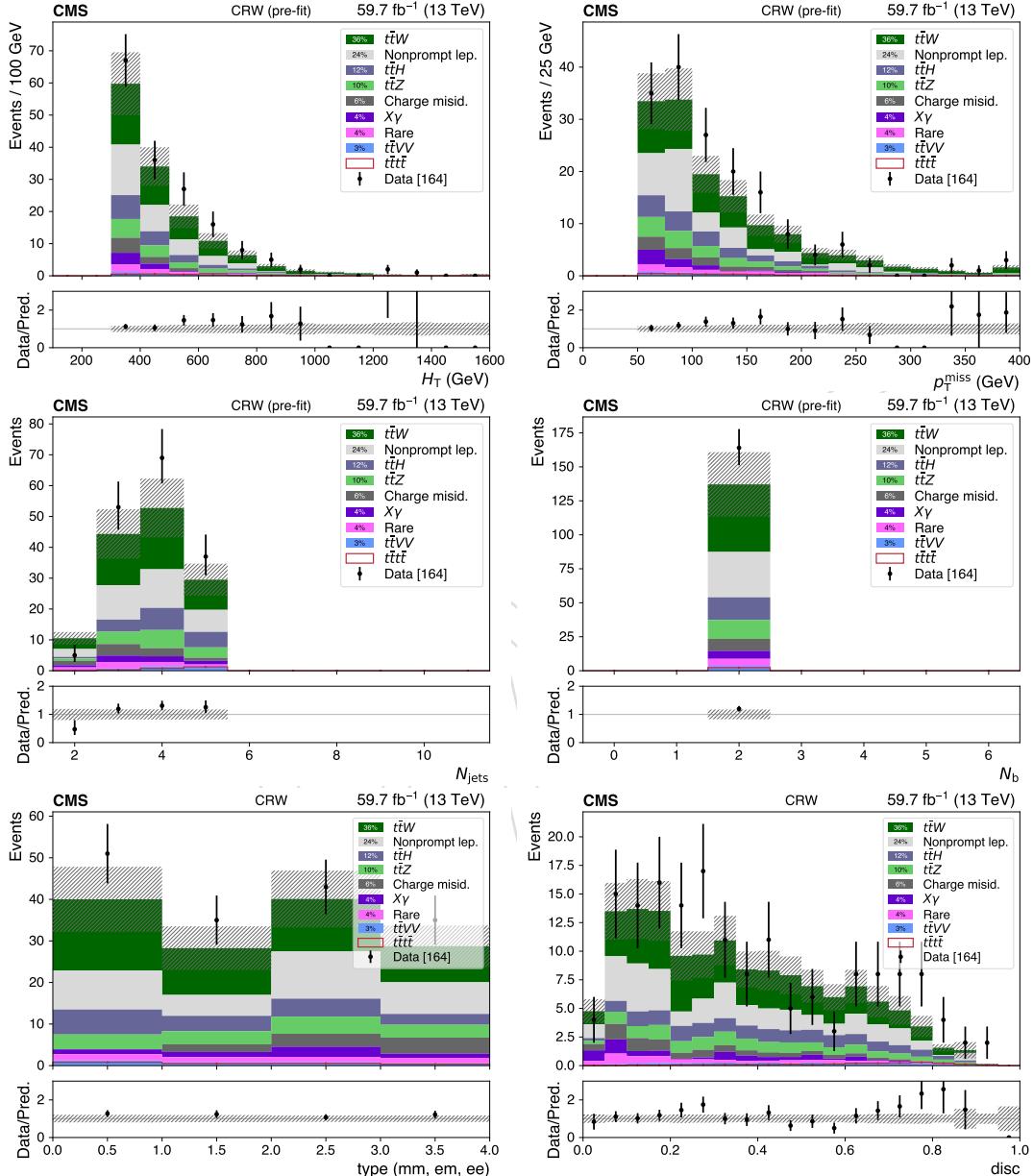


Figure 45: 2018 data and predictions: Prefit distributions of the main analysis variables in the  $t\bar{t}W$  control region:  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_b$  jets, lepton flavor, and raw BDT discriminant where the last bin includes the overflow.

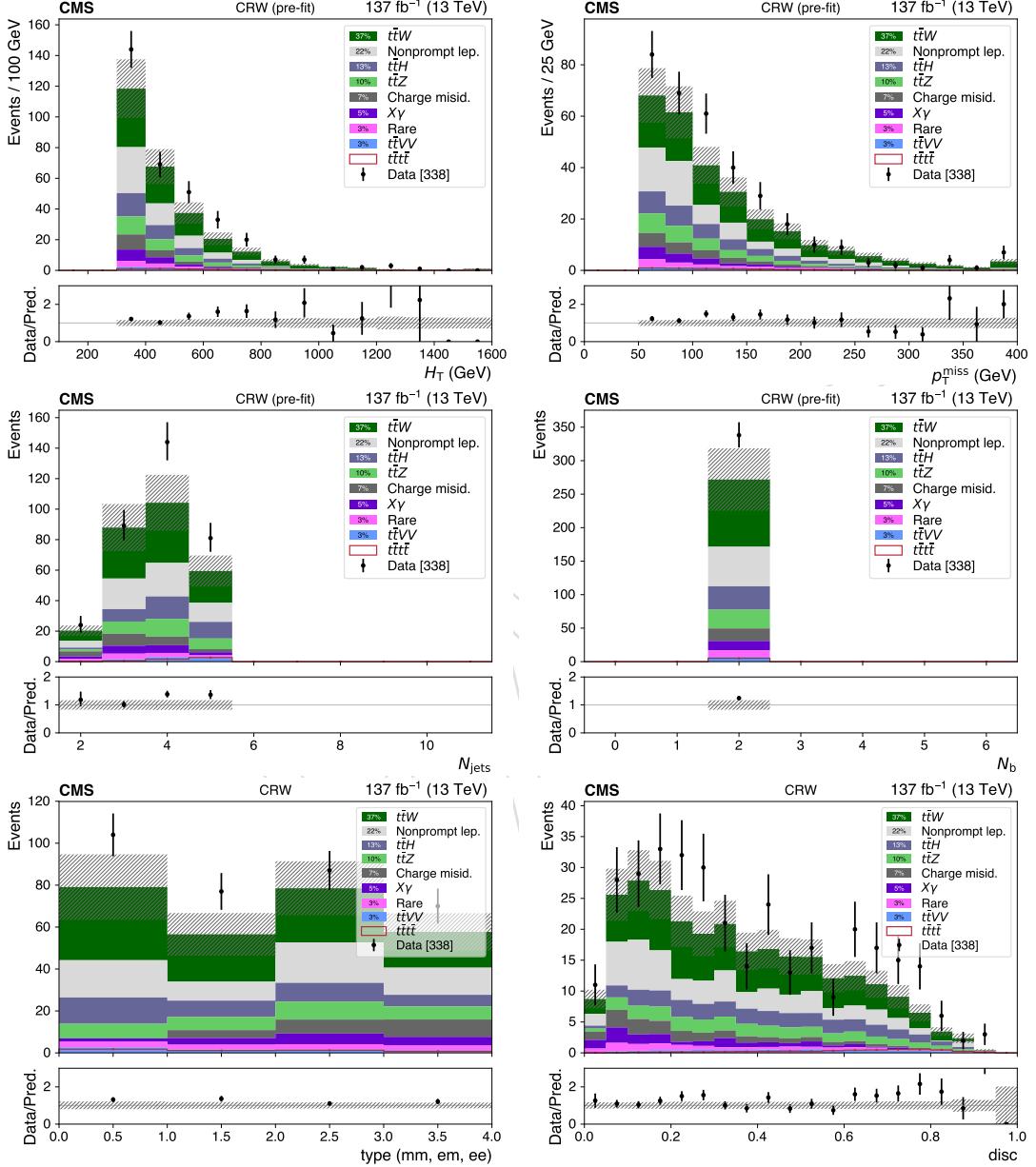


Figure 46: Run2 data and predictions: Prefit distributions of the main analysis variables in the  $t\bar{t}W$  control region:  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_b$  jets, lepton flavor, and raw BDT discriminant where the last bin includes the overflow.

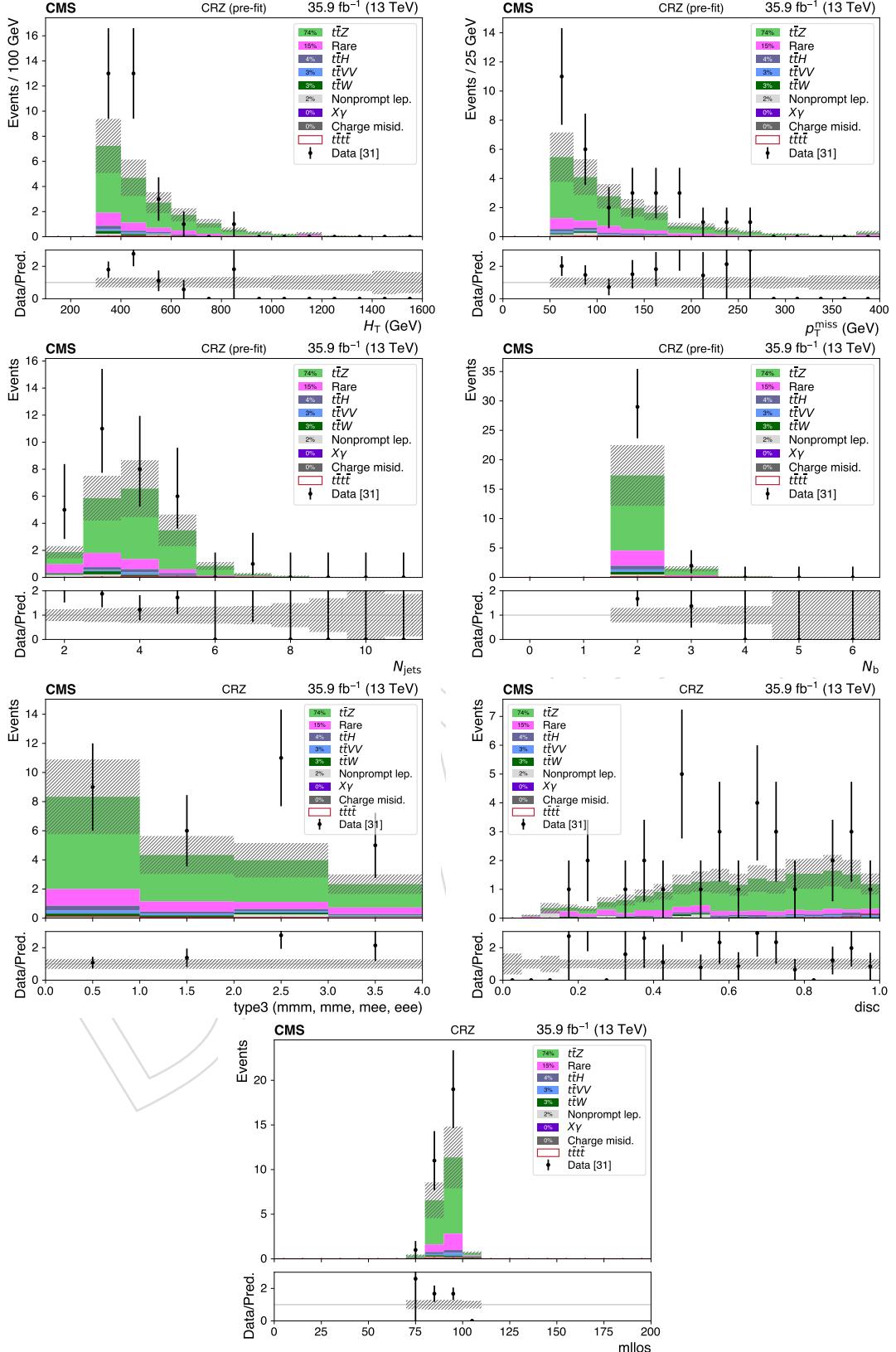


Figure 47: 2016 data and predictions: Prefit distributions of the main analysis variables in the  $t\bar{t}Z$  control region:  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_b$  jets, lepton flavor, raw BDT discriminant, and  $m_{\ell\ell}$  of the OSSF pair, where the last bin includes the overflow.

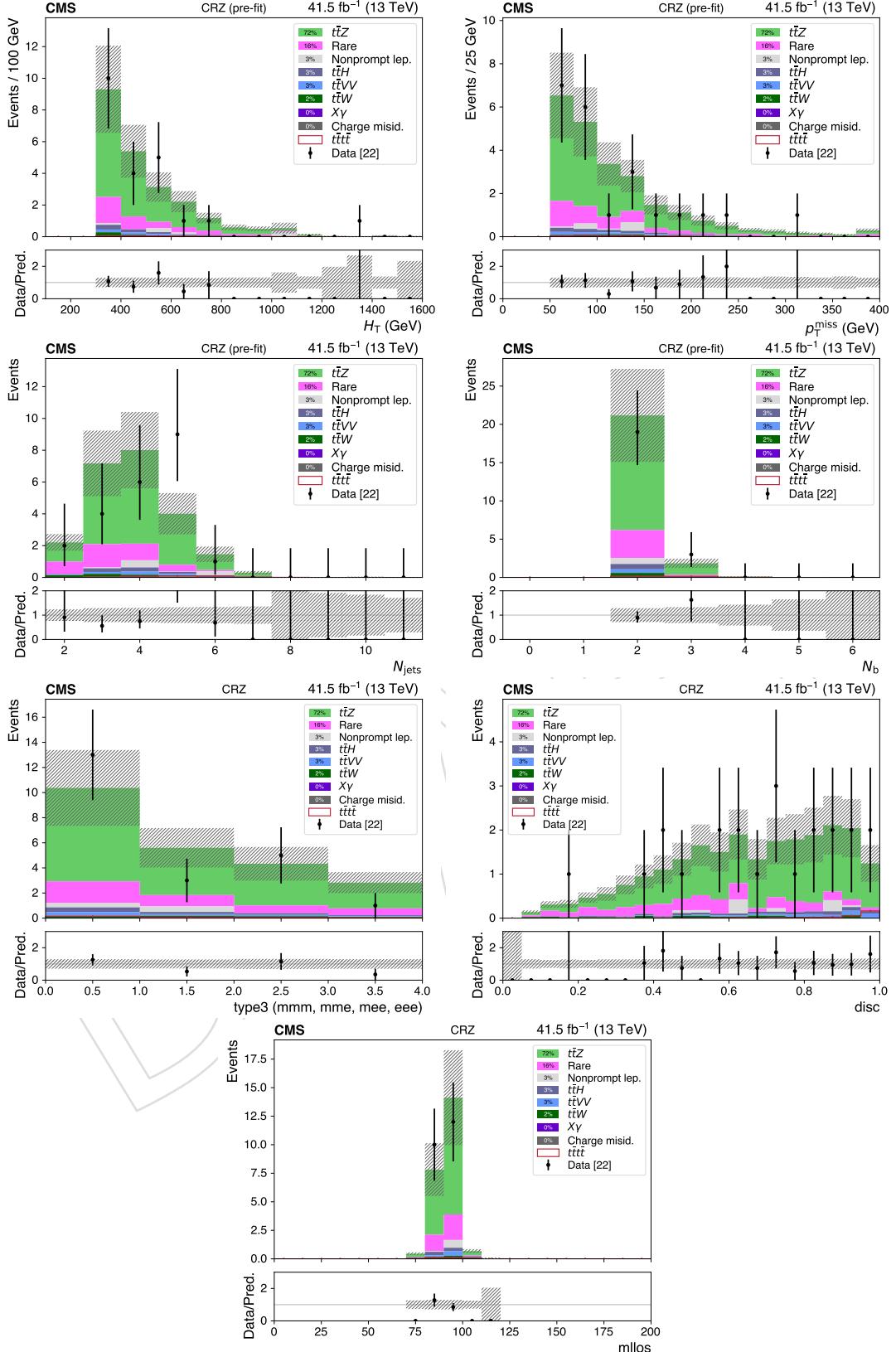


Figure 48: 2017 data and predictions: Prefit distributions of the main analysis variables in the  $t\bar{t}Z$  control region:  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_b$  jets, lepton flavor, raw BDT discriminant, and  $m_{\ell\ell}$  of the OSSF pair, where the last bin includes the overflow.

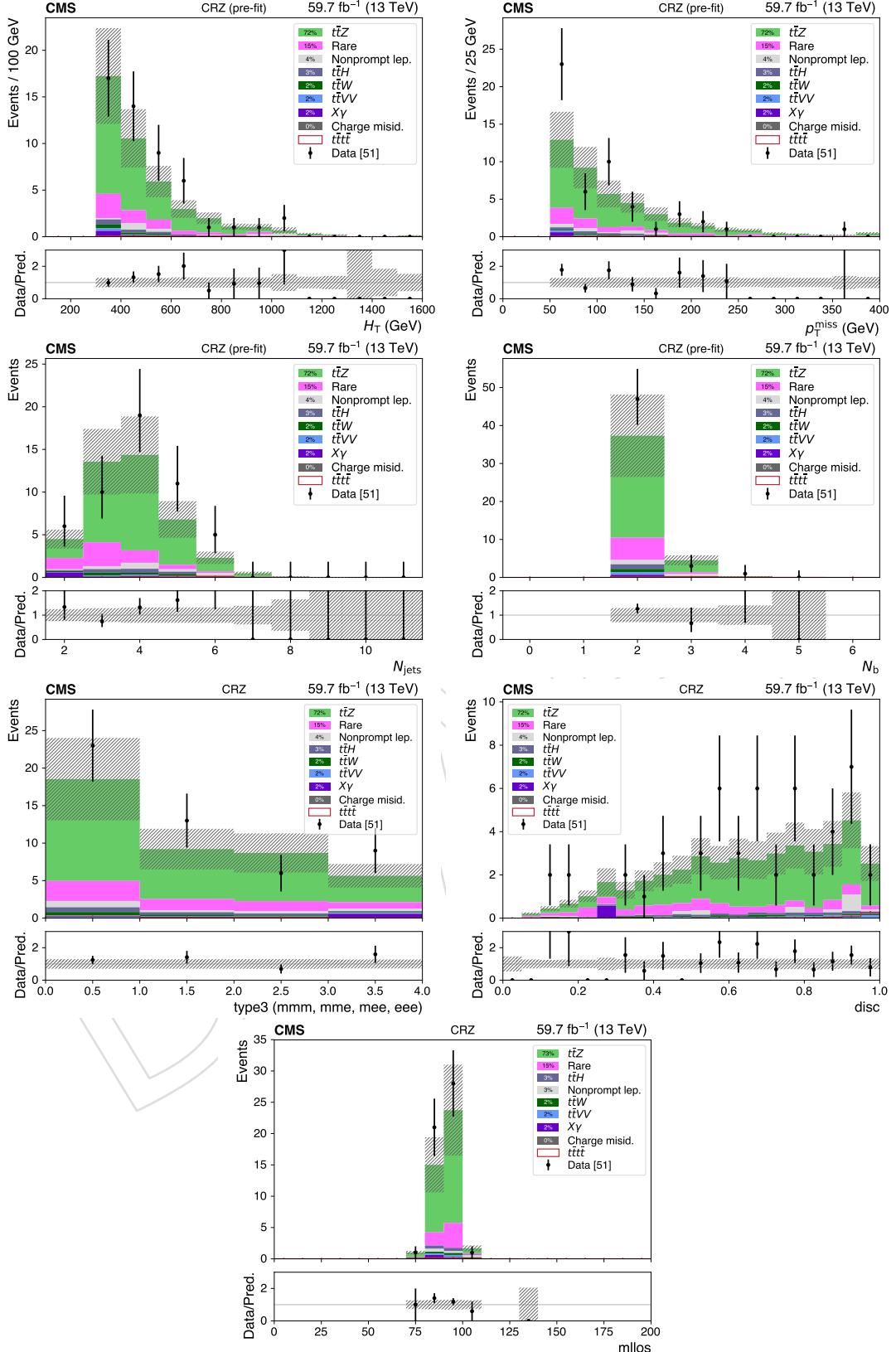


Figure 49: 2018 data and predictions: Prefit distributions of the main analysis variables in the  $t\bar{t}Z$  control region:  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_b$  jets, lepton flavor, raw BDT discriminant, and  $m_{\ell\ell}$  of the OSSF pair, where the last bin includes the overflow.

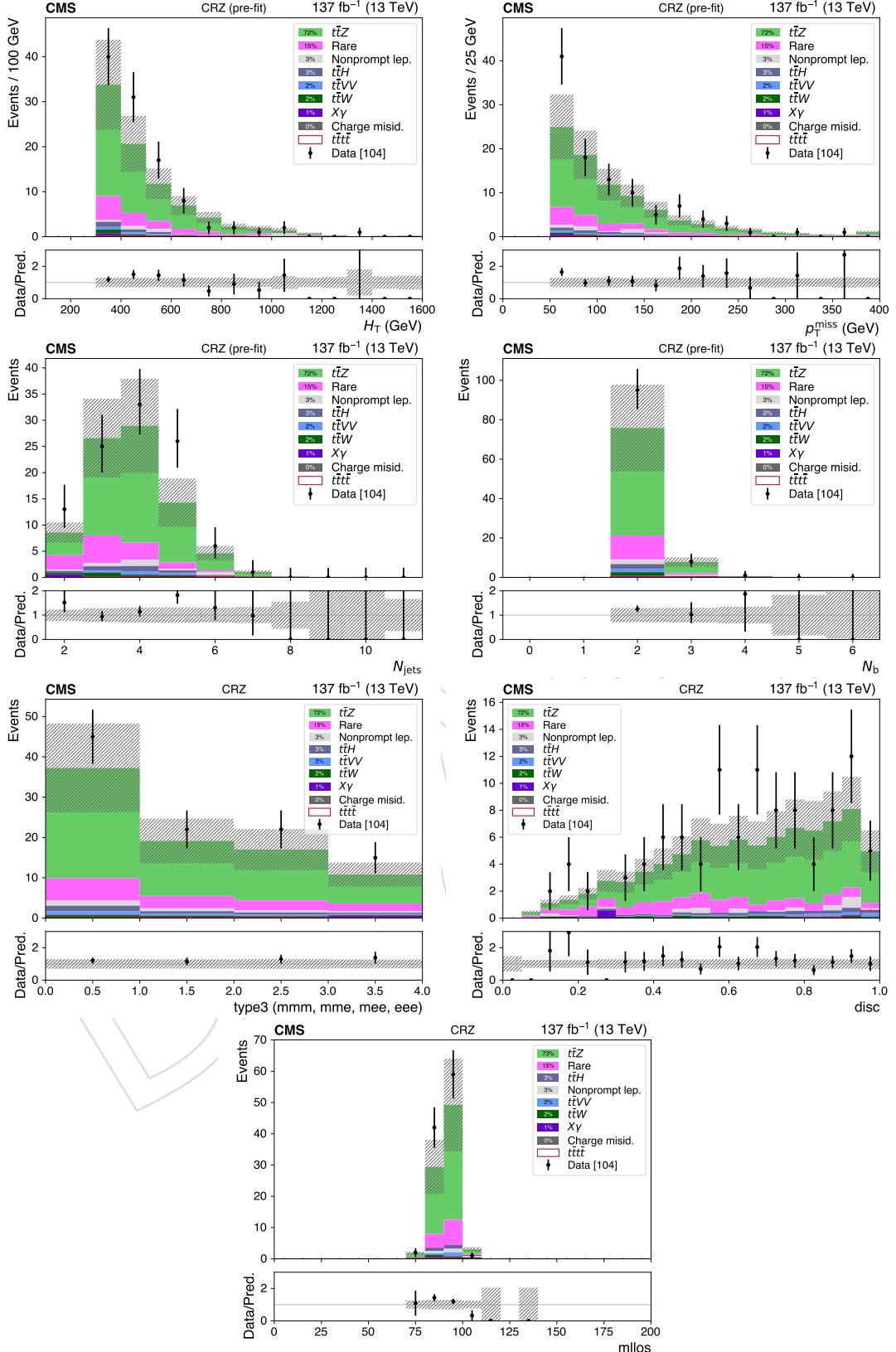


Figure 50: Run2 data and predictions: Prefit distributions of the main analysis variables in the  $t\bar{t}Z$  control region:  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_b$  jets, lepton flavor, raw BDT discriminant, and  $m_{\ell\ell}$  of the OSSF pair, where the last bin includes the overflow.

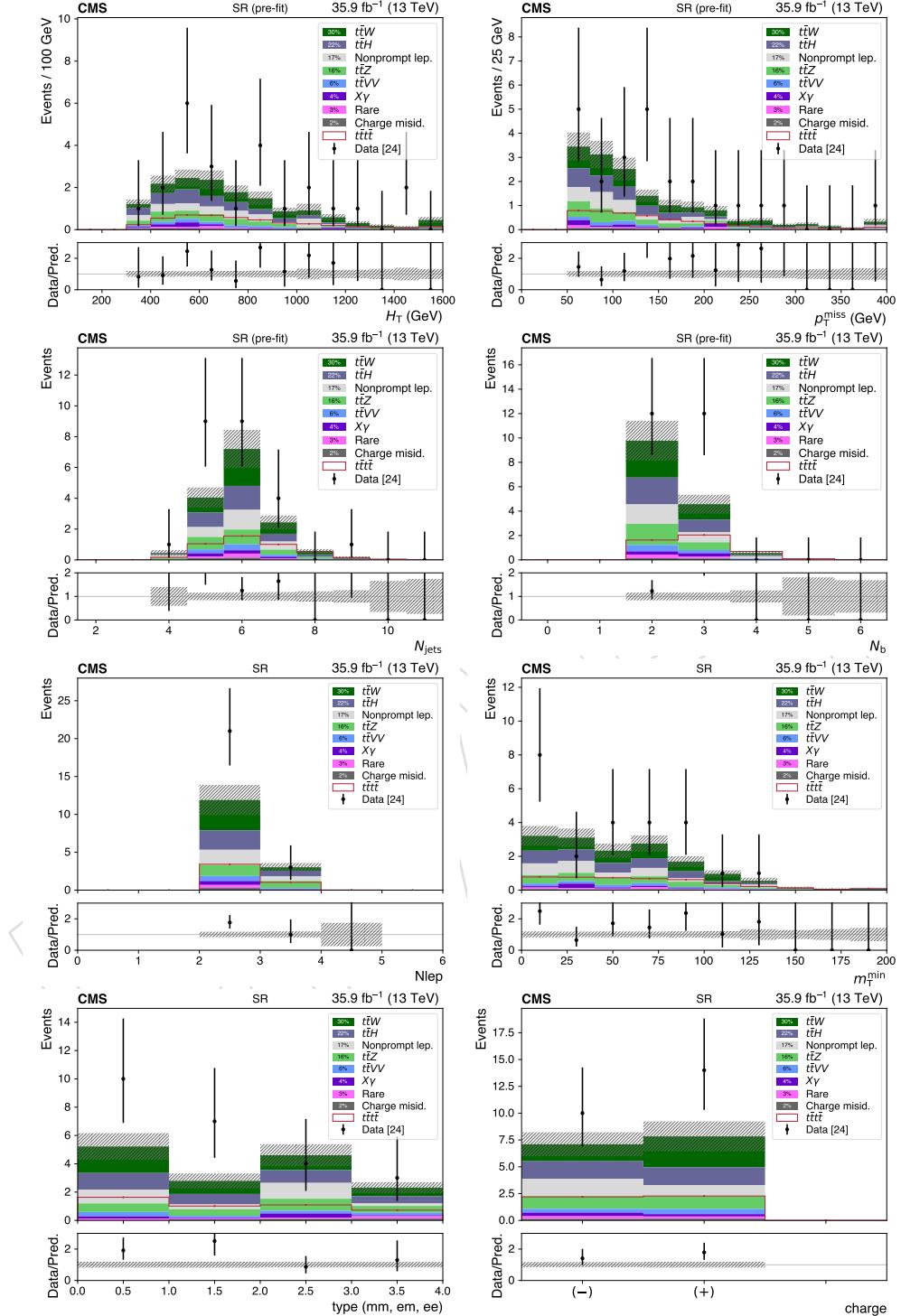


Figure 51: 2016 predictions: prefit distributions of the main analysis variables in the sum of signal regions:  $H_T$ ,  $E_T^{\text{miss}}$ ,  $M_T^{\text{min}}$ ,  $N_{\text{jets}}$ ,  $N_b$ , lepton flavor, and  $m_{\ell\ell}$  of the OSSF pair, where the last bin includes the overflow.

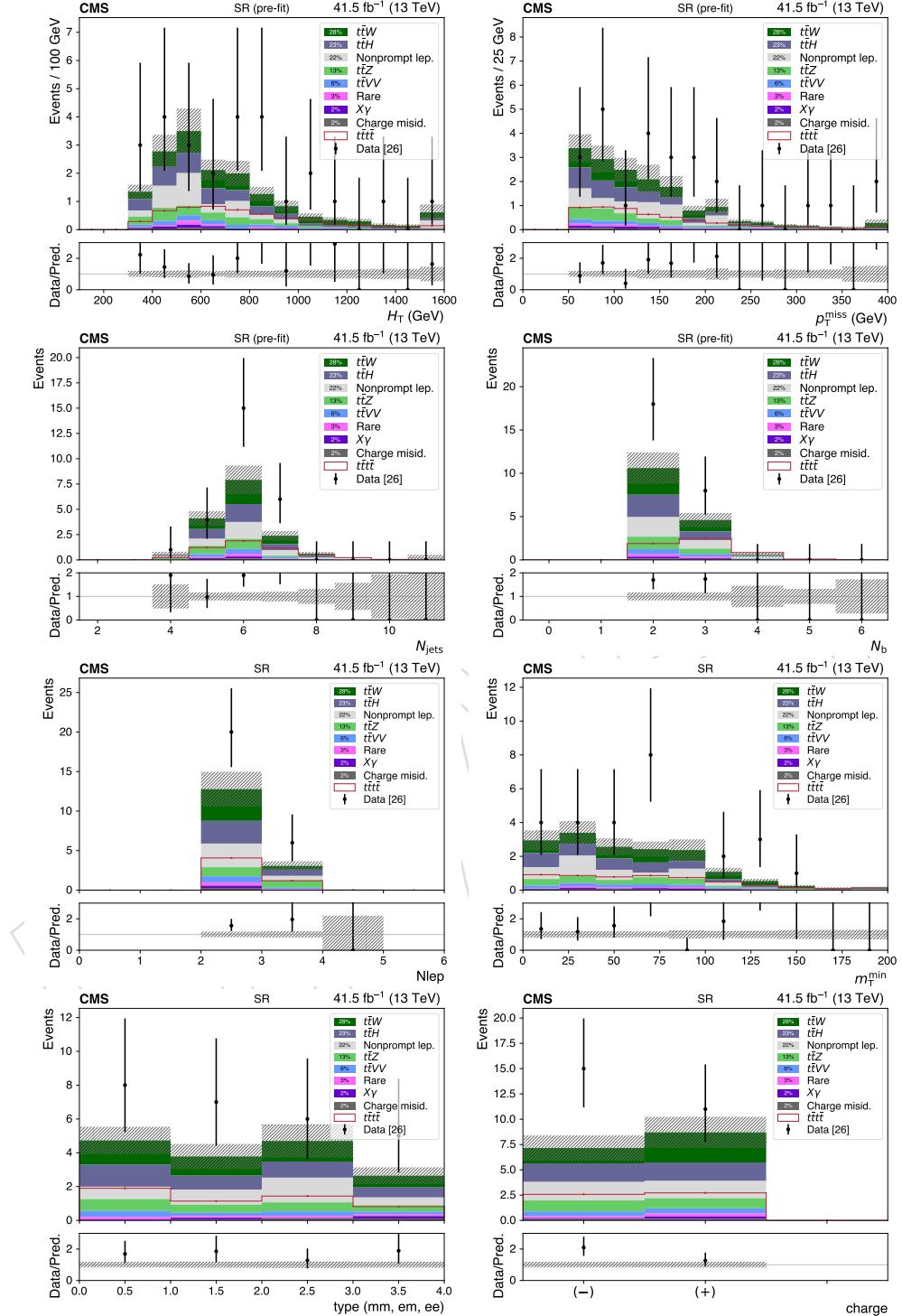


Figure 52: 2017 predictions: prefit distributions of the main analysis variables in the sum of signal regions:  $H_T$ ,  $E_T^{\text{miss}}$ ,  $M_T^{\text{min}}$ ,  $N_{\text{jets}}$ ,  $N_b$  jets, lepton flavor, and  $m_{\ell\ell}$  of the OSSF pair, where the last bin includes the overflow.

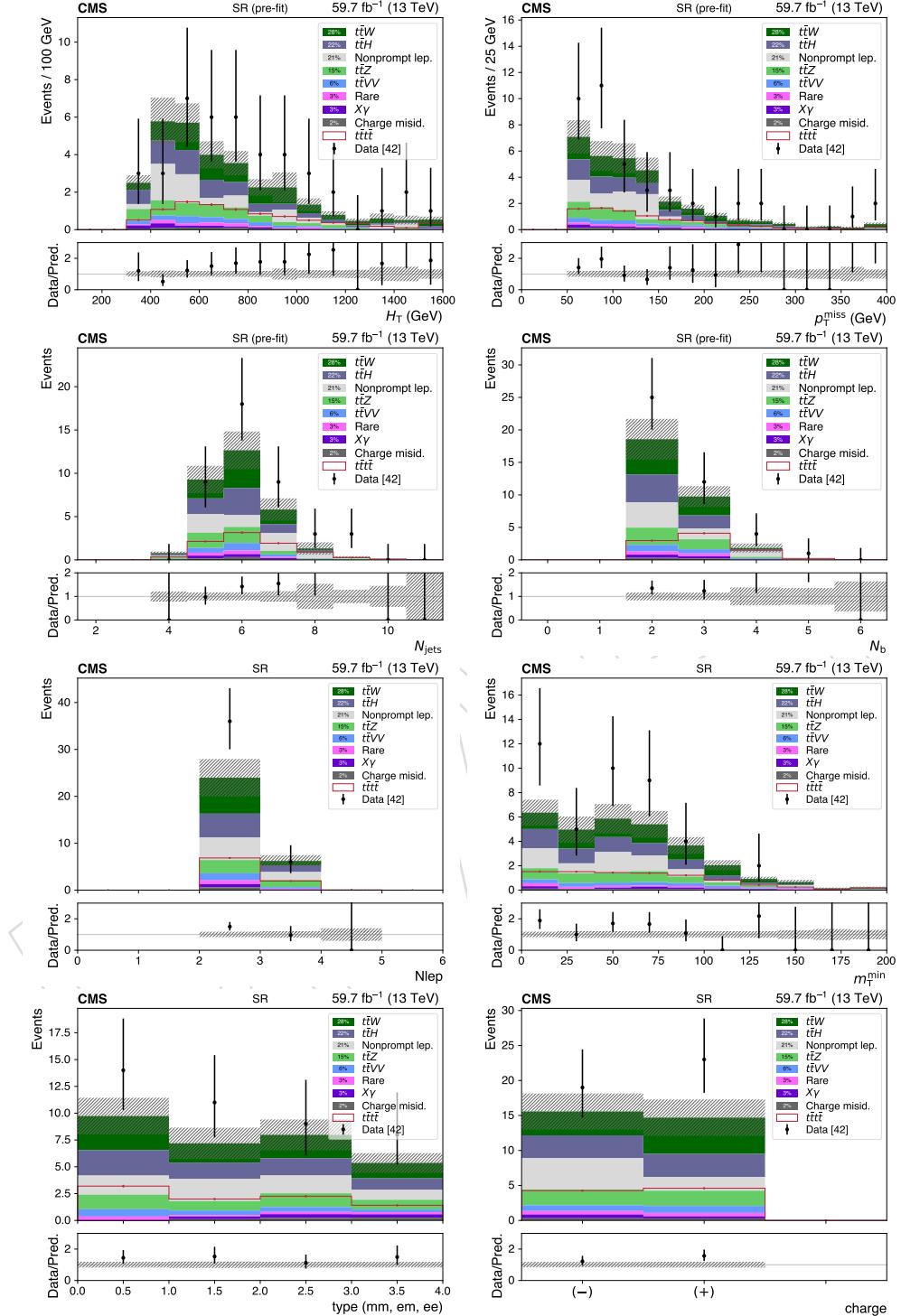


Figure 53: 2018 predictions: prefit distributions of the main analysis variables in the sum of signal regions:  $H_T$ ,  $E_T^{\text{miss}}$ ,  $M_T^{\text{min}}$ ,  $N_{\text{jets}}$ ,  $N_b$ , lepton flavor, and  $m_{\ell\ell}$  of the OSSF pair, where the last bin includes the overflow.

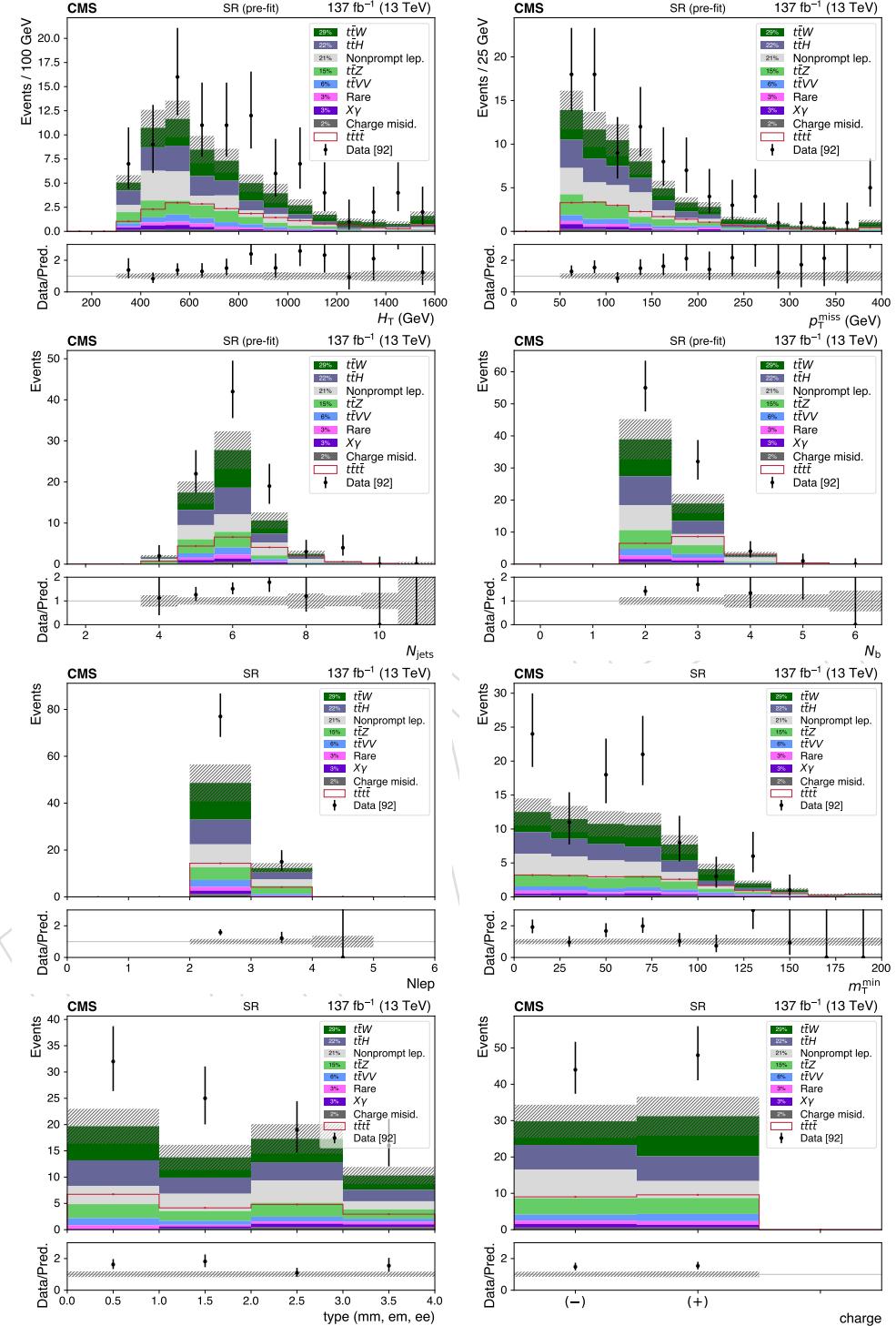


Figure 54: run2 predictions: prefit distributions of the main analysis variables in the sum of signal regions:  $H_T$ ,  $E_T^{\text{miss}}$ ,  $M_T^{\text{min}}$ ,  $N_{\text{jets}}$ ,  $N_b$  jets, lepton flavor, and  $m_{\ell\ell}$  of the OSSF pair, where the last bin includes the overflow.

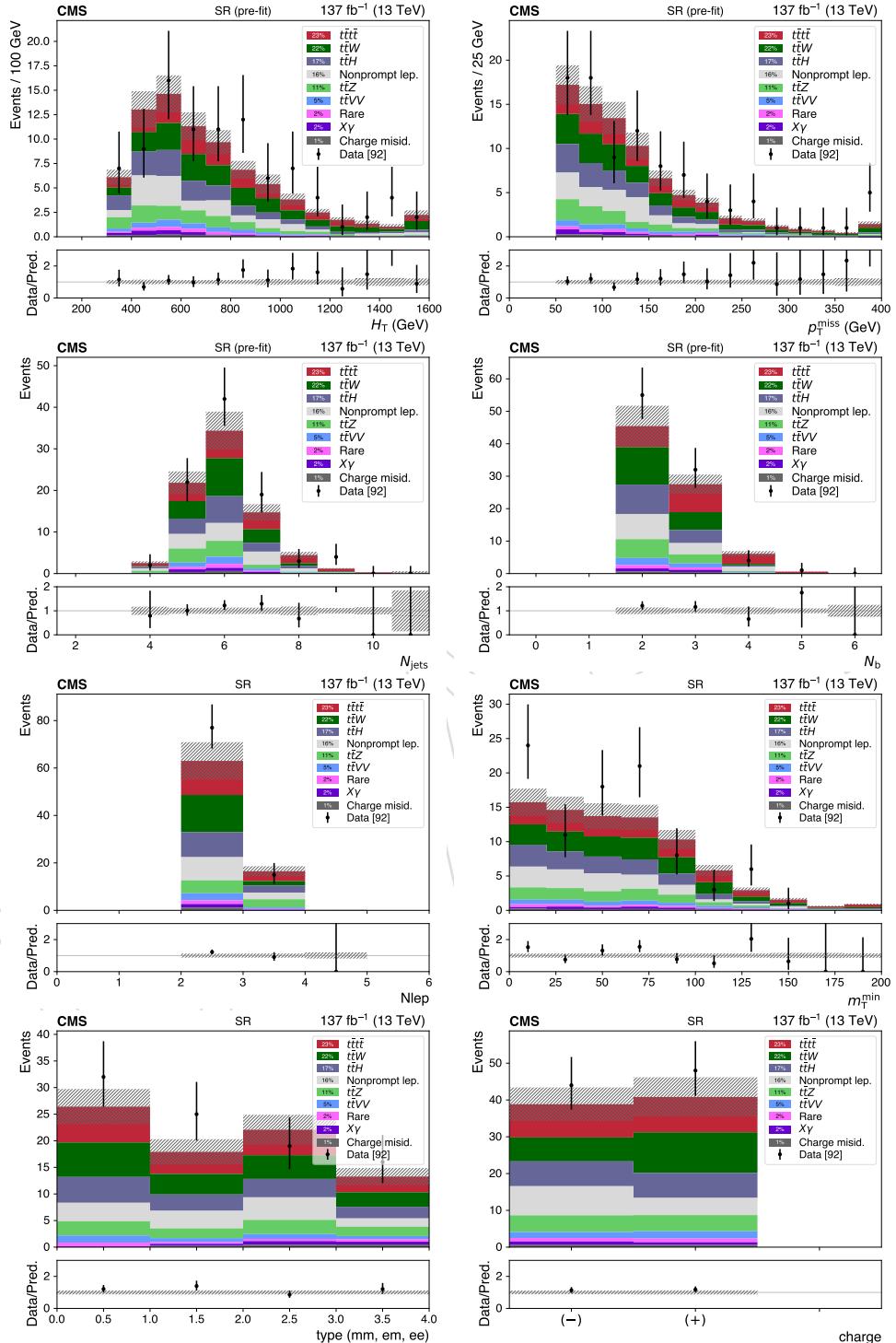


Figure 55: Same as Figure 54 but with stacked signal, assuming signal strength of 1.

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## 924 11 Results

925 Results for the 2016+2017+2018 (“Run 2”) data are obtained with the data in “golden JSON”  
 926 given in Sec.2 , for a total luminosity of  $77.4 \text{ fb}^{-1}$ . The expected yields in the signal regions  
 927 as well as the  $t\bar{t}W$  and  $t\bar{t}Z$  control regions, and BDT regions are shown in Figure 56 for prefit  
 928 for 2016/2017/2018/Run2. The postfits plot for the two analyses are shown in Figure 57. The  
 929 2016,2017,2018,Run 2 cut-based numerical yields can be found in Table 18,19,20,21 for prefit,  
 930 and analogous tables for the BDT analysis can be found in Table 22,23,24,25. Postfit tables for  
 931 Run 2 cut-based and BDT can be found in Table 26 and Table 27, respectively. . Note that with  
 932 respect to 2016, also the theoretical cross section assumed for the  $t\bar{t}t\bar{t}$  process has changed, from  
 933 9.2 to 11.96 fb.

934 A binned likelihood fit is performed over signal regions using the Higgs Combine tool [22];  
 935 exclusion limits at 95% CL are calculated with the Asymptotic CLs method.

936 For reference, the commands used to extract upper limits, significances, measurement values,  
 937 and obtain the NLL vs  $\mu$ , respectively, are:

```
938 • combine -M Asymptotic card.txt --noFitAsimov
939 • combine -M Significance card.txt -t -1 --expectSignal=1 --significance
940 • combine -M FitDiagnostics card.txt -t -1 --expectSignal=1 --robustFit=1
941   --saveShapes --saveOverallShapes --saveWithUncertainties -n name
942 • combine -M MultiDimFit card.txt -t -1 --expectSignal=1 --algo grid
943   --centeredRange=2.0 --saveFitResult --redefineSignalPOI r --robustFit=1
944   --saveNLL
```

945 -t -1 --expectSignal=1 and --noFitAsimov are omitted for observed results.

946 Using the Run2 data, the *cut-based* analysis sets an observed (expected) upper limit on the  
 947 production cross section of  $20.04 \text{ fb}$  ( $9.35^{+4.29}_{-2.88} \text{ fb}$ ), assuming the signal process does not exist.  
 948 The observed (expected) significance is  $1.712$  ( $2.478$ ) standard deviations, corresponding to a  
 949 measured observed (expected) signal strength parameter of  $0.784^{+0.514}_{-0.469}$  ( $1.000^{+0.480}_{-0.433}$ )

950 The *BDT* analysis sets an observed (expected) upper limit on the production cross section of  
 951  $22.51 \text{ fb}$  ( $8.46^{+3.91}_{-2.57} \text{ fb}$ ), assuming the signal process does not exist. The observed (expected)  
 952 significance is  $2.561$  ( $2.699$ ) standard deviations, corresponding to a measured observed (ex-  
 953 pected) signal strength parameter of  $1.052^{+0.483}_{-0.437}$  ( $1.000^{+0.442}_{-0.401}$ )

954 The likelihood scan for Run2 is shown in Fig. 58.

955 Several interpretations of these results are presented in the following sections

- 956 • Section 12.1: 2HDM with a heavy scalar or pseudoscalar boson decaying to on-shell  
 957  $t\bar{t}$
- 958 • Section 12.2: top yukawa coupling constant
- 959 • Section 12.3: off-shell top-philic scalar or pseudoscalar boson
- 960 • Section 12.4: EFT oblique Higgs parameter  $\hat{H}$
- 961 • Section 12.5: Simplified dark matter with scalar or pseudoscalar mediator

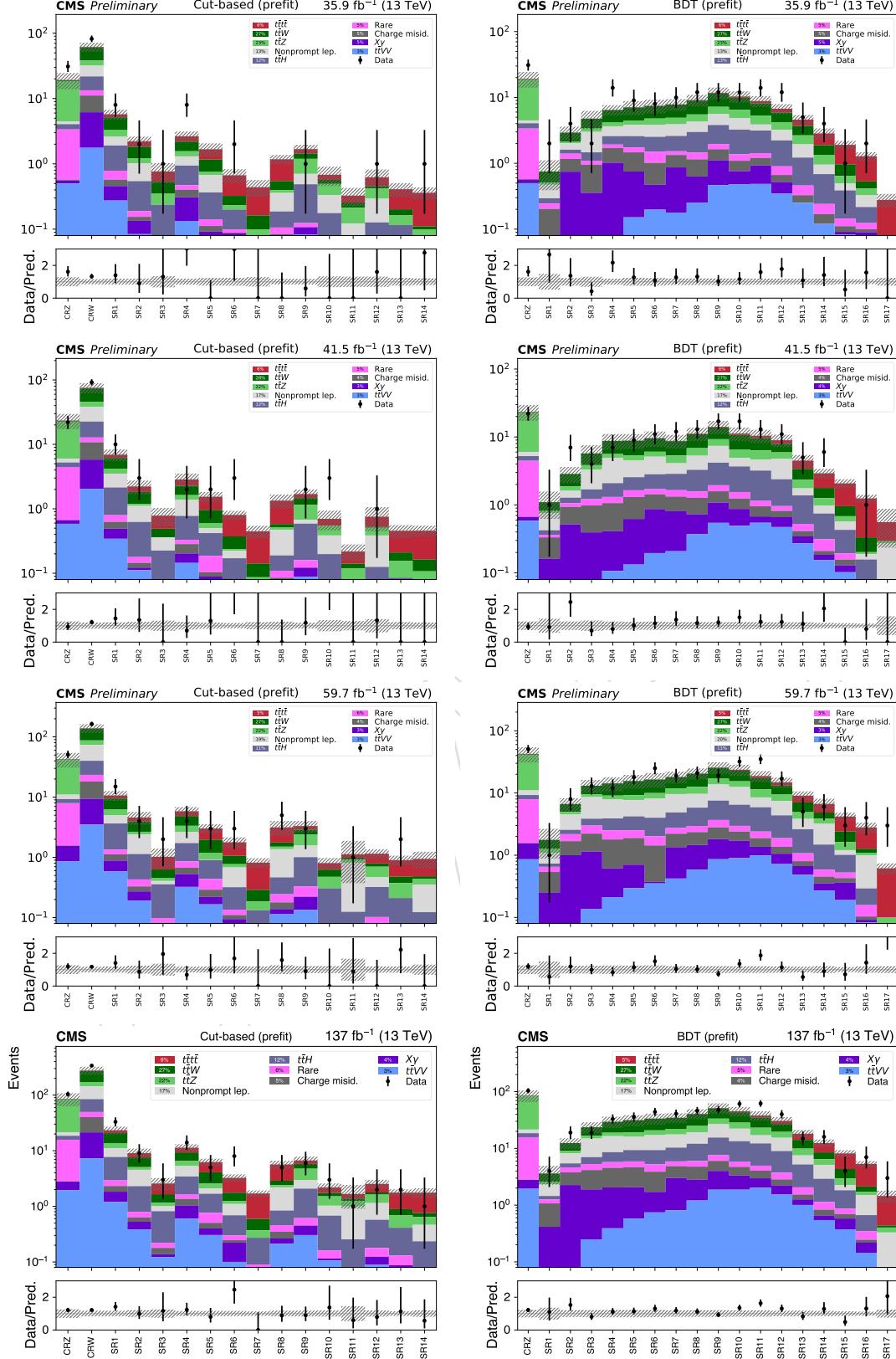


Figure 56: Data yields compared to prefit for cut-based (left) and BDT (right) analyses separately for data periods 2016, 2017, 2018, Run2 from top to bottom.

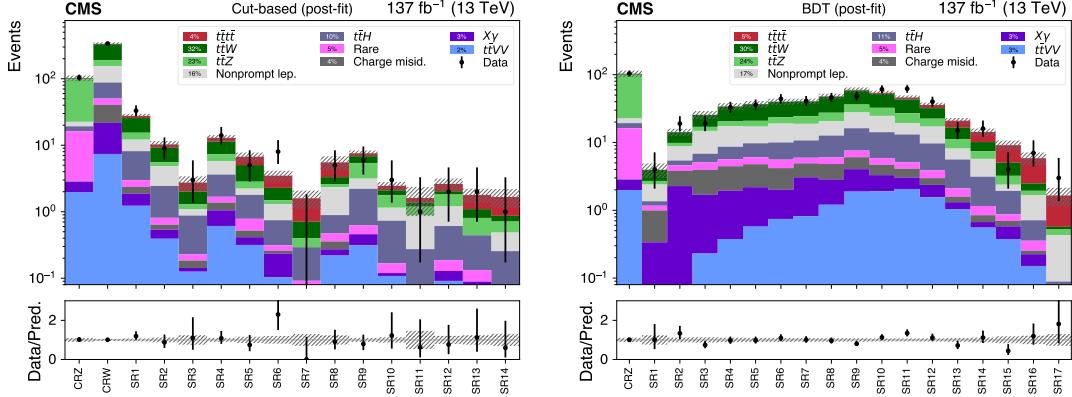


Figure 57: Data yields compared to postfit for cut-based (left) and BDT (right) analyses.

Table 18: Prefit event yields in SR+CR regions for 2016.

	$t\bar{t}W$	$t\bar{t}Z$	$t\bar{t}H$	$t\bar{t}VV$	$X+\gamma$	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	$t\bar{t}\bar{t}$
CRZ	$0.50 \pm 0.18$	$14.01 \pm 5.20$	$0.69 \pm 0.18$	$0.50 \pm 0.06$	$0.06 \pm 0.01$	$2.79 \pm 0.58$	$0.02 \pm 0.00$	$0.42 \pm 0.17$	$18.98 \pm 5.34$	31	$0.25 \pm 0.04$
CRW	$22.56 \pm 8.01$	$6.31 \pm 2.31$	$8.44 \pm 2.14$	$1.76 \pm 0.21$	$4.41 \pm 0.71$	$2.24 \pm 0.49$	$4.95 \pm 0.95$	$10.05 \pm 3.82$	$60.72 \pm 9.61$	82	$1.23 \pm 0.10$
SR1	$1.87 \pm 0.70$	$0.68 \pm 0.26$	$1.12 \pm 0.31$	$0.27 \pm 0.04$	$0.18 \pm 0.07$	$0.21 \pm 0.06$	$0.12 \pm 0.02$	$0.70 \pm 0.32$	$5.16 \pm 0.97$	8	$0.58 \pm 0.07$
SR2	$0.61 \pm 0.25$	$0.20 \pm 0.10$	$0.35 \pm 0.10$	$0.08 \pm 0.02$	$0.05 \pm 0.01$	$0.03 \pm 0.01$	$0.02 \pm 0.00$	$0.48 \pm 0.23$	$1.83 \pm 0.39$	2	$0.39 \pm 0.02$
SR3	$0.18 \pm 0.12$	$0.11 \pm 0.06$	$0.17 \pm 0.07$	$0.05 \pm 0.01$	$0.00 \pm 0.01$	$0.01 \pm 0.01$	$0.01 \pm 0.00$	$0.00 \pm 0.11$	$0.53 \pm 0.24$	1	$0.23 \pm 0.05$
SR4	$0.63 \pm 0.25$	$0.26 \pm 0.14$	$0.47 \pm 0.13$	$0.13 \pm 0.03$	$0.17 \pm 0.12$	$0.07 \pm 0.01$	$0.09 \pm 0.02$	$0.36 \pm 0.21$	$2.19 \pm 0.48$	8	$0.46 \pm 0.05$
SR5	$0.40 \pm 0.19$	$0.09 \pm 0.03$	$0.22 \pm 0.08$	$0.06 \pm 0.02$	$0.03 \pm 0.01$	$0.03 \pm 0.01$	$0.02 \pm 0.00$	$0.31 \pm 0.29$	$1.17 \pm 0.39$	0	$0.51 \pm 0.05$
SR6	$0.09 \pm 0.06$	$0.03 \pm 0.05$	$0.10 \pm 0.04$	$0.01 \pm 0.01$	$0.07 \pm 0.05$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.00 \pm 0.04$	$0.32 \pm 0.15$	2	$0.34 \pm 0.04$
SR7	$0.06 \pm 0.04$	$0.04 \pm 0.02$	$0.05 \pm 0.02$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.00 \pm 0.05$	$0.16 \pm 0.10$	0	$0.28 \pm 0.08$
SR8	$0.15 \pm 0.08$	$0.04 \pm 0.03$	$0.08 \pm 0.03$	$0.05 \pm 0.01$	$0.00 \pm 0.00$	$0.02 \pm 0.01$	$0.03 \pm 0.01$	$0.17 \pm 0.11$	$0.55 \pm 0.17$	0	$0.63 \pm 0.06$
SR9	$0.25 \pm 0.10$	$0.49 \pm 0.18$	$0.36 \pm 0.10$	$0.08 \pm 0.02$	$0.02 \pm 0.00$	$0.02 \pm 0.01$	$0.00 \pm 0.00$	$0.22 \pm 0.11$	$1.44 \pm 0.28$	1	$0.24 \pm 0.06$
SR10	$0.08 \pm 0.03$	$0.16 \pm 0.07$	$0.14 \pm 0.04$	$0.02 \pm 0.01$	$0.00 \pm 0.00$	$0.01 \pm 0.01$	$0.00 \pm 0.00$	$0.16 \pm 0.19$	$0.57 \pm 0.22$	0	$0.11 \pm 0.03$
SR11	$0.04 \pm 0.02$	$0.09 \pm 0.05$	$0.05 \pm 0.02$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.01 \pm 0.01$	$0.00 \pm 0.00$	$0.05 \pm 0.04$	$0.26 \pm 0.10$	0	$0.07 \pm 0.02$
SR12	$0.04 \pm 0.02$	$0.13 \pm 0.05$	$0.09 \pm 0.03$	$0.02 \pm 0.00$	$0.00 \pm 0.00$	$0.02 \pm 0.00$	$0.00 \pm 0.00$	$0.17 \pm 0.18$	$0.46 \pm 0.19$	1	$0.17 \pm 0.02$
SR13	$0.04 \pm 0.02$	$0.05 \pm 0.03$	$0.08 \pm 0.02$	$0.01 \pm 0.00$	$0.02 \pm 0.00$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.00 \pm 0.06$	$0.20 \pm 0.09$	0	$0.21 \pm 0.02$
SR14	$0.01 \pm 0.02$	$0.05 \pm 0.03$	$0.04 \pm 0.02$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.00 \pm 0.00$	$0.00 \pm 0.00$	$0.00 \pm 0.00$	$0.11 \pm 0.05$	1	$0.25 \pm 0.05$

Table 19: Prefit event yields in SR+CR regions for 2017.

	$t\bar{t}W$	$t\bar{t}Z$	$t\bar{t}H$	$t\bar{t}VV$	$X+\gamma$	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	$t\bar{t}\bar{t}$
CRZ	$0.53 \pm 0.20$	$16.57 \pm 6.03$	$0.77 \pm 0.22$	$0.59 \pm 0.08$	$0.08 \pm 0.04$	$3.79 \pm 0.81$	$0.00 \pm 0.00$	$0.79 \pm 0.42$	$23.12 \pm 6.24$	22	$0.32 \pm 0.02$
CRW	$28.09 \pm 9.95$	$7.70 \pm 2.83$	$9.81 \pm 2.67$	$2.04 \pm 0.25$	$3.77 \pm 0.57$	$1.99 \pm 0.42$	$4.98 \pm 0.95$	$15.62 \pm 8.52$	$73.99 \pm 13.76$	92	$1.48 \pm 0.07$
SR1	$1.97 \pm 0.72$	$0.70 \pm 0.27$	$1.35 \pm 0.38$	$0.34 \pm 0.05$	$0.15 \pm 0.06$	$0.17 \pm 0.06$	$0.14 \pm 0.03$	$1.35 \pm 1.05$	$6.17 \pm 1.42$	10	$0.75 \pm 0.04$
SR2	$0.64 \pm 0.28$	$0.08 \pm 0.07$	$0.45 \pm 0.14$	$0.11 \pm 0.02$	$0.01 \pm 0.01$	$0.03 \pm 0.01$	$0.02 \pm 0.00$	$0.48 \pm 0.37$	$1.83 \pm 0.52$	3	$0.39 \pm 0.04$
SR3	$0.13 \pm 0.09$	$0.00 \pm 0.02$	$0.17 \pm 0.06$	$0.03 \pm 0.01$	$0.00 \pm 0.00$	$0.02 \pm 0.00$	$0.01 \pm 0.00$	$0.14 \pm 0.16$	$0.49 \pm 0.24$	0	$0.29 \pm 0.05$
SR4	$0.65 \pm 0.25$	$0.26 \pm 0.14$	$0.45 \pm 0.14$	$0.14 \pm 0.02$	$0.06 \pm 0.04$	$0.05 \pm 0.02$	$0.08 \pm 0.02$	$0.60 \pm 0.49$	$2.28 \pm 0.67$	2	$0.57 \pm 0.05$
SR5	$0.36 \pm 0.20$	$0.10 \pm 0.04$	$0.23 \pm 0.08$	$0.08 \pm 0.01$	$0.01 \pm 0.01$	$0.08 \pm 0.02$	$0.01 \pm 0.00$	$0.08 \pm 0.09$	$0.96 \pm 0.27$	2	$0.58 \pm 0.04$
SR6	$0.14 \pm 0.06$	$0.01 \pm 0.04$	$0.11 \pm 0.04$	$0.03 \pm 0.00$	$0.01 \pm 0.00$	$0.02 \pm 0.01$	$0.01 \pm 0.00$	$0.05 \pm 0.06$	$0.37 \pm 0.12$	3	$0.43 \pm 0.03$
SR7	$0.05 \pm 0.03$	$0.01 \pm 0.01$	$0.04 \pm 0.02$	$0.01 \pm 0.00$	$0.01 \pm 0.01$	$0.01 \pm 0.00$	$0.01 \pm 0.00$	$0.00 \pm 0.06$	$0.14 \pm 0.08$	0	$0.30 \pm 0.04$
SR8	$0.07 \pm 0.05$	$0.04 \pm 0.03$	$0.08 \pm 0.04$	$0.05 \pm 0.01$	$0.02 \pm 0.01$	$0.02 \pm 0.01$	$0.01 \pm 0.00$	$0.28 \pm 0.28$	$0.58 \pm 0.30$	0	$0.76 \pm 0.08$
SR9	$0.32 \pm 0.16$	$0.48 \pm 0.18$	$0.43 \pm 0.12$	$0.09 \pm 0.01$	$0.03 \pm 0.01$	$0.04 \pm 0.01$	$0.00 \pm 0.00$	$0.08 \pm 0.07$	$1.47 \pm 0.30$	2	$0.23 \pm 0.01$
SR10	$0.00 \pm 0.03$	$0.16 \pm 0.07$	$0.15 \pm 0.05$	$0.03 \pm 0.01$	$0.01 \pm 0.01$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.19 \pm 0.21$	$0.55 \pm 0.23$	3	$0.15 \pm 0.02$
SR11	$0.02 \pm 0.01$	$0.06 \pm 0.03$	$0.05 \pm 0.02$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.00 \pm 0.00$	$0.00 \pm 0.00$	$0.14 \pm 0.06$	0	$0.08 \pm 0.02$	
SR12	$0.02 \pm 0.01$	$0.12 \pm 0.05$	$0.08 \pm 0.03$	$0.03 \pm 0.00$	$0.01 \pm 0.00$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.26 \pm 0.30$	$0.53 \pm 0.30$	1	$0.23 \pm 0.02$
SR13	$0.05 \pm 0.03$	$0.07 \pm 0.03$	$0.06 \pm 0.02$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.00 \pm 0.06$	$0.21 \pm 0.09$	0	$0.25 \pm 0.02$
SR14	$0.06 \pm 0.03$	$0.05 \pm 0.03$	$0.04 \pm 0.02$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.01 \pm 0.00$	$0.00 \pm 0.00$	$0.16 \pm 0.08$	0	$0.29 \pm 0.03$	

Table 20: Prefit event yields in SR+CR regions for 2018.

	$t\bar{t}W$	$t\bar{t}Z$	$t\bar{t}H$	$t\bar{t}VV$	$X+\gamma$	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	$t\bar{t}\bar{t}$
CRZ	$0.88 \pm 0.33$	$30.20 \pm 11.07$	$1.34 \pm 0.36$	$0.86 \pm 0.12$	$0.70 \pm 0.40$	$6.36 \pm 1.28$	$0.01 \pm 0.00$	$1.72 \pm 0.60$	$42.07 \pm 11.49$	51	$0.49 \pm 0.03$
CRW	$49.69 \pm 18.72$	$13.94 \pm 4.93$	$16.64 \pm 4.19$	$3.50 \pm 0.43$	$5.88 \pm 1.15$	$5.25 \pm 1.10$	$8.80 \pm 1.75$	$33.41 \pm 9.80$	$137.09 \pm 23.04$	164	$2.24 \pm 0.09$
SR1	$3.27 \pm 1.39$	$1.25 \pm 0.47$	$2.31 \pm 0.63$	$0.58 \pm 0.09$	$0.30 \pm 0.09$	$0.22 \pm 0.05$	$0.24 \pm 0.05$	$1.31 \pm 0.62$	$9.48 \pm 1.96$	15	$1.18 \pm 0.07$
SR2	$1.29 \pm 0.53$	$0.26 \pm 0.16$	$0.71 \pm 0.24$	$0.19 \pm 0.04$	$0.08 \pm 0.04$	$0.10 \pm 0.03$	$0.06 \pm 0.01$	$1.27 \pm 0.79$	$3.95 \pm 1.09$	4	$0.65 \pm 0.05$
SR3	$0.18 \pm 0.17$	$0.04 \pm 0.08$	$0.26 \pm 0.09$	$0.05 \pm 0.01$	$0.01 \pm 0.02$	$0.02 \pm 0.00$	$0.03 \pm 0.01$	$0.05 \pm 0.04$	$0.63 \pm 0.31$	2	$0.39 \pm 0.07$
SR4	$1.51 \pm 0.63$	$0.72 \pm 0.32$	$0.87 \pm 0.27$	$0.32 \pm 0.05$	$0.20 \pm 0.06$	$0.18 \pm 0.05$	$0.15 \pm 0.03$	$0.89 \pm 0.42$	$4.85 \pm 1.01$	4	$0.97 \pm 0.05$
SR5	$0.77 \pm 0.32$	$0.28 \pm 0.18$	$0.46 \pm 0.15$	$0.17 \pm 0.03$	$0.05 \pm 0.03$	$0.15 \pm 0.05$	$0.07 \pm 0.01$	$0.01 \pm 0.05$	$1.97 \pm 0.54$	3	$1.05 \pm 0.06$
SR6	$0.29 \pm 0.15$	$0.09 \pm 0.04$	$0.18 \pm 0.07$	$0.05 \pm 0.01$	$0.04 \pm 0.01$	$0.02 \pm 0.00$	$0.02 \pm 0.00$	$0.36 \pm 0.24$	$1.06 \pm 0.33$	3	$0.71 \pm 0.08$
SR7	$0.10 \pm 0.05$	$0.05 \pm 0.03$	$0.09 \pm 0.04$	$0.02 \pm 0.00$	$0.00 \pm 0.00$	$0.02 \pm 0.01$	$0.00 \pm 0.00$	$0.00 \pm 0.04$	$0.29 \pm 0.11$	0	$0.53 \pm 0.07$
SR8	$0.26 \pm 0.16$	$0.14 \pm 0.09$	$0.22 \pm 0.09$	$0.11 \pm 0.02$	$0.03 \pm 0.01$	$0.07 \pm 0.01$	$0.04 \pm 0.01$	$0.90 \pm 0.75$	$1.76 \pm 0.83$	5	$1.39 \pm 0.13$
SR9	$0.47 \pm 0.23$	$0.84 \pm 0.36$	$$								

Table 21: Prefit event yields in SR+CR regions for 2016+2017+2018.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	1.92± 0.70	60.77± 21.92	2.80± 0.72	1.94± 0.22	0.84± 0.40	12.94± 2.74	0.03± 0.00	2.94± 1.10	84.18± 22.57	104	1.05± 0.06
CRW	100.33± 36.18	27.95± 10.15	34.89± 8.58	7.30± 0.79	14.06± 1.83	9.48± 2.08	18.73± 2.14	59.08± 21.64	271.81± 46.38	338	4.94± 0.17
SR1	7.11± 2.70	2.63± 0.96	4.78± 1.29	1.20± 0.15	0.63± 0.14	0.60± 0.15	0.49± 0.06	3.36± 1.41	20.80± 3.77	33	2.50± 0.12
SR2	2.53± 0.98	0.53± 0.28	1.51± 0.46	0.39± 0.06	0.14± 0.05	0.17± 0.04	0.10± 0.01	2.23± 1.24	7.60± 1.88	9	1.43± 0.07
SR3	0.49± 0.33	0.16± 0.11	0.60± 0.22	0.12± 0.02	0.02± 0.02	0.04± 0.01	0.04± 0.01	0.19± 0.18	1.65± 0.61	3	0.92± 0.12
SR4	2.78± 1.02	1.24± 0.53	1.78± 0.48	0.60± 0.07	0.44± 0.14	0.29± 0.07	0.32± 0.04	1.85± 0.87	9.31± 1.75	14	2.00± 0.09
SR5	1.54± 0.64	0.47± 0.21	0.92± 0.29	0.31± 0.04	0.10± 0.04	0.26± 0.07	0.10± 0.01	0.40± 0.28	4.10± 0.97	5	2.13± 0.09
SR6	0.52± 0.20	0.14± 0.10	0.39± 0.14	0.10± 0.02	0.12± 0.05	0.05± 0.01	0.02± 0.00	0.41± 0.27	1.76± 0.48	8	1.49± 0.09
SR7	0.22± 0.09	0.10± 0.04	0.18± 0.07	0.04± 0.01	0.01± 0.01	0.03± 0.01	0.01± 0.00	0.00± 0.08	0.59± 0.19	0	1.11± 0.10
SR8	0.48± 0.24	0.22± 0.12	0.38± 0.14	0.21± 0.03	0.05± 0.02	0.11± 0.02	0.08± 0.01	1.36± 0.81	2.89± 0.93	5	2.77± 0.16
SR9	1.03± 0.44	1.80± 0.66	1.47± 0.39	0.30± 0.04	0.15± 0.02	0.16± 0.09	0.00± 0.00	0.86± 0.45	5.78± 1.16	6	0.90± 0.08
SR10	0.23± 0.10	0.49± 0.17	0.52± 0.16	0.11± 0.02	0.01± 0.01	0.04± 0.01	0.00± 0.00	0.35± 0.24	1.74± 0.39	3	0.44± 0.06
SR11	0.13± 0.07	0.25± 0.11	0.19± 0.06	0.04± 0.01	0.01± 0.01	0.02± 0.01	0.00± 0.00	0.74± 0.74	1.38± 0.77	1	0.29± 0.03
SR12	0.15± 0.08	0.48± 0.17	0.39± 0.10	0.09± 0.01	0.04± 0.01	0.06± 0.02	0.00± 0.00	0.57± 0.37	1.78± 0.46	2	0.75± 0.04
SR13	0.20± 0.10	0.28± 0.11	0.28± 0.09	0.07± 0.01	0.02± 0.01	0.04± 0.01	0.00± 0.00	0.00± 0.09	0.89± 0.24	2	0.88± 0.05
SR14	0.11± 0.05	0.18± 0.09	0.17± 0.07	0.04± 0.01	0.00± 0.00	0.03± 0.01	0.00± 0.00	0.23± 0.26	0.76± 0.32	1	1.00± 0.09

Table 22: Prefit event yields in BDT regions for 2016.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	0.50± 0.20	14.01± 5.02	0.69± 0.15	0.50± 0.06	0.06± 0.01	2.79± 0.59	0.02± 0.00	0.42± 0.19	18.98± 5.15	31	0.25± 0.04
SR1	0.24± 0.11	0.04± 0.03	0.05± 0.04	0.00± 0.00	0.00± 0.15	0.04± 0.11	0.20± 0.04	0.18± 0.12	0.75± 0.36	2	0.00± 0.00
SR2	0.77± 0.44	0.13± 0.08	0.21± 0.09	0.02± 0.01	0.72± 0.29	0.25± 0.10	0.42± 0.08	0.41± 0.22	2.92± 0.68	4	0.00± 0.00
SR3	1.38± 0.79	0.38± 0.15	0.37± 0.14	0.05± 0.02	0.30± 0.36	0.14± 0.09	0.60± 0.11	1.53± 0.72	4.76± 1.48	2	0.01± 0.00
SR4	2.23± 0.96	0.57± 0.26	0.66± 0.20	0.07± 0.04	0.95± 0.45	0.22± 0.08	0.69± 0.13	1.06± 0.42	6.44± 1.17	14	0.02± 0.02
SR5	2.61± 1.19	0.62± 0.29	0.84± 0.24	0.15± 0.03	0.60± 0.30	0.27± 0.13	0.72± 0.14	1.24± 0.51	7.05± 1.72	9	0.04± 0.01
SR6	2.65± 1.23	0.92± 0.38	1.02± 0.24	0.20± 0.04	0.27± 0.27	0.49± 0.14	0.54± 0.10	1.33± 0.67	7.43± 1.72	8	0.05± 0.04
SR7	2.87± 1.17	1.00± 0.48	1.08± 0.24	0.18± 0.04	0.69± 0.17	0.18± 0.11	0.47± 0.09	1.30± 0.56	7.78± 1.57	10	0.08± 0.04
SR8	3.44± 1.42	1.23± 0.60	1.53± 0.32	0.25± 0.05	0.38± 0.13	0.30± 0.11	0.52± 0.10	1.29± 0.61	8.96± 1.82	12	0.17± 0.03
SR9	4.24± 1.77	1.63± 0.73	1.85± 0.41	0.47± 0.06	0.65± 0.18	0.32± 0.13	0.52± 0.10	1.76± 0.72	11.44± 2.37	12	0.32± 0.04
SR10	3.40± 1.43	1.61± 0.77	1.88± 0.40	0.48± 0.06	0.28± 0.16	0.28± 0.06	0.33± 0.06	1.62± 0.77	9.88± 2.05	12	0.45± 0.04
SR11	2.83± 1.21	1.43± 0.64	1.72± 0.39	0.49± 0.05	0.45± 0.19	0.21± 0.05	0.28± 0.05	0.74± 0.47	8.14± 1.61	14	0.67± 0.07
SR12	1.68± 0.73	0.98± 0.47	1.33± 0.31	0.35± 0.06	0.17± 0.07	0.24± 0.06	0.14± 0.03	0.99± 0.51	5.88± 1.18	12	0.86± 0.04
SR13	1.03± 0.44	0.72± 0.31	0.76± 0.18	0.25± 0.03	0.00± 0.02	0.13± 0.03	0.08± 0.01	0.72± 0.38	3.68± 0.80	5	0.92± 0.04
SR14	0.56± 0.33	0.37± 0.14	0.44± 0.14	0.12± 0.02	0.00± 0.06	0.07± 0.01	0.05± 0.01	0.27± 0.18	1.89± 0.54	4	0.95± 0.05
SR15	0.32± 0.16	0.15± 0.09	0.24± 0.07	0.08± 0.01	0.01± 0.01	0.03± 0.01	0.02± 0.00	0.10± 0.07	0.96± 0.27	1	0.94± 0.06
SR16	0.12± 0.06	0.09± 0.05	0.09± 0.04	0.03± 0.01	0.06± 0.06	0.02± 0.00	0.01± 0.00	0.11± 0.09	0.53± 0.18	2	0.75± 0.07
SR17	0.03± 0.03	0.00± 0.01	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.05± 0.04	0	0.23± 0.05

Table 23: Prefit event yields in BDT regions for 2017.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	0.53± 0.22	16.57± 5.93	0.77± 0.18	0.59± 0.07	0.08± 0.04	3.79± 0.86	0.00± 0.00	0.79± 0.46	23.12± 6.17	22	0.32± 0.02
SR1	0.28± 0.22	0.06± 0.04	0.06± 0.04	0.01± 0.01	0.16± 0.12	0.03± 0.02	0.17± 0.03	0.33± 0.39	1.10± 0.47	1	0.00± 0.00
SR2	0.99± 0.48	0.26± 0.14	0.21± 0.08	0.02± 0.01	0.50± 0.22	0.13± 0.06	0.42± 0.08	0.34± 0.31	2.86± 0.75	7	0.00± 0.00
SR3	1.87± 1.00	0.53± 0.25	0.48± 0.15	0.06± 0.01	0.33± 0.20	0.23± 0.18	0.59± 0.11	1.63± 1.11	5.72± 1.81	4	0.01± 0.01
SR4	3.13± 1.31	0.67± 0.28	0.78± 0.22	0.11± 0.03	0.30± 0.34	0.25± 0.13	0.65± 0.12	2.86± 1.58	8.75± 2.28	7	0.02± 0.01
SR5	3.42± 1.49	0.88± 0.39	0.92± 0.23	0.13± 0.02	0.49± 0.18	0.33± 0.12	0.69± 0.13	1.94± 1.20	8.80± 2.18	9	0.05± 0.01
SR6	3.48± 1.51	0.94± 0.49	1.12± 0.28	0.20± 0.03	0.67± 0.33	0.27± 0.09	0.57± 0.11	2.31± 1.44	9.54± 2.45	11	0.07± 0.02
SR7	3.34± 1.47	1.23± 0.57	1.22± 0.29	0.21± 0.03	0.60± 0.14	0.32± 0.07	0.52± 0.10	1.23± 0.86	8.68± 1.99	12	0.08± 0.02
SR8	4.11± 1.76	1.55± 0.83	1.70± 0.38	0.37± 0.05	0.36± 0.13	0.32± 0.14	0.54± 0.10	2.01± 1.25	10.96± 2.47	13	0.21± 0.02
SR9	4.48± 1.96	1.89± 0.91	2.24± 0.52	0.54± 0.07	0.54± 0.07	0.36± 0.11	0.51± 0.10	3.20± 1.82	13.76± 3.12	17	0.35± 0.03
SR10	4.15± 1.66	1.63± 0.82	2.13± 0.48	0.49± 0.07	0.38± 0.20	0.30± 0.10	0.37± 0.07	1.25± 1.43	10.72± 2.54	17	0.52± 0.04
SR11	3.29± 1.44	1.69± 0.69	2.06± 0.47	0.55± 0.07	0.13± 0.13	0.19± 0.06	0.28± 0.05	1.46± 1.21	9.64± 2.27	13	0.79± 0.05
SR12	1.98± 0.90	1.17± 0.47	1.42± 0.33	0.47± 0.06	0.19± 0.03	0.14± 0.10	0.16± 0.03	2.38± 1.37	7.92± 1.87	11	1.03± 0.05
SR13	1.04± 0.51	0.60± 0.34	0.86± 0.22	0.27± 0.03	0.07± 0.03	0.11± 0.02	0.07± 0.01	0.37± 0.42	3.39± 0.94	5	1.09± 0.05
SR14	0.52± 0.26	0.28± 0.13	0.44± 0.12	0.16± 0.02	0.03± 0.02	0.10± 0.02	0.03± 0.00	0.23± 0.26	1.78± 0.47	6	1.14± 0.05
SR15	0.28± 0.16	0.11± 0.06	0.24± 0.08	0.10± 0.02	0.03± 0.01	0.04± 0.01	0.03± 0.01	0.14± 0.16	0.96± 0.29	0	1.12± 0.06
SR16	0.13± 0.09	0.01± 0.02	0.09± 0.03	0.04± 0.00	0.00± 0.00	0.05± 0.00	0.01± 0.00	0.03± 0.06	0.33± 0.14	1	0.91± 0.06
SR17	0.00± 0.00	0.02± 0.02	0.01± 0.00	0.00± 0.00	0.00± 0.00	0.01± 0.00	0.00± 0.00	0.25± 0.32	0.29± 0.32	0	0.26± 0.03

Table 24: Prefit event yields in BDT regions for 2018.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	0.88± 0.31	30.20± 11.24	1.34± 0.35	0.86± 0.13	0.70± 0.42	6.36± 1.55	0.01± 0.00	1.72± 0.66	42.07± 11.64	51	0.49± 0.03
SR1	0.50± 0.38	1.89± 0.10	0.08± 0.06	0.01± 0.01	0.24± 0.20	0.11± 0.17	0.29± 0.05	0.36± 0.36	1.76± 0.89	1	0.00± 0.00
SR2	1.68± 0.87	0.44± 0.28	0.38± 0.20	0.03± 0.02	0.97± 0.49	0.50± 0.19	0.68± 0.13	1.97± 1.04	6.65± 1.59</td		

Table 25: Prefit event yields in BDT regions for 2016+2017+2018.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	1.92± 0.76	60.77±19.35	2.80± 0.61	1.94± 0.25	0.84± 0.35	12.94± 2.43	0.03± 0.00	2.94± 0.97	84.18±19.84	104	1.05± 0.05
SR1	1.02± 0.66	0.28± 0.15	0.19± 0.13	0.02± 0.01	0.40± 0.27	0.18± 0.24	0.66± 0.07	0.86± 0.55	3.61± 1.31	4	0.00± 0.00
SR2	3.45± 1.75	0.83± 0.45	0.80± 0.35	0.07± 0.03	2.19± 0.62	0.85± 0.27	1.52± 0.17	2.71± 1.28	12.43± 2.62	19	0.01± 0.01
SR3	6.83± 3.29	1.82± 0.80	1.72± 0.57	0.25± 0.05	1.65± 0.44	1.14± 0.38	2.26± 0.26	7.90± 2.79	23.58± 5.19	19	0.04± 0.02
SR4	10.53± 4.72	2.47± 1.13	2.83± 0.80	0.39± 0.08	1.66± 0.65	1.18± 0.45	2.51± 0.29	7.86± 2.55	29.44± 6.65	33	0.09± 0.03
SR5	12.38± 5.22	3.04± 1.56	3.45± 0.96	0.58± 0.10	1.49± 0.58	1.23± 0.54	2.59± 0.29	6.64± 2.23	31.41± 7.00	36	0.15± 0.03
SR6	12.60± 5.48	3.71± 1.54	4.12± 1.03	0.74± 0.09	0.97± 0.87	1.35± 0.37	2.17± 0.25	7.76± 2.77	33.41± 7.42	44	0.22± 0.06
SR7	12.35± 5.14	4.50± 2.09	4.43± 1.11	0.81± 0.14	2.21± 0.36	1.02± 0.26	1.86± 0.21	7.27± 2.42	34.45± 6.85	41	0.31± 0.07
SR8	15.32± 6.19	5.66± 2.87	6.10± 1.45	1.21± 0.15	1.61± 0.37	1.21± 0.43	1.95± 0.22	7.16± 2.61	40.22± 8.27	46	0.68± 0.05
SR9	17.01± 7.14	6.82± 3.69	7.95± 1.84	1.87± 0.25	2.14± 0.35	1.56± 0.41	2.02± 0.23	11.11± 4.16	50.48± 10.65	48	1.12± 0.08
SR10	14.93± 6.06	6.20± 3.08	7.53± 1.71	1.87± 0.24	1.46± 0.46	1.14± 0.36	1.40± 0.16	8.86± 3.31	43.40± 8.48	61	1.81± 0.07
SR11	11.98± 5.15	6.11± 2.67	7.27± 1.67	2.03± 0.26	0.95± 0.19	0.92± 0.22	1.07± 0.12	4.79± 2.15	35.12± 7.15	62	2.83± 0.14
SR12	7.43± 3.07	4.18± 1.74	5.24± 1.22	1.54± 0.19	0.70± 0.22	0.89± 0.26	0.61± 0.07	6.41± 2.60	27.00± 5.09	40	3.54± 0.12
SR13	4.20± 1.87	2.40± 1.18	3.15± 0.77	1.01± 0.14	0.26± 0.04	0.52± 0.14	0.28± 0.03	2.31± 0.88	14.13± 3.00	15	4.01± 0.14
SR14	2.06± 1.01	1.32± 0.51	1.78± 0.48	0.55± 0.07	0.11± 0.08	0.30± 0.05	0.18± 0.02	2.21± 1.13	8.50± 1.93	16	3.93± 0.15
SR15	1.38± 0.73	0.68± 0.29	0.96± 0.28	0.37± 0.05	0.21± 0.11	0.17± 0.03	0.12± 0.01	0.44± 0.24	4.33± 1.13	4	3.84± 0.20
SR16	0.43± 0.25	0.20± 0.12	0.32± 0.13	0.15± 0.02	0.07± 0.04	0.10± 0.02	0.03± 0.00	0.84± 0.48	2.13± 0.66	7	3.19± 0.17
SR17	0.03± 0.03	0.07± 0.05	0.03± 0.02	0.02± 0.00	0.00± 0.00	0.02± 0.00	0.00± 0.00	0.25± 0.30	0.44± 0.33	3	1.01± 0.10

Table 26: Postfit event yields in SR+CR regions for 2016+2017+2018.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	2.72± 0.56	75.97±10.55	3.00± 0.67	1.98± 0.24	0.85± 0.37	13.37± 2.24	0.03± 0.00	3.24± 1.03	101.17±10.12	104	0.83± 0.49
CRW	142.07±27.77	34.87± 4.83	37.34± 8.06	7.39± 0.83	14.31± 1.79	9.82± 1.74	18.78± 2.53	66.66±19.89	331.24±18.64	338	3.88± 2.28
SR1	10.22± 2.17	3.29± 0.47	5.16± 1.17	1.23± 0.15	0.64± 0.13	0.62± 0.12	0.49± 0.07	3.99± 1.61	25.64± 2.09	33	1.98± 1.18
SR2	3.61± 0.86	0.69± 0.23	1.63± 0.41	0.39± 0.06	0.15± 0.05	0.17± 0.04	0.10± 0.01	2.40± 1.04	9.15± 1.26	9	1.13± 0.65
SR3	0.72± 0.34	0.20± 0.10	0.65± 0.20	0.13± 0.02	0.02± 0.02	0.04± 0.01	0.04± 0.01	0.21± 0.17	2.01± 0.58	3	0.73± 0.42
SR4	4.02± 0.93	1.58± 0.34	1.94± 0.46	0.61± 0.07	0.44± 0.13	0.31± 0.06	0.32± 0.04	2.13± 0.84	11.34± 1.25	14	1.58± 0.90
SR5	2.20± 0.57	0.61± 0.13	1.01± 0.27	0.32± 0.04	0.10± 0.04	0.26± 0.05	0.10± 0.02	0.44± 0.27	5.03± 0.77	5	1.68± 0.95
SR6	0.80± 0.21	0.18± 0.09	0.43± 0.13	0.10± 0.02	0.13± 0.05	0.05± 0.01	0.02± 0.00	0.56± 0.29	2.29± 0.40	8	1.20± 0.67
SR7	0.31± 0.11	0.12± 0.04	0.20± 0.06	0.04± 0.01	0.01± 0.01	0.03± 0.01	0.01± 0.00	0.00± 0.08	0.71± 0.20	0	0.88± 0.48
SR8	0.70± 0.28	0.28± 0.12	0.42± 0.15	0.22± 0.03	0.05± 0.02	0.11± 0.02	0.08± 0.01	1.44± 0.81	3.31± 0.95	5	2.20± 1.27
SR9	1.46± 0.41	2.24± 0.34	1.58± 0.36	0.31± 0.05	0.14± 0.02	0.16± 0.09	0.00± 0.00	0.94± 0.46	6.84± 0.80	6	0.71± 0.39
SR10	0.33± 0.11	0.63± 0.14	0.56± 0.14	0.11± 0.02	0.01± 0.01	0.05± 0.01	0.00± 0.00	0.42± 0.26	2.10± 0.31	3	0.35± 0.22
SR11	0.18± 0.07	0.32± 0.08	0.20± 0.05	0.04± 0.01	0.01± 0.01	0.02± 0.01	0.00± 0.00	0.60± 0.72	1.38± 0.75	1	0.23± 0.14
SR12	0.22± 0.08	0.61± 0.12	0.42± 0.10	0.09± 0.01	0.04± 0.01	0.06± 0.01	0.00± 0.00	0.59± 0.39	2.03± 0.48	2	0.59± 0.34
SR13	0.29± 0.11	0.36± 0.12	0.31± 0.09	0.07± 0.01	0.02± 0.01	0.04± 0.01	0.00± 0.00	0.00± 0.11	1.09± 0.28	2	0.69± 0.39
SR14	0.16± 0.05	0.23± 0.07	0.18± 0.06	0.04± 0.01	0.00± 0.00	0.03± 0.01	0.00± 0.00	0.23± 0.27	0.87± 0.30	1	0.80± 0.45

Table 27: Postfit event yields in BDT regions for 2016+2017+2018.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	2.50± 0.54	77.26±11.71	3.09± 0.64	1.99± 0.21	0.85± 0.35	13.33± 2.56	0.03± 0.00	3.23± 1.18	102.28±11.58	104	1.11± 0.43
SR1	1.25± 0.49	0.31± 0.10	0.19± 0.10	0.02± 0.01	0.32± 0.25	0.17± 0.19	0.66± 0.07	1.04± 0.64	4	0.00± 0.00	
SR2	4.35± 1.28	0.97± 0.32	0.83± 0.29	0.08± 0.02	2.20± 0.54	0.79± 0.23	1.53± 0.17	3.46± 1.44	14.19± 1.76	19	0.01± 0.01
SR3	8.53± 2.40	2.17± 0.50	1.75± 0.53	0.23± 0.05	1.46± 0.47	1.12± 0.38	2.26± 0.26	8.01± 2.77	25.53± 3.53	19	0.04± 0.03
SR4	13.52± 3.37	3.04± 0.68	3.00± 0.76	0.38± 0.08	1.57± 0.51	1.19± 0.40	2.52± 0.29	8.74± 3.07	33.96± 4.01	33	0.08± 0.05
SR5	15.80± 3.69	3.73± 0.93	3.68± 0.92	0.57± 0.08	1.61± 0.49	1.27± 0.46	2.59± 0.29	7.41± 2.61	36.67± 3.96	36	0.15± 0.07
SR6	16.34± 3.87	4.50± 0.99	4.41± 1.03	0.75± 0.08	1.26± 0.69	1.40± 0.33	2.18± 0.25	8.99± 3.45	39.81± 4.16	44	0.23± 0.12
SR7	15.96± 3.61	5.53± 1.27	4.75± 1.10	0.82± 0.12	2.25± 0.31	1.03± 0.26	1.86± 0.21	8.12± 3.05	40.32± 3.73	41	0.31± 0.16
SR8	19.81± 4.39	6.88± 1.74	6.61± 1.47	1.21± 0.14	1.65± 0.32	1.21± 0.34	1.95± 0.22	7.97± 2.80	47.28± 4.33	46	0.72± 0.28
SR9	22.22± 4.94	8.57± 2.22	8.63± 1.89	1.88± 0.21	2.15± 0.30	1.53± 0.42	2.02± 0.24	11.50± 3.92	58.51± 5.22	48	1.18± 0.46
SR10	19.68± 4.28	7.79± 1.84	8.31± 1.78	1.91± 0.21	1.39± 0.37	1.18± 0.29	1.40± 0.16	10.48± 3.93	52.15± 4.28	61	1.91± 0.74
SR11	15.96± 3.65	7.82± 1.52	8.03± 1.73	2.06± 0.23	0.98± 0.17	0.97± 0.22	1.07± 0.12	6.10± 2.62	43.00± 3.52	62	2.98± 1.19
SR12	9.66± 2.22	5.30± 1.00	5.81± 1.26	1.56± 0.18	0.81± 0.21	0.89± 0.24	0.61± 0.07	7.46± 3.03	32.10± 3.04	40	3.74± 1.41
SR13	5.57± 1.34	3.01± 0.66	3.50± 0.77	1.04± 0.12	0.27± 0.05	0.51± 0.12	0.28± 0.03	2.52± 0.93	16.71± 1.62	15	4.25± 1.63
SR14	2.74± 0.74	1.67± 0.28	1.97± 0.48	0.56± 0.07	0.11± 0.07	0.31± 0.06	0.18± 0.02	2.61± 1.23	10.14± 1.24	16	4.17± 1.59
SR15	1.81± 0.53	0.84± 0.23	1.06± 0.26	0.38± 0.04	0.20± 0.10	0.17± 0.03	0.12± 0.01	0.45± 0.22	5.03± 0.77	4	4.09± 1.55
SR16	0.55± 0.21	0.26± 0.10	0.36± 0.12	0.15± 0.02	0.07± 0.04	0.10± 0.02	0.03± 0.00	0.97± 0.54	2.49± 0.61	7	3.37± 1.25
SR17	0.05± 0.03	0.09± 0.05	0.04± 0.02	0.02± 0.00	0.00± 0.00	0.02± 0.00	0.00± 0.00	0.34± 0.35	0.57± 0.36	3	1.08± 0.42

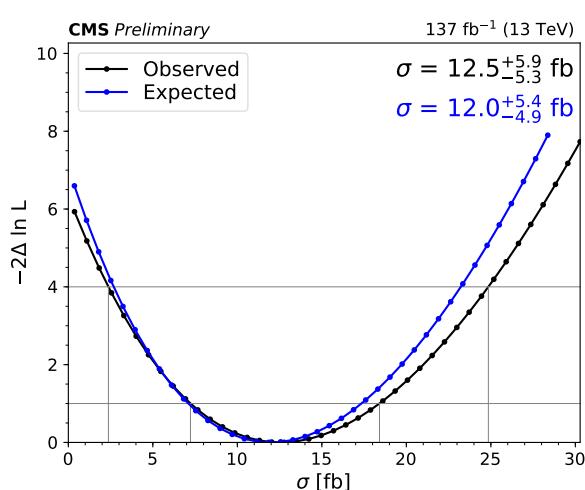


Figure 58: Observed/expected likelihood scans for Run 2 with the BDT analysis

## 962 12 Results: interpretations

### 963 12.1 Type-II 2HDM

#### 964 12.1.1 Introduction

965 Final states with 3 or more top quarks occur in many scenarios of new physics.

966 In in two Higgs doublet models (2HDM) [23, 24] of type-II, the couplings of the CP-even scalar  $h$   
 967 become exactly SM-like in the so-called "alignment limit",  $\sin(\beta - \alpha) \rightarrow 1$ . In such models, the  
 968 couplings of the heavy scalar and pseudoscalar to the SM vector bosons are suppressed, van-  
 969 ishing as  $\cos(\beta - \alpha) \rightarrow 0$ . In such a limit, production is predominantly via gluon-fusion and  
 970 then associated production with either  $b\bar{b}$  or  $t\bar{t}$ . The sensitivity of the direct search for resonant  
 971  $t\bar{t}$  production is significantly reduced due to interference with SM QCD production [25, 26]. As  
 972 a result, at low  $\tan\beta$ , three and four top final states as seen in the diagrams in Figure 59 provide  
 973 a promising window to probe this scenario [27].

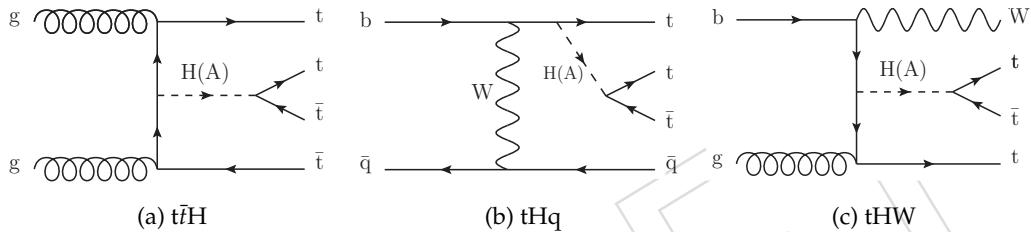


Figure 59: Diagrams for scalar (pseudoscalar) production in association with top quarks. In some scenarios, these heavy bosons will preferentially decay to a pair of top quarks, generating final states with three or four top quarks.

#### 974 12.1.2 Simulation

975 We interpret the results of the SM  $t\bar{t}\bar{t}\bar{t}$  search as limits on the top-associated production of a  
 976 scalar or pseudoscalar with subsequent decay to a pair of top quarks.

977 We generate a one-dimensional grid of points for scalar masses between 350 and 550 GeV in  
 978 20 GeV steps for the three processes shown in Figure 59 at LO with MadGraph in the 5-flavor  
 979 scheme using the NN23LO pdf set. The  $t\bar{t}H$  and  $tWH$  processes were generated with up to one  
 980 additional parton, while the  $tHq$  process was generated with no additional partons due to the  
 981 presence of a light flavor quark in the lowest order ME.

982 For the purpose of interpretation, we use LO cross sections for the production of a heavy scalar  
 983 or pseudoscalar boson using the 2HDMtII\_NLO MadGraph model with the NNPDF30\_lo\_as\_0130  
 984 PDF. Processes involving H and A mediators are generated separately and charged higgses  
 985 ( $H^\pm$ ) are decoupled by setting their mass to 10TeV. We fix  $\sin(\beta - \alpha) = 1$  for the alignment  
 986 condition, and use  $\tan\beta$  as an additional parameter. For consistency with previous results and  
 987 after verifying consistency with the results of [27],  $\tan\beta = 1$  is used for one-dimensional ex-  
 988 clusions as a function of mediator mass. To construct a two-dimensional exclusion plane, we  
 989 calculate cross-sections for the same mediator mass grid with  $\tan\beta$  values ranging from 0.5 to  
 990 3. Note that with the type-II 2HDM, the branching ratio to up-type quarks (e.g., top quark) is  
 991 proportional to  $\frac{1}{\tan\beta}$ .

992 For the pseudoscalar case, we use the same events as the scalar scenario but with cross sections  
 993 determined for the case of a CP-odd scalar with the same mass.

**994 12.1.3 Kinematic comparison of scalar and pseudoscalar diagrams**

995 Figure 60 shows analysis-level quantities for  $t\bar{t}H$  and  $t\bar{t}A$ , with the mediator mass set to 450 GeV  
996 at  $\tan \beta = 1$ . The events were simulated with the 2HDMt II\_NLO MG5 model and reconstructed  
997 with the 2017 MINIAODSIM workflow. In general, acceptance and kinematics are similar be-  
998 tween scalar and pseudoscalar processes, especially the final distribution of yields with the  
999 BDT-based regions, where a KS statistical test yields a p-value of 0.8. For this reason, scalar  
1000 samples can be used for both processes as the production cross-section is the primary differ-  
1001 ence.

DRAFT

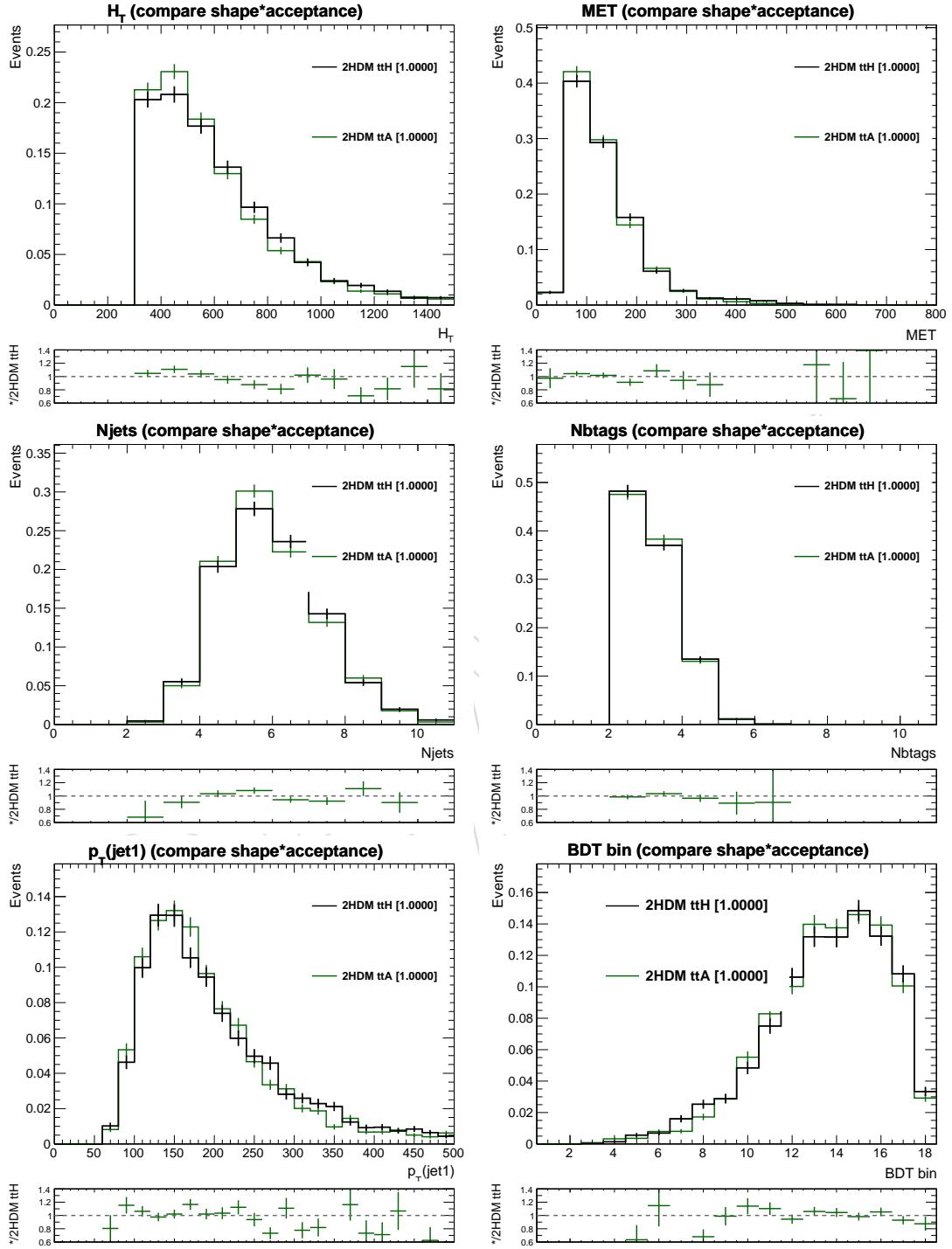


Figure 60: Distributions comparing  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_{\text{b jets}}$ ,  $p_T$  of the leading jet, and the BDT signal region yields for scalar and pseudoscalar processes. The cross-section has been normalized away, so both shapes and acceptance are relevant here.

### 1002 12.1.4 Cross-sections vs $\tan\beta$

1003 Cross sections for the scalar and pseudoscalar processes are plotted in Figure 61 and summarized  
 1004 in Table 28, separated by the value of  $\tan\beta$ . Comparing these limits with the Type-2  
 1005 2HDM cross sections with  $\tan\beta = 1$  in the alignment limit,  $\sin(\beta - \alpha) = 1$ , we exclude scalar  
 1006 (pseudoscalar) masses up to 470 (550) GeV. Alternatively, we consider the simplified model of  
 1007 dark matter defined in Ref. [28], which includes a Dirac fermion dark matter candidate,  $\chi$ , in  
 1008 addition to H/A, and where the couplings of H/A to SM fermions and  $\chi$  are determined by  
 1009 parameters  $g_{\text{SM}}$  and  $g_{\text{DM}}$ , respectively. In this model, exclusions similar to the 2HDM ones are  
 1010 reached by assuming  $g_{\text{SM}} = 1$  and  $g_{\text{DM}} = 1$ , and taking  $m_{\text{H/A}} > 2m_\chi$

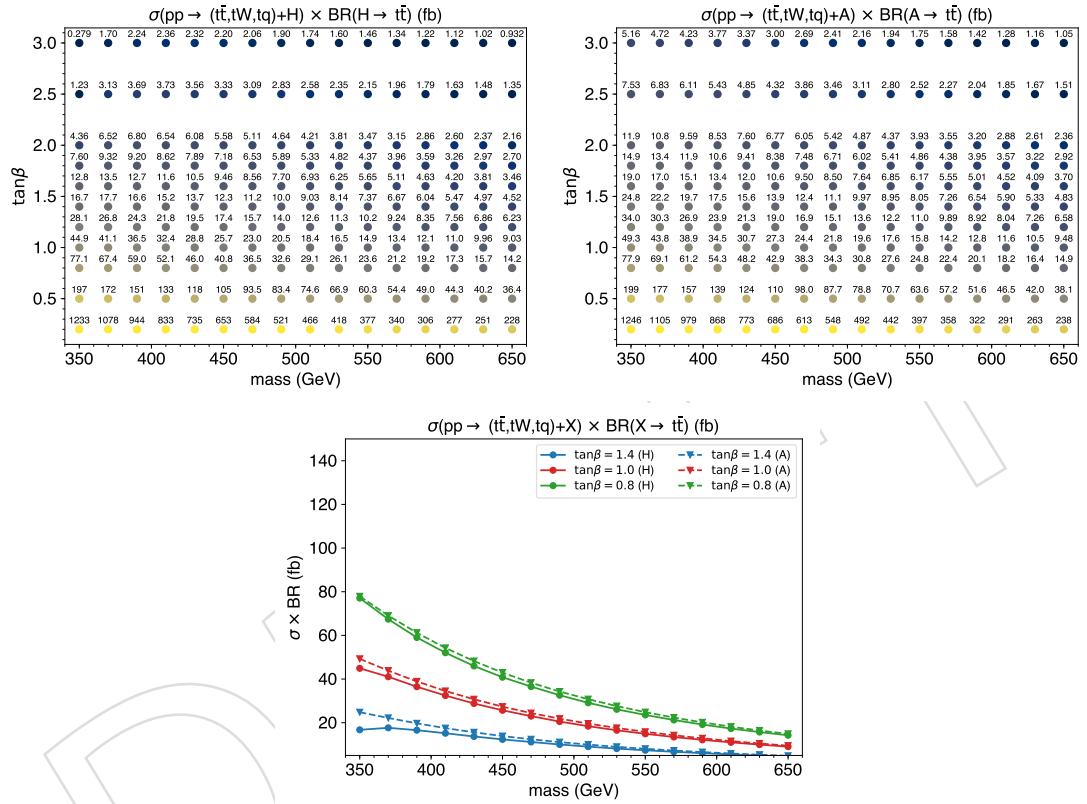


Figure 61: Cross-sections in units of fb in the plane of  $\tan\beta$  vs mediator mass for a heavy scalar boson (top left) and heavy pseudoscalar boson (top right). The bottom plot contains 1-D projections for both mediators for a few values of  $\tan\beta$ .

Table 28: Cross sections for the case of scalar (H) and pseudoproduction (A), assuming a heavy higgs boson with SM-like top quark couplings. All cross-sections are reported in units of fb, and masses in units of GeV. Based on Figure 5 of arXiv:1605.08744, the LO scale uncertainties are close to 40% for  $t\bar{t}H \rightarrow t\bar{t}tt$ ,  $m_H = 400$ . This is consistent with the LO uncertainty on SM  $t\bar{t}tt$ . Single top processes,  $t\bar{t}H \rightarrow t\bar{t}tX$  have a smaller uncertainty (15%) as they have one less power of  $\alpha_S$ .

$\tan \beta$	mass	$\sigma_{t\bar{t}H}$	$\sigma_{t\bar{t}WH}$	$\sigma_{t\bar{t}qH}$	$\sigma_{t\bar{t}A}$	$\sigma_{t\bar{t}WA}$	$\sigma_{t\bar{t}qA}$	$\tan \beta$	mass	$\sigma_{t\bar{t}H}$	$\sigma_{t\bar{t}WH}$	$\sigma_{t\bar{t}qH}$	$\sigma_{t\bar{t}A}$	$\sigma_{t\bar{t}WA}$	$\sigma_{t\bar{t}qA}$
0.5	350	81.4	38.0	77.9	117	36.3	45.7	0.8	350	31.8	14.9	30.4	45.9	14.2	17.8
	370	71.4	34.3	66.7	103	33.0	40.4		370	27.9	13.4	26.1	40.4	12.9	15.8
	390	62.5	31.0	57.5	91.1	29.9	35.7		390	24.4	12.1	22.5	35.6	11.7	13.9
	410	55.2	28.1	49.9	80.1	27.1	31.7		410	21.6	11.0	19.5	31.3	10.6	12.4
	430	48.8	25.4	43.4	70.8	24.7	28.2		430	19.1	9.94	17.0	27.6	9.63	11.0
	450	43.6	23.1	37.9	62.4	22.3	25.1		450	17.0	9.02	14.8	24.4	8.73	9.83
	470	39.1	21.0	33.4	55.4	20.3	22.3		470	15.3	8.22	13.0	21.6	7.94	8.74
	490	35.0	19.1	29.3	49.2	18.5	20.0		490	13.7	7.47	11.4	19.2	7.23	7.82
	510	31.5	17.4	25.7	43.9	16.9	18.0		510	12.3	6.78	10.0	17.2	6.62	7.02
	530	28.3	15.8	22.7	39.1	15.4	16.1		530	11.1	6.17	8.89	15.3	6.03	6.29
	550	25.6	14.4	20.2	35.0	14.1	14.5		550	10.0	5.63	7.90	13.7	5.53	5.65
	570	23.2	13.1	18.0	31.3	12.9	13.0		570	9.07	5.13	7.03	12.2	5.05	5.07
	590	21.0	12.0	16.0	28.1	11.8	11.7		590	8.21	4.70	6.25	11.0	4.61	4.57
	610	19.1	11.0	14.2	25.2	10.8	10.6		610	7.46	4.29	5.56	9.83	4.21	4.13
	630	17.3	10.1	12.8	22.6	9.89	9.54		630	6.77	3.95	4.99	8.83	3.86	3.72
	650	15.8	9.24	11.4	20.4	9.09	8.64		650	6.16	3.61	4.46	7.96	3.55	3.37
1.0	350	18.5	8.66	17.7	29.0	8.98	11.3	1.2	350	11.6	5.40	11.1	20.0	6.20	7.80
	370	17.0	8.19	15.9	25.6	8.18	10.0		370	11.1	5.34	10.4	17.7	5.67	6.92
	390	15.1	7.49	13.9	22.6	7.43	8.85		390	10.0	4.98	9.25	15.7	5.15	6.14
	410	13.4	6.85	12.2	19.9	6.76	7.87		410	9.01	4.60	8.18	13.8	4.68	5.45
	430	11.9	6.24	10.6	17.6	6.13	7.00		430	8.07	4.21	7.18	12.2	4.25	4.86
	450	10.7	5.68	9.32	15.5	5.56	6.26		450	7.25	3.86	6.33	10.8	3.86	4.34
	470	9.61	5.18	8.22	13.8	5.06	5.57		470	6.54	3.53	5.60	9.55	3.51	3.86
	490	8.63	4.70	7.19	12.2	4.61	4.98		490	5.88	3.21	4.92	8.49	3.20	3.46
	510	7.77	4.28	6.35	10.9	4.20	4.47		510	5.31	2.94	4.34	7.58	2.92	3.10
	530	7.00	3.91	5.62	9.74	3.85	4.01		530	4.79	2.67	3.85	6.75	2.66	2.78
	550	6.34	3.56	5.00	8.71	3.51	3.60		550	4.35	2.44	3.43	6.04	2.45	2.50
	570	5.75	3.25	4.44	7.80	3.21	3.23		570	3.94	2.23	3.07	5.41	2.23	2.24
	590	5.20	2.98	3.96	6.99	2.94	2.92		590	3.57	2.05	2.73	4.85	2.05	2.02
	610	4.73	2.72	3.54	6.27	2.69	2.63		610	3.25	1.87	2.43	4.35	1.86	1.83
	630	4.29	2.50	3.16	5.63	2.47	2.38		630	2.95	1.72	2.18	3.91	1.71	1.65
	650	3.91	2.28	2.83	5.08	2.25	2.15		650	2.69	1.58	1.95	3.52	1.57	1.49
1.4	350	6.89	3.23	6.62	14.6	4.52	5.69	1.6	350	5.28	2.48	5.08	11.2	3.47	4.36
	370	7.29	3.52	6.84	13.0	4.15	5.07		370	5.58	2.70	5.24	9.92	3.18	3.89
	390	6.85	3.42	6.32	11.5	3.77	4.49		390	5.24	2.62	4.85	8.77	2.89	3.44
	410	6.27	3.20	5.70	10.1	3.43	4.00		410	4.80	2.45	4.37	7.72	2.63	3.07
	430	5.67	2.97	5.05	8.92	3.12	3.56		430	4.34	2.27	3.88	6.83	2.39	2.73
	450	5.13	2.73	4.48	7.88	2.83	3.17		450	3.93	2.09	3.44	6.03	2.17	2.44
	470	4.66	2.52	3.99	6.99	2.57	2.83		470	3.57	1.93	3.06	5.35	1.97	2.17
	490	4.21	2.30	3.52	6.22	2.34	2.53		490	3.22	1.77	2.71	4.76	1.80	1.94
	510	3.81	2.10	3.12	5.55	2.15	2.28		510	2.91	1.62	2.39	4.25	1.64	1.75
	530	3.44	1.92	2.78	4.95	1.96	2.04		530	2.64	1.48	2.13	3.79	1.50	1.56
	550	3.13	1.76	2.48	4.43	1.79	1.83		550	2.40	1.35	1.90	3.39	1.37	1.41
	570	2.85	1.61	2.21	3.97	1.64	1.65		570	2.18	1.24	1.70	3.04	1.25	1.26
	590	2.58	1.48	1.97	3.55	1.50	1.49		590	1.98	1.14	1.51	2.72	1.15	1.14
	610	2.35	1.36	1.76	3.19	1.37	1.34		610	1.80	1.04	1.35	2.44	1.05	1.03
	630	2.14	1.25	1.58	2.86	1.26	1.21		630	1.64	0.956	1.21	2.19	0.964	0.929
	650	1.95	1.15	1.42	2.58	1.15	1.10		650	1.49	0.877	1.09	1.98	0.884	0.840
1.8	350	3.12	1.47	3.01	8.73	2.72	3.42	2.0	350	1.79	0.845	1.73	7.00	2.19	2.75
	370	3.84	1.86	3.62	7.79	2.51	3.06		370	2.68	1.31	2.53	6.26	2.03	2.47
	390	3.79	1.90	3.51	6.89	2.28	2.72		390	2.80	1.41	2.60	5.55	1.84	2.19
	410	3.55	1.83	3.24	6.07	2.07	2.42		410	2.69	1.39	2.46	4.89	1.68	1.96
	430	3.26	1.72	2.92	5.37	1.88	2.15		430	2.50	1.32	2.25	4.33	1.53	1.74
	450	2.97	1.59	2.61	4.75	1.71	1.92		450	2.31	1.24	2.03	3.83	1.39	1.56
	470	2.72	1.48	2.34	4.21	1.56	1.71		470	2.12	1.16	1.83	3.40	1.26	1.39
	490	2.46	1.35	2.08	3.75	1.42	1.54		490	1.93	1.07	1.64	3.02	1.15	1.24
	510	2.24	1.25	1.84	3.35	1.30	1.38		510	1.76	0.987	1.46	2.70	1.05	1.12
	530	2.03	1.14	1.64	2.98	1.19	1.23		530	1.60	0.904	1.31	2.41	0.963	1.00
	550	1.85	1.05	1.47	2.67	1.08	1.11		550	1.47	0.833	1.17	2.16	0.878	0.900
	570	1.69	0.961	1.31	2.39	0.992	0.998		570	1.34	0.766	1.05	1.93	0.807	0.807
	590	1.53	0.884	1.18	2.14	0.910	0.900		590	1.22	0.706	0.940	1.73	0.736	0.730
	610	1.40	0.811	1.05	1.92	0.830	0.812		610	1.11	0.648	0.841	1.55	0.673	0.658
	630	1.27	0.746	0.945	1.73	0.762	0.734		630	1.02	0.597	0.756	1.40	0.618	0.594
	650	1.16	0.687	0.849	1.56	0.696	0.664		650	0.928	0.551	0.681	1.26	0.567	0.538
2.5	350	0.498	0.240	0.488	4.38	1.40	1.75	3.0	350	0.112	0.056	0.112	2.97	0.978	1.22
	370	1.28	0.633	1.22	3.95	1.30	1.58		370	0.682	0.349	0.668	2.69	0.915	1.11
	390	1.50	0.768	1.41	3.51	1.19	1.41		390	0.899	0.476	0.866	2.40	0.836	0.992
	410	1.52	0.800	1.41	3.10	1.08	1.25		410	0.948	0.514	0.902	2.12	0.764	0.886
	430	1.46	0.784	1.32	2.74	0.985	1.12		430	0.933	0.514	0			

## 1011 12.1.5 Exclusions

With the Run2 BDT analysis and the previous model assumptions, heavy scalar (pseudoscalar) bosons are excluded as shown in Figure 62, as a function of mediator mass for  $\tan \beta = 1$ , and in Figure 63, as a function of mediator mass and  $\tan \beta$ , considering one new particle at a time and also including a scenario with  $m_H = m_A$  inspired by a special case of Type-2 2HDM, the hMSSM [? ]. For this interpretation the SM  $t\bar{t}t\bar{t}$  process is treated as a background and assigned its SM cross section and uncertainty  $12.0^{+2.2}_{-2.5}$  fb [6].

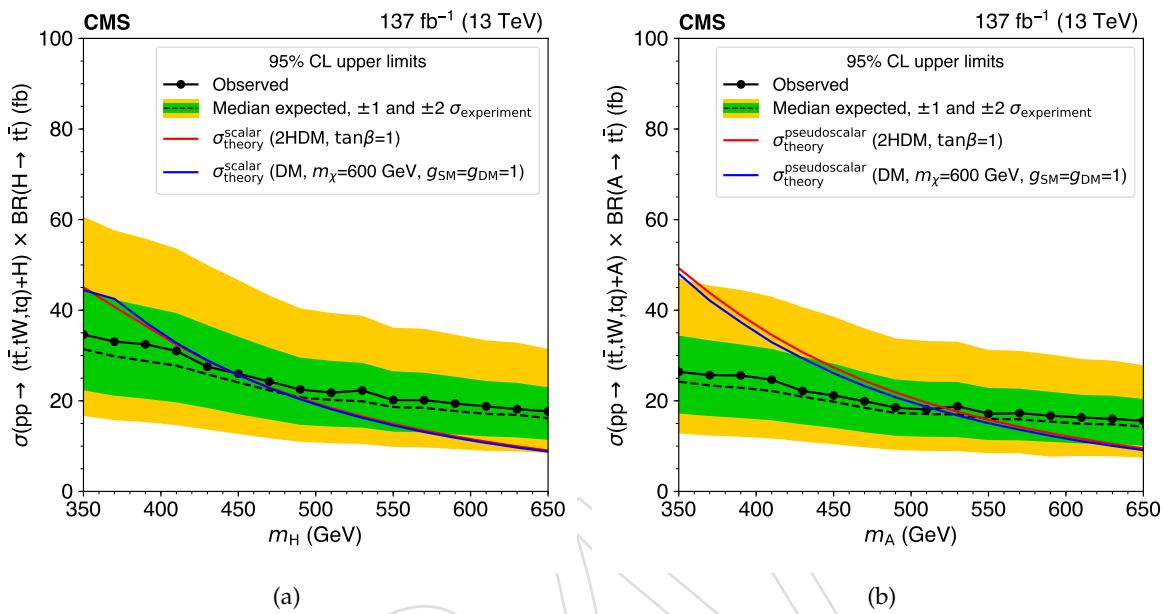


Figure 62: Observed and expected limits on heavy scalar (left) and pseudoscalar (right) processes as a function of the (pseudo)scalar mass using FullSim samples. The parameter  $\tan \beta$  is assumed to be 1 here. The DM model is described in Section 12.5.

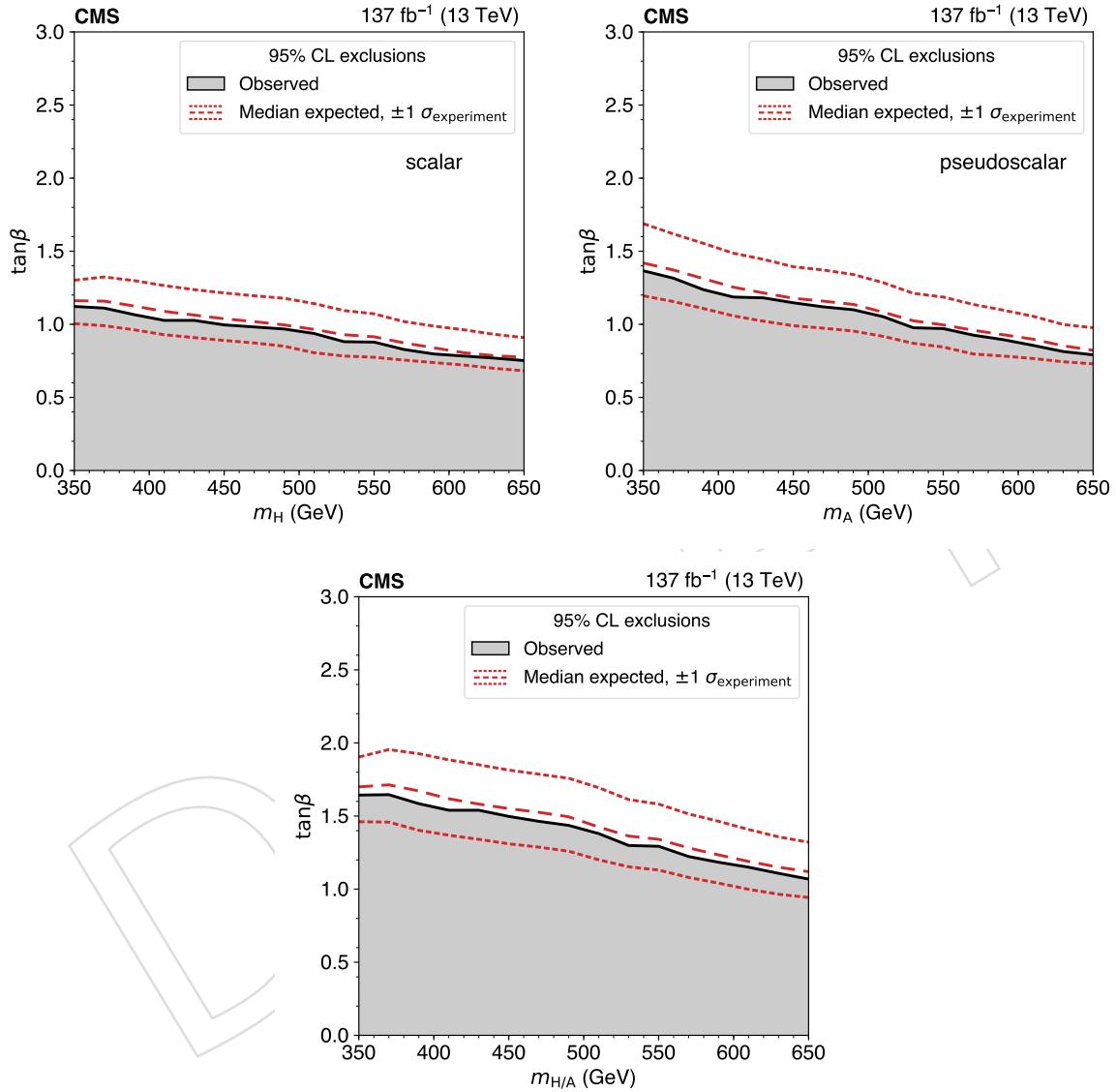


Figure 63: Two-dimensional observed and expected exclusions of heavy scalar only (top left), pseudoscalar only (top right), and both scalar+pseudoscalar (both) as a function of the mediator mass and  $\tan\beta$ . In each case, other 2HDM particles (except SM higgs) are decoupled.

1018 **12.2 Top Yukawa coupling**

1019 **12.2.1 Introduction**

1020 In the SM there are contributions to  $pp \rightarrow t\bar{t}t\bar{t}$  from diagrams with virtual Higgs bosons, see  
 1021 for example Figure 64. The amplitude corresponding to these diagrams is proportional to the  
 1022 square of the top Yukawa coupling.

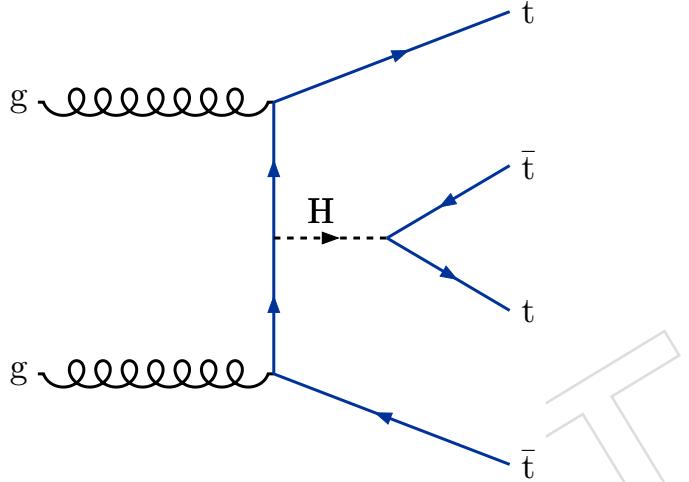


Figure 64: One of the Feynman diagrams for  $t\bar{t}t\bar{t}$  including a virtual Higgs.

1023 Using the notation of Reference [29] the  $t\bar{t}t\bar{t}$  cross-section can be written as

$$\sigma(t\bar{t}t\bar{t}) = \sigma^{SM}(t\bar{t}t\bar{t})_{g+Z/\gamma} + k_t^4 \sigma^{SM}(t\bar{t}t\bar{t})_H + k_t^2 \sigma_{int}^{SM} \quad (6)$$

1024 where  $k_t \equiv y_t/y_t^{SM}$ ,  $y_t$  is the top Yukawa coupling, and  $y_t^{SM}$  is its value in the SM. In equation 6  
 1025 the first term on the right hand side corresponds to the SM contribution to the cross section  
 1026 from diagrams with virtual gluons or  $Z/\gamma$ , the second term is the contribution from diagrams  
 1027 with virtual  $H$  bosons, and the third term is the interference between the two. Therefore, given  
 1028 a theoretical calculation and a measurement of  $\sigma(t\bar{t}t\bar{t})$ , one can put constraints on  $|y_t/y_t^{SM}|$ .

1029 The authors of Reference [29] have calculated the cross-section terms at LO, and have provided  
 1030 us privately with the uncertainties under variations of the factorization and renormalizations  
 1031 scales. These are given in Table 29.

	$[\mu/2, \mu, 2\mu]$
$\sigma^{SM}(t\bar{t}t\bar{t})_{g+Z/\gamma}$	[14.104, 9.997, 6.378] fb
$\sigma^{SM}(t\bar{t}t\bar{t})_H$	[1.625, 1.167, 0.7655] fb
$\sigma_{int}^{SM}$	[-2.152, -1.547, -0.999] fb

Table 29: LO calculation of the terms in equation 6 from Reference [29]. The scale variations are private communications from the authors.

1032 We have investigated the possibility of a full NLO treatment of the interpretation, and dis-  
 1033 cussed with the authors of the NLO calculation [6], but Madgraph could not provide what we  
 1034 need yet. So, we have decided to continue using the LO calculation for this interpretation. The  
 1035 authors of Reference [29] have also argued that it is appropriate to apply the overall NLO/LO  
 1036  $k$ -factor of 1.27 calculated at 14 TeV for the **total**  $\sigma(t\bar{t}t\bar{t})$  cross-section [30] to the individual

1037 components in equation 6. This would then result in an NLO cross-section of  $12.2^{+5.0}_{-4.4}$  fb. This  
 1038 is in agreement with the NLO calculation of  $11.97^{+2.15}_{-2.51}$  fb [6].<sup>1</sup> The authors of Reference [29]  
 1039 then go on to extract a limit on  $|y_t/y_t^{SM}|$  based on their calculation and the then-available ex-  
 1040 perimental limit on  $\sigma(t\bar{t}t\bar{t})$ . Following this procedure in the 2016 result, in Fig. 65 (left) we  
 1041 show the measurement of the cross section and its upper limit, as well as its SM prediction as  
 1042 a function of the absolute value of the ratio of the top quark Yukawa to its SM value ( $|y_t/y_t^{SM}|$ ).  
 1043 The central (upper,lower) value of the theoretical cross section band resulted in a 95% CL limit  
 1044  $|y_t/y_t^{SM}| < 2.27$  (2.03,2.56) (in the 2016 result).

1045 Noting that ttH is a non-negligible background, showing a flat line for the observed cross sec-  
 1046 tion of  $t\bar{t}t\bar{t}$  (and observed upper limit) as a function of  $\kappa_t$  is not completely correct, as it neglects  
 1047 the  $\kappa_t^2$  scaling of ttH, since  $\sigma(t\bar{t}H) \propto (y_t/y_t^{SM})^2 / (\Gamma_H/\Gamma_H^{SM})$ , where  $\Gamma_H$  is the total width of the  
 1048 Higgs. The limit plot therefore is also shown, on the right, for the case of the ttH background is  
 1049 scaled by  $(y_t/y_t^{SM})^2$  for each value of  $|y_t/y_t^{SM}|$ . As the ttH background becomes larger, fewer  
 1050  $t\bar{t}t\bar{t}$  events are allowed by the fit, and the observed  $t\bar{t}t\bar{t}$  upper limit and the measured  $t\bar{t}t\bar{t}$   
 1051 cross sections becomes smaller. With this prescription, the central (upper,lower) value of the  
 1052 theoretical cross section band results in a 95% CL limit  $|y_t/y_t^{SM}| < 2.10^{+0.22}_{-0.27}$  (in the 2016 result).

1053 For the full Run2 BDT analysis, this procedure resulted in a 95% CL limit of the central, upper,  
 1054 and lower values of the theoretical cross section provide respective 95% CL limits for  $|y_t/y_t^{SM}|$   
 1055  $< 1.7$ ,  $< 1.4$ , and  $< 2.0$ , as shown in Figure 66.

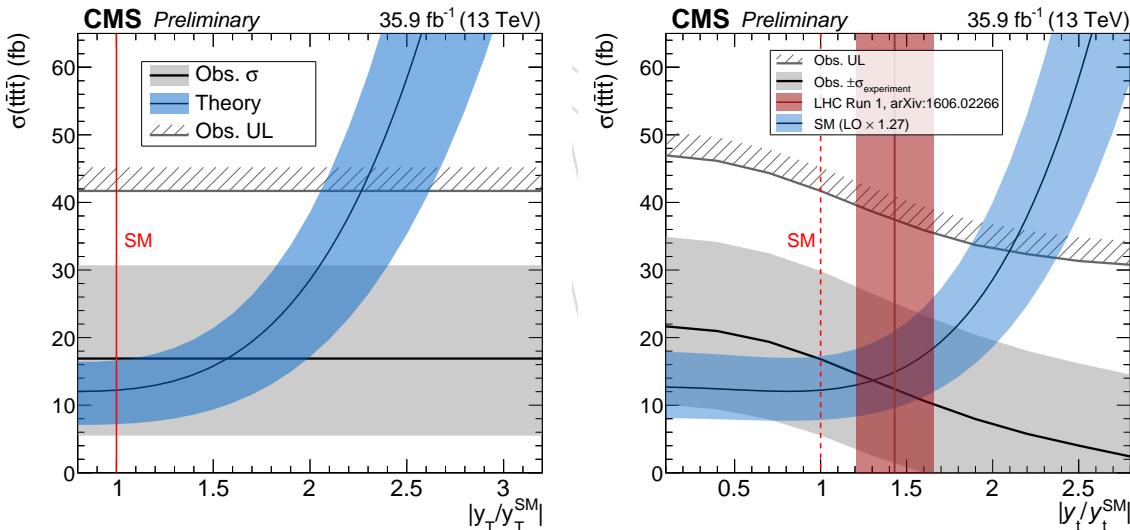


Figure 65: (2016-only results, the next plot shows the full run 2 results). The expected  $t\bar{t}t\bar{t}$  cross-section,  $\sigma(t\bar{t}t\bar{t})$ , as a function of  $|y_T/y_T^{SM}|$ , the absolute value of the ratio of the top quark Yukawa coupling and its SM value (diagonal band), compared with the measured  $\sigma(t\bar{t}t\bar{t})$  (horizontal band), and its 95% CL upper limit (horizontal line). The right plot includes the scaling of the ttH background as a function of  $y_T$  as discussed in the text.

### 1056 12.2.2 Alternative statistical treatment

1057 Note that this section is from 2016, but is left here as a reference

1058 An alternative analysis, still based on the LO cross-section (and its uncertainty) from Refer-  
 1059 ence [29], scaled up by the k-factor of 1.27, consists of interpreting the experimental likelihood

<sup>1</sup>Note that the uncertainties in the full NLO calculation are smaller, but this is to be expected.

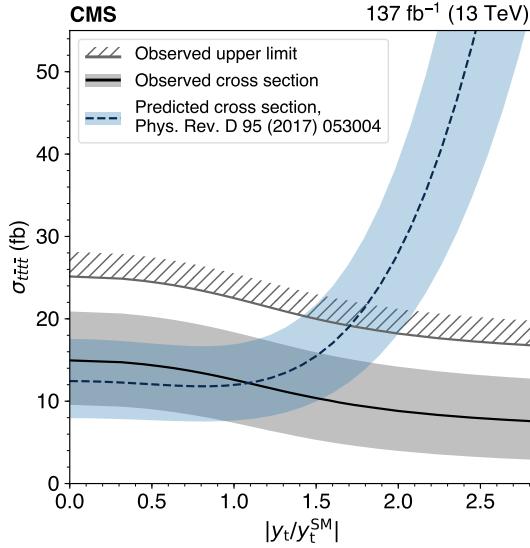


Figure 66: The expected  $t\bar{t}t\bar{t}$  cross-section,  $\sigma(t\bar{t}t\bar{t})$ , as a function of  $|y_T/y_T^{SM}|$ , the absolute value of the ratio of the top quark Yukawa coupling and its SM value (diagonal band), compared with the measured  $\sigma(t\bar{t}t\bar{t})$  (horizontal band), and its 95% CL upper limit (horizontal line)

as the posterior pdf for the cross-section.<sup>2</sup> We then construct (correlated) pdfs for the three terms on the right side of equation 6, and extract a pdf for  $|y_T/y_T^{SM}|$  by propagating uncertainties using equation 6 and a toy MC method. We do this in two ways, by either taking the correlated pdfs as bifurcated gaussians (truncated to not allow a change of sign), or as “flat” within the limits of Table 29. The 95% CL limit would then be  $|y_T/y_T^{SM}| < 2.45$  (bifurcated gaussians pdfs) or  $|y_T/y_T^{SM}| < 2.25$  (flat pdfs). These results are to be compared with the “constant ttH” interpretation of Figure 65 (left), which resulted in a 95% CL limits of  $|y_T/y_T^{SM}| < 2.27$  (2.03,2.56).

We prefer to present the results according to Figure 65. This is in the same spirit as what is customary in the SUSY group, where the experimental limits are clearly separated from the theoretical uncertainties in the interpretation (but maybe the TOP group does it differently). In addition, we are reluctant to give too much weight to an error analysis that includes the systematic uncertainty from Reference [29], which is somewhat of a hack (a LO calculation with the pieces of the LO calculation scaled by the same k-factor as obtained by another group for the full calculation at a slightly different energy...). Incidentally, the limit on  $|y_T/y_T^{SM}|$  quoted in Reference [29] ignores the uncertainty on the theoretical calculation. In any case, we are open to suggestions.

### 12.2.3 Kinematic dependence on top yukawa coupling

Since MC is taken with a nominal top yukawa coupling and only the cross-section is scaled in the scan above, it is worth verifying that there is no significant kinematic dependence on the top yukawa coupling value. This is checked with LO MG5 (MG5\_aMC\_v2\_6\_3\_2) with default parameters, including a dynamical scale choice) using the following proc card.

```
1082 set default_unset_coupleings 99
```

<sup>2</sup>This would correspond to using a Bayesian method with a flat prior in the cross-section. It is not entirely kosher since the likelihood has been profiled with respect to the nuisances, while Bayesian approaches usually require marginalization.

```

1083 set group_subprocesses Auto
1084 set ignore_six_quark_processes False
1085 set loop_optimized_output True
1086 set loop_color_flows False
1087 set gauge unitary
1088 set complex_mass_scheme False
1089 set max_npoint_for_channel 0
1090 set nb_core 4
1091
1092
1093 import model sm
1094 define p = g u c d s u~ c~ d~ s~
1095 define p = p b b~
1096 generate p p > t t~ t t~ QED=99
1097
1098
1099 output ftlo_ytscan
1100 launch
1101
1102 # >>> import numpy as np
1103 # >>> x = np.arange(0.4,2.2,0.1)
1104 # >>> print ",".join(map(lambda y: str(round(y,1)),x))
1105 # paste output in scan brackets below
1106
1107 set param_card yukawa 6 scan:[69.2,...,363.3]

```

1108 The LHE for  $|y_T/y_T^{SM}|$  between 0.4 and 2.2 in steps of 0.1 was carried through the 2016 MINIAOD-SIM workflow with a slightly modified Pythia fragment which imposes a dilepton (or more) filter for better statistical uncertainties. The relevant part of the fragment is

```

1111 'JetMatching:setMad = off',
1112 'JetMatching:scheme = 1',
1113 'JetMatching:merge = on',
1114 'JetMatching:jetAlgorithm = 2',
1115 'JetMatching:etaJetMax = 5.',
1116 'JetMatching:coneRadius = 1.',
1117 'JetMatching:slowJetPower = 1',
1118 'JetMatching:qCut = 59.,
1119 'JetMatching:nQmatch = 5', #4 corresponds to 4-flavour scheme (no matching of b-
1120 'JetMatching:nJetMax = 0', #number of partons in born matrix element for highest
1121 'JetMatching:doShowerKt = off', #off for MLM matching, turn on for shower-kT mat
1122 '6:m0 = 172.5',
1123 '24:mMin = 0.1',
1124 '23:mMin = 0.1',
1125 'ResonanceDecayFilter:filter = on',
1126 'ResonanceDecayFilter:exclusive = off', #off: require at least the specified num
1127 'ResonanceDecayFilter:eMuAsEquivalent = off', #on: treat electrons and muons as
1128 'ResonanceDecayFilter:eMuTauAsEquivalent = on', #on: treat electrons, muons , an
1129 'ResonanceDecayFilter:allNuAsEquivalent = on', #on: treat all three neutrino fla
1130 'ResonanceDecayFilter:daughters = 11,11',

```

1131 'Check:abortIfVeto = on',

1132 To validate the MadGraph setup, cross-sections for points along the scan are compared with  
 1133 Equation 6 (using values from Table 29) and they are found to agree within 2%, as shown in  
 1134 Figure 67

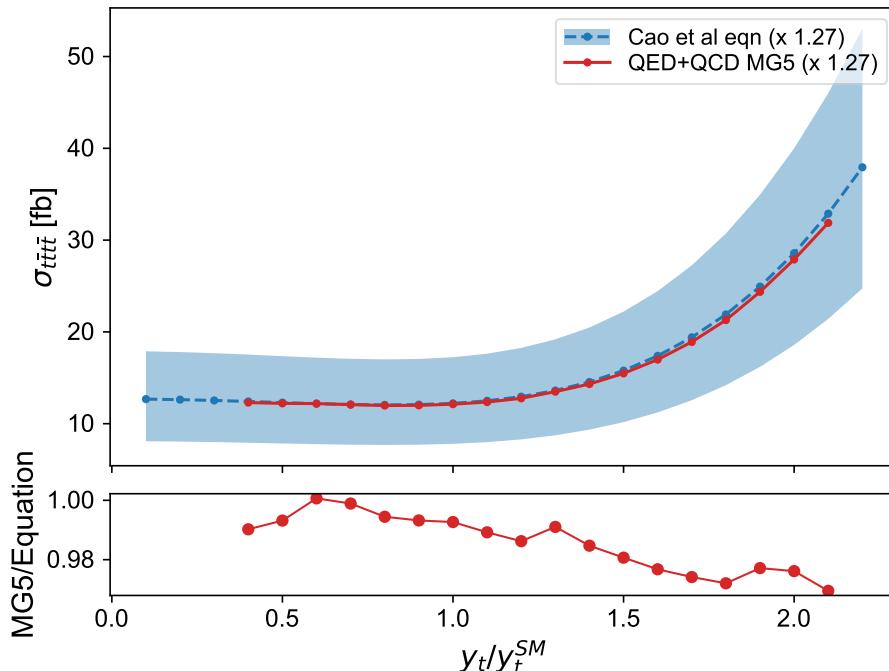


Figure 67: Calculated cross-section from Equation 6 and from MG5 for varying  $|y_T/y_T^{SM}|$  values

1135 Finally, plots of various kinematic quantities are shown in Figure 68 for the analysis baseline  
 1136 selection, also allowing opposite-sign events in addition to same-sign events for augmented  
 1137 statistics. There is no visible significant difference in the kinematic plots for different values of  
 1138  $|y_T/y_T^{SM}|$ .

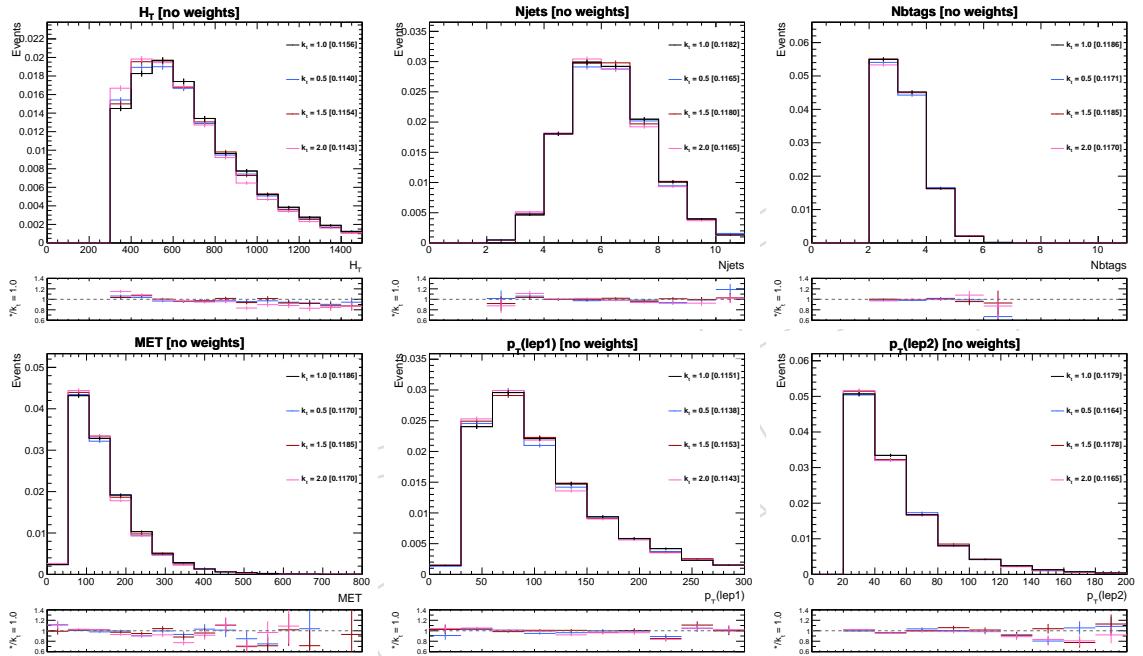


Figure 68: Kinematic quantities (left to right, top to bottom:  $H_T$ ,  $N_{\text{jets}}$ ,  $N_{\text{btags}}$ ,  $\text{MET}$ ,  $p_T(\text{lep1})$ ,  $p_T(\text{lep2})$ ) for the baseline selection with same-sign and opposite-sign dileptons for  $|y_T/y_T^{SM}|$  of 0.5, 1.0, 1.5, 2.0. All histograms are normalized to the same cross-section.

1139 **12.3 Off-shell mediators decaying to top quark pairs**

1140 **12.3.1 Introduction**

The production of  $t\bar{t}t\bar{t}$  may also be influenced by a neutral scalar mediator ( $\phi$ ) or neutral vector mediator ( $Z'$ ) which couple to top quarks and have masses less than twice the mass of the top quark, distinguishing them from similar processes within the 2HDM framework, for example. The off-shell contributions to the SM  $t\bar{t}t\bar{t}$  production can be large, as shown in Ref. [31]. For a large range of masses, the authors have shown that kinematics are identical when considering these additional processes, so that the total  $t\bar{t}t\bar{t}$  cross-section is subject to a simple rescaling. We consider coupling terms in the lagrangian of the form

$$\mathcal{L}_{Z'} = -g_{tZ'} \bar{t}_R Z' t_R \quad \mathcal{L}_\phi = -g_{t\phi} \bar{t}_L \phi t_R$$

1141 and calculate leading order cross-sections for the process

1142 `generate p p > t~ t t~ t QED=2`

1143 with MadGraph UFO models provided by the authors of Ref. [31], which have been copied to  
 1144 <https://github.com/aminnj/FTInterpretations/tree/master/models>. We also  
 1145 considered single top processes, similar to those for the 2HDM exclusions, but found their  
 1146 cross-section to be small compared to  $t\bar{t}$ -associated production, so they are not included.

1147 Due to the approximate independence of kinematics on the coupling strength and mediator  
 1148 mass, we are able to use the upper limit result from the nominal analysis to place constraints  
 1149 on couplings  $g_{tZ'}$  and  $g_{t\phi}$  as a function of masses  $m_{Z'}$  and  $m_\phi$ , respectively, without the use  
 1150 of dedicated signal samples. The nominal analysis uses a NLO  $t\bar{t}t\bar{t}$  sample, so to justify this  
 1151 procedure, we first show that the nominal NLO  $t\bar{t}t\bar{t}$  sample indeed has good shape agreement  
 1152 with a LO  $t\bar{t}t\bar{t}$  generated with these models (setting coupling strengths to 0) in Figure 69. Next,  
 1153 we show various couplings and mass points near the exclusion boundary compared to the LO  
 1154 SM  $t\bar{t}t\bar{t}$  sample in Figure 70 for the  $Z'$  mediator, and Figure 71 for the  $\phi$  mediator. Based on  
 1155 the level of agreement for both mediator types, we include a 10% additional normalization  
 1156 uncertainty on the SM  $t\bar{t}t\bar{t}$  signal to conservatively cover minor acceptance differences.

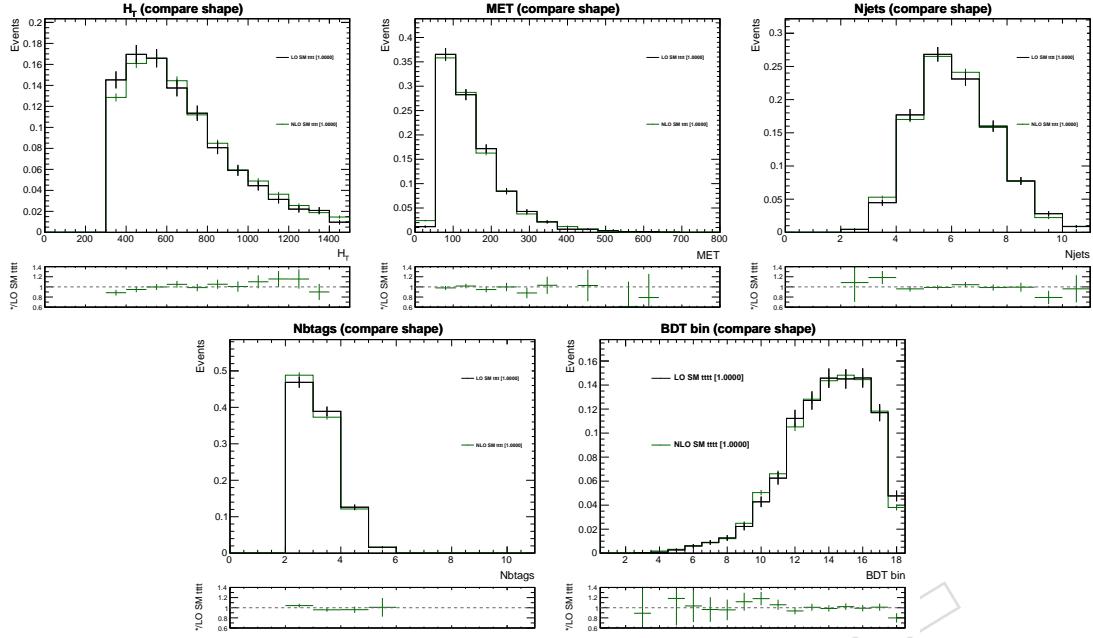


Figure 69: Distributions comparing  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_{\text{b jets}}$ , and the BDT signal region yields for the nominal NLO and LO  $t\bar{t}t\bar{t}$  samples. Only shapes are relevant here.

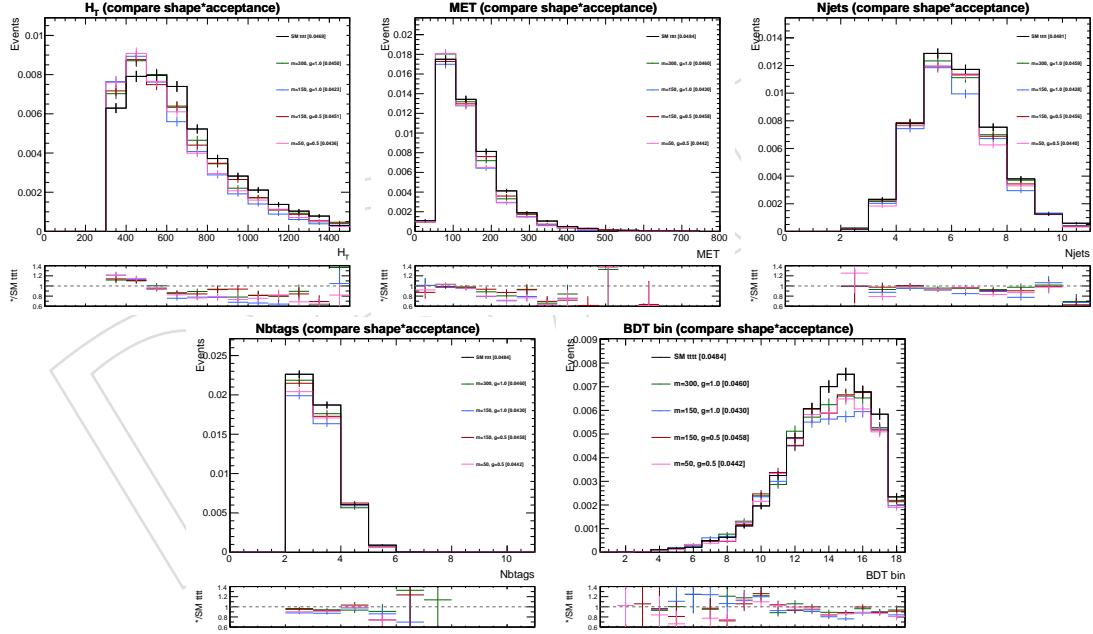


Figure 70: Distributions comparing  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_{\text{b jets}}$ , and the BDT signal region yields for the LO  $t\bar{t}t\bar{t}$  samples and various vector mediator mass points. Both shape and acceptance are relevant here.

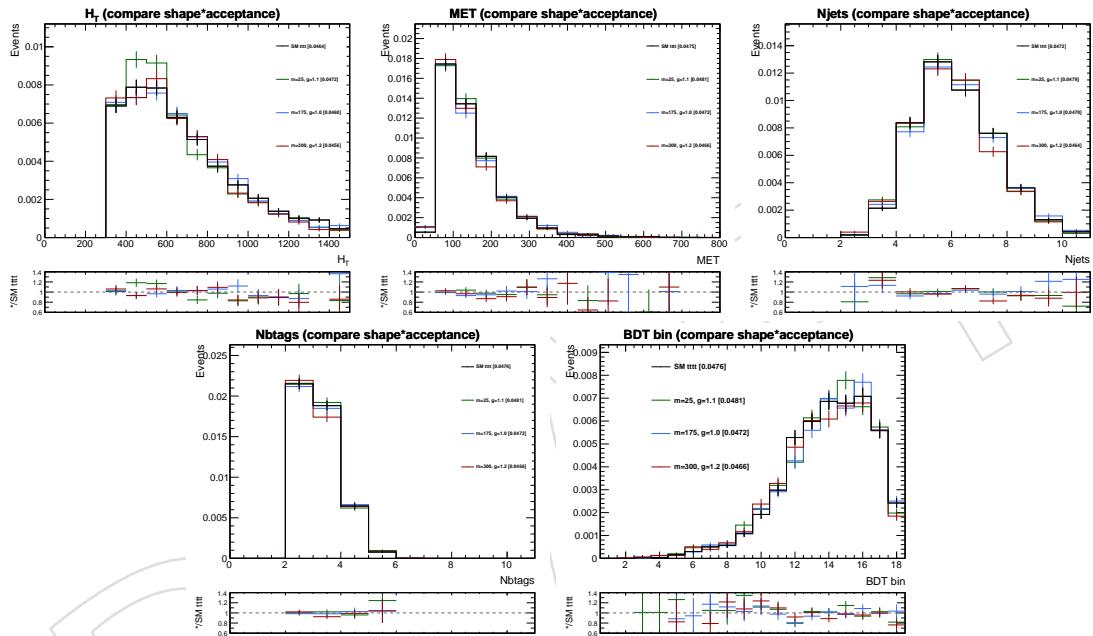


Figure 71: Distributions comparing  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_{\text{b jets}}$ , and the BDT signal region yields for the LO  $t\bar{t}t\bar{t}$  samples and various scalar mediator mass points. Both shape and acceptance are relevant here.

1157 The BDT observed upper limit on  $t\bar{t}t\bar{t}$  production, including the extra 10% uncertainty previ-  
 1158 ously motivated, is  $23 \text{ fb}^{-1}$ . Taking the ratio with the SM cross-section gives approximately 1.9.  
 1159 This is represented as the horizontal dashed line in Figure 72, which includes curves for cross-  
 1160 sections (normalized to SM) for various mediator masses as a function of coupling strengths.  
 1161 Intersections between the cross-section curves and the horizontal dashed line are marked with  
 1162 vertical dashed lines. To obtain smoother values, a quadratic interpolation between generated  
 1163 points is used. These intersections form the exclusion boundary as shown in Figure 73.

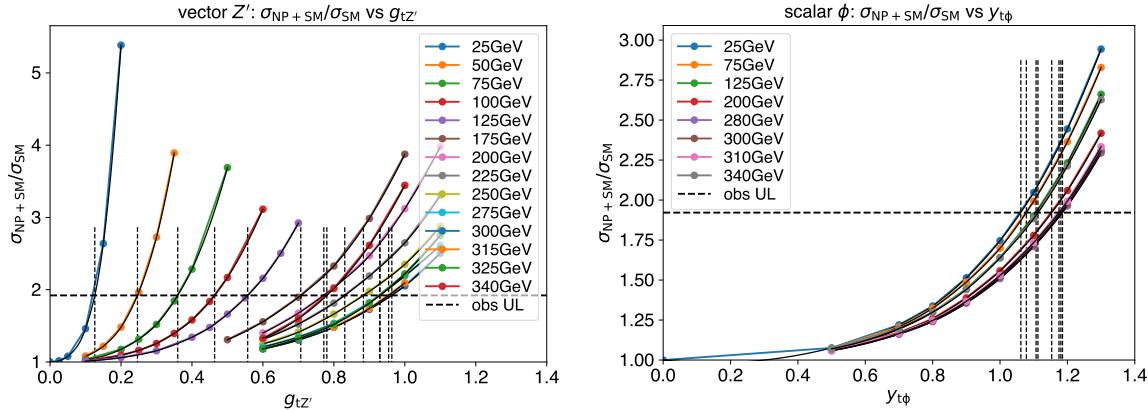


Figure 72: Cross-section (normalized to SM) as a function of coupling values with varying mediator masses for vector (left) and scalar (right) mediators. The observed analysis upper limit is shown as a horizontal dashed line.

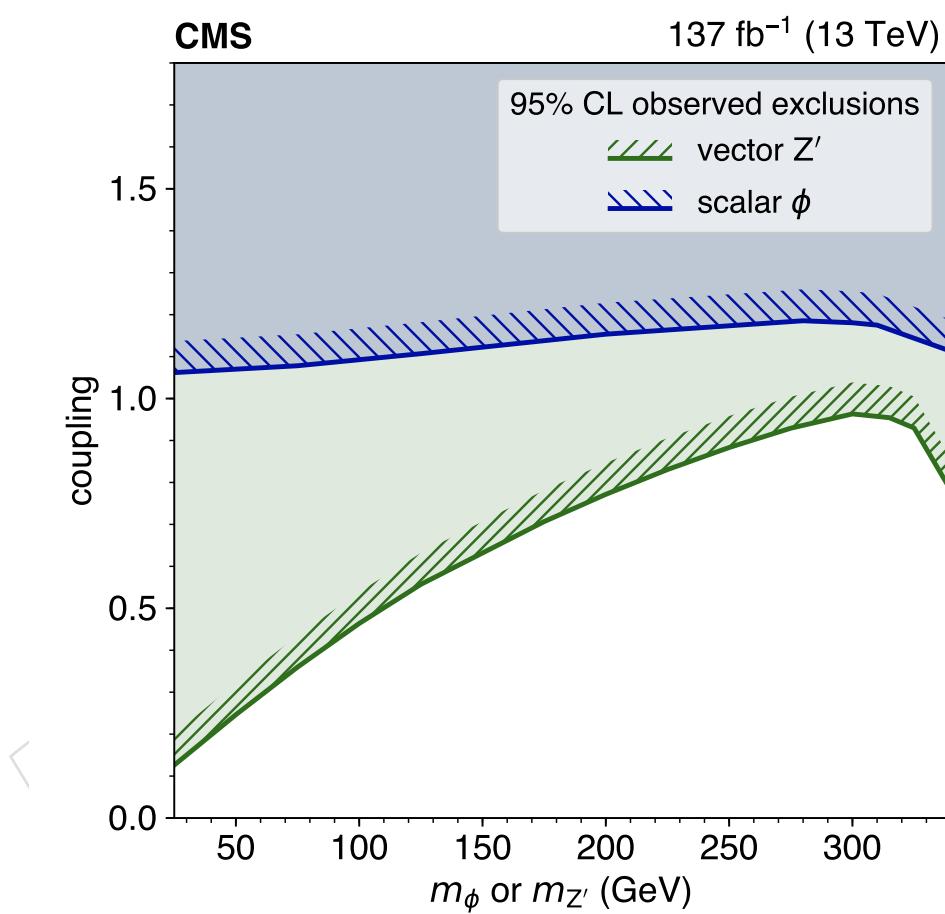


Figure 73: Observed 95% CL exclusions in the plane of coupling-mass for off-shell vector and scalar mediators.

## 1164 12.4 Oblique Higgs parameter

### 1165 12.4.1 Cross-section calculation

1166 In a universal effective field theory framework, the Higgs oblique parameter  $\hat{H}$ , defined as the  
 1167 Wilson coefficient of the dimension-6 operator modifying the Higgs boson propagator, can re-  
 1168 sult in deviations of the SM  $t\bar{t}t\bar{t}$  cross-section, as shown in Ref. [32]. These (off-shell) deviations  
 1169 can be constrained to a level which is competitive with constraints from on-shell processes.

The two main characteristic effects of this oblique parameter are an additional contact term in the SM Higgs boson propagator

$$P_h(p^2) \approx \frac{i}{p^2 - m_h^2} - \frac{i\hat{H}}{m_h^2},$$

and a rescaling of the fermionic higgs couplings

$$\kappa_f = 1 - \hat{H}.$$

1170 Using the latest combined fits of ATLAS for the (on-shell) fermionic couplings, with  $80\text{ fb}^{-1}$  of  
 1171  $13\text{ TeV}$  data, the authors of Ref. [32] find a constraint on the oblique parameter of  $\hat{H} < 0.16$  at  
 1172 95% CL.

The authors also calculate that the cross-section of (off-shell)  $t\bar{t}t\bar{t}$  is subject to a fractional modification (with respect to the SM cross-section) at  $14\text{ TeV}$ , given by,

$$\frac{\sigma_{\hat{H}+\text{SM}}}{\sigma_{\text{SM}}} = 1 + 0.03 \left( \frac{\hat{H}}{0.04} \right) + 0.15 \left( \frac{\hat{H}}{0.04} \right)^2.$$

1173 For an oblique parameter value of 0.1, the formula predicts a doubling of the SM cross-section  
 1174 of  $t\bar{t}t\bar{t}$  with MG 2.6.1 and 2.6.5

The SM model within MadGraph was modified to take into account the **extra term in the propagator**, as well as the **rescaling of the top-yukawa coupling**, by changing only the numerator of the scalar propagator affecting the Higgs boson:

$$\frac{i}{p^2 - m^2 + im\Gamma} \rightarrow \frac{(1 - \hat{H})^2 \left( 1 - \frac{\hat{H}}{m^2} (p^2 - m^2) \right)}{p^2 - m^2 + im\Gamma}.$$

1175 Note that the complex  $i$  in the numerator of the default propagator is removed due to an inter-  
 1176 internal inconsistency found in the latest versions of MadGraph, as verified by the authors. Explic-  
 1177 itely, we took the SM model UFO file and modified the numerator of S in file propagators.py  
 1178 to be

```
1179     "(1-hhat)*(1-hhat)*(1-(hhat/(Mass(id)*Mass(id))))" +
1180     "* (P('mu', id) * P('mu', id) - Mass(id) * Mass(id)))"
```

1181 instead of " $i$ ", and added the line

```
1182     propagator = Prop.S,
```

1183 to the SM Higgs particle definition in file `particles.py`. Finally, for convenience of scanning  
 1184 the parameter  $\hat{H}$ , we added

```
1185 hhat = Parameter(  

1186     name = 'hhat',  

1187     nature = 'external',  

1188     type = 'real',  

1189     value = 0.,  

1190     texname = '\\text{hhat}',  

1191     lhablock = 'PROP',  

1192     lhacode = [ 1 ],  

1193 )
```

1194 to file `parameters.py`. The final model can be found at [https://github.com/aminnj/FTInterpretations/tree/master/models/Oblique\\_UFO](https://github.com/aminnj/FTInterpretations/tree/master/models/Oblique_UFO).

### 1196 12.4.2 Private generation details

1197 We privately generated five different values of  $\hat{H}$  with the above modifications to the SM model  
 1198 in MadGraph 2.6.5 with the nn23lo1 PDF and default (dynamic) scale choices at leading order.  
 1199 Each parameter point consists of three datasets of 50k events each with LHE→MINIAODSIM  
 1200 configurations matching the RunIISummer16MiniAODv3, RunIIFall17MiniAODv2, and  
 1201 RunIIAutumn18MiniAOD campaigns. To increase the baseline selection efficiency, we used  
 1202 the dilepton Pythia filter described in Section 12.2, which requires events to have at least 2 gen-  
 1203 erator leptons (of any flavor) with a filter efficiency of 0.3944. The cross-sections (times filter  
 efficincies) for these parameter points are tabulated below.

$\hat{H}$	$\sigma \times \epsilon_{\text{filter}}$ (fb)
0.0	4.721
0.04	5.411
0.08	6.752
0.12	8.456
0.16	10.235

1204

### 1205 12.4.3 Comparison of kinematics/acceptance

1206 Figure 74 shows analysis-level quantities for five different values of  $\hat{H}$ , one of which is analo-  
 1207 gous to SM  $t\bar{t}t\bar{t}$  ( $\hat{H} = 0$ ). In the last BDT bin, the acceptance for the two values of non-zero  $\hat{H}$  are  
 1208 within approximately 10% of the SM. In general, acceptance for those non-zero points is 10-20%  
 1209 higher than the SM in the higher BDT bins, with visible trends in jet/lepton momentum, and  
 1210  $H_T$ . The nominal analysis upper limit cannot be directly used to place a constraint on  $\hat{H}$  as ac-  
 1211 ceptance increases for higher values of  $\hat{H}$ . However, the next section outlines the method of ex-  
 1212 clusion while directly using the nominal analysis result (assuming SM acceptance/kinematics).  
 1213 The subsequent section then takes into account acceptance/kinematic differences by showing  
 1214 the result with dedicated samples for each  $\hat{H}$  point.

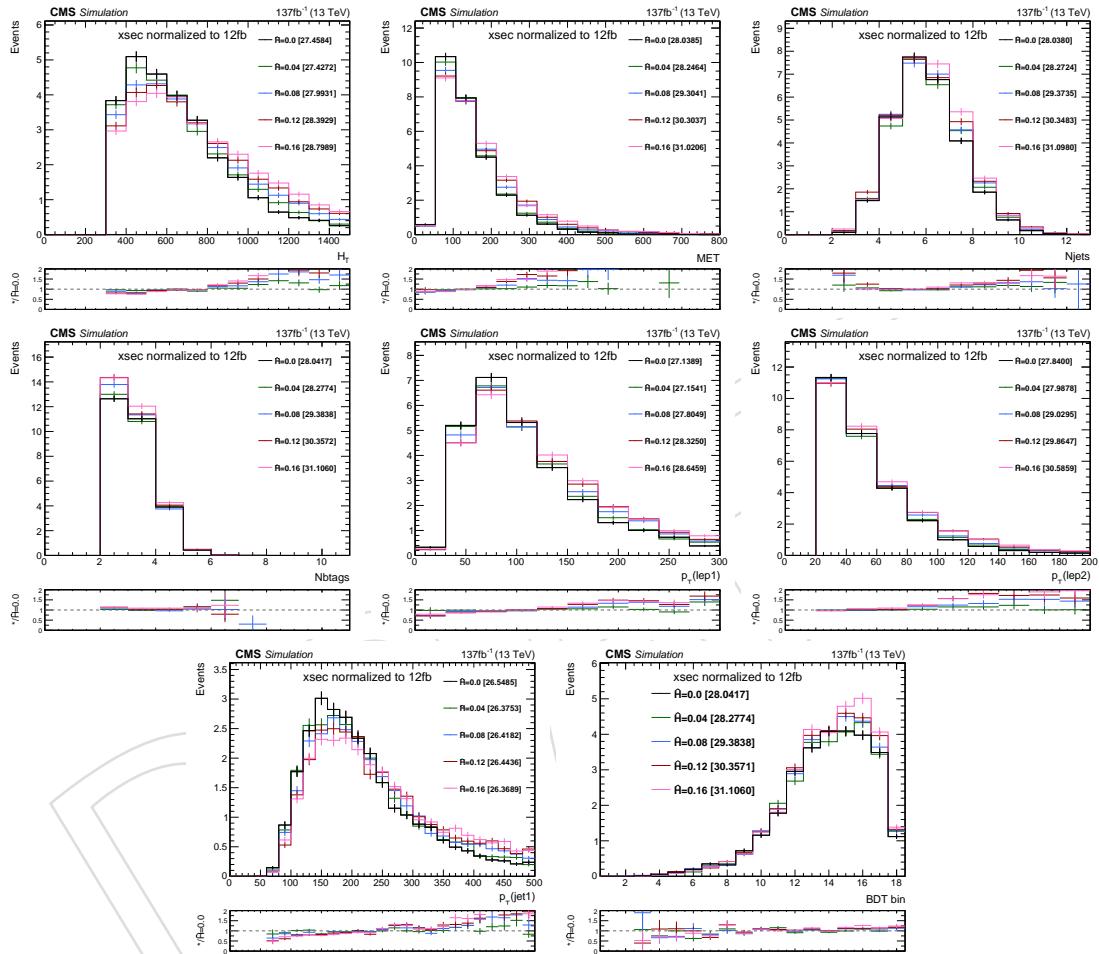


Figure 74: Distributions comparing  $H_T$ ,  $E_T^{\text{miss}}$ ,  $N_{\text{jets}}$ ,  $N_{\text{bjets}}$ ,  $p_T$  of the leading lepton,  $p_T$  of the subleading lepton,  $p_T$  of the leading jet, and the BDT signal region yields for three values of the oblique parameter  $\hat{R}$ . Both shapes and acceptance are relevant here, as cross-section has been normalized to  $12 \text{ fb}^{-1}$ .

#### 12.4.4 Exclusion, assuming SM kinematics

1216 Through correspondence with the authors, we recalculated the 14TeV formula and find cross-  
 1217 sections to match closely with the predicted formula above, for values  $\hat{H} < 0.04$ , as the refer-  
 1218 ence expands only to second order due to their tighter region of interest (HL-LHC projections).  
 1219 Figure 75 shows a recalculation of the 14TeV formula instead at 13TeV, as well as a cubic fit  
 1220 to the values in order to extract an upper limit on the oblique parameter using the nominal  
 1221 analysis BDT upper limit. We find  $\hat{H} < 0.13$  at 95% CL.

1222 As is the case for the top Yukawa interpretation from Section 12.2, the  $t\bar{t}H$  background is also  
 1223 subject to modifications due to non-zero  $\hat{H}$ . The  $t\bar{t}H$  background is affected by a cross-section  
 1224 scaling due to the reduction of the top yukawa constant ( $y_t \rightarrow y_t - \hat{H}$ ). Thus, we can utilize the  
 1225 observed upper limits calculated for the  $y_t < 1$  section of the curve in Figure 66. An increase in  
 1226  $\hat{H}$  corresponds to slight increase of the observed upper limit, which translates into a weakening  
 1227 of the  $\hat{H}$  constraint by 0.005 (4% relative).

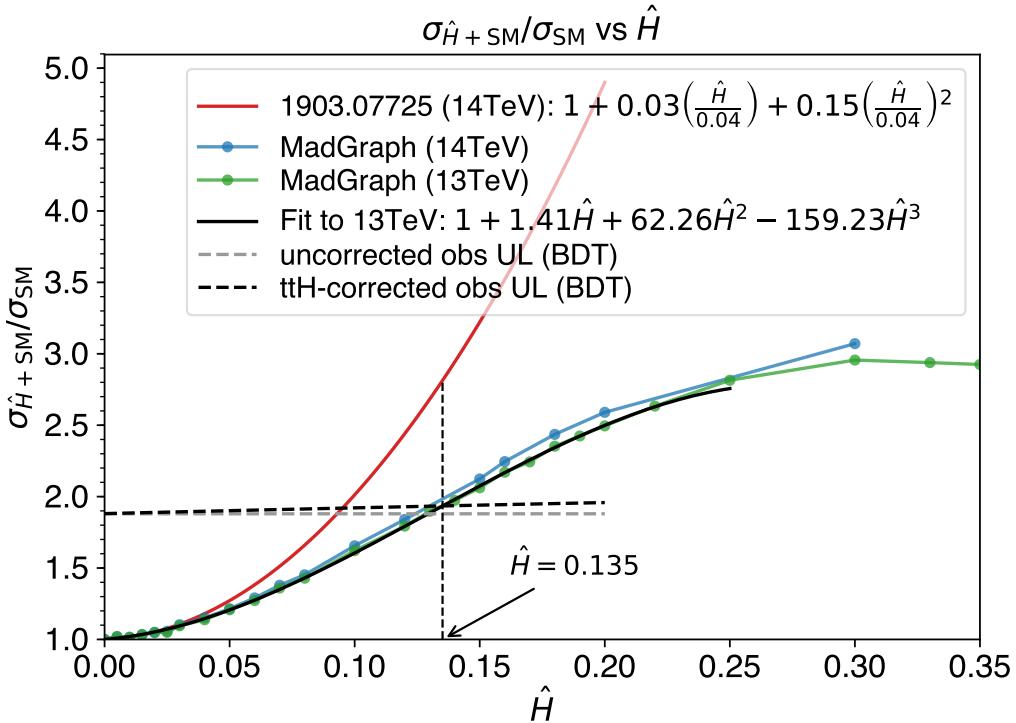


Figure 75: Cross-section (normalized to SM) as a function of oblique parameter  $\hat{H}$ . The red curve represents equation 5.4 from Ref. [32] at 14TeV. The blue and green curves are private calculations from MadGraph at 14TeV and 13TeV, respectively. The solid black curve is a cubic fit to the 13TeV private calculation, which is used to find the intersection with the (nearly) horizontal black dotted line (nominal analysis upper limit, corrected to account for deviations from  $t\bar{t}H$ , as described in the text. For the sake of comparison, the (completely) horizontal gray dotted line does not take this effect into account. The intersection of the black dotted lines provides the upper limit on  $\hat{H}$ .

#### 12.4.5 Exclusion using dedicated samples

1229 Repeating the previous subsection with privately-generated dedicated samples yields Figure 76  
 1230 and excludes  $\hat{H} > 0.12$ . Explicitly, we substitute the privately-generated  $\hat{H}$ -modified  $t\bar{t}t\bar{t}$  sam-

1231 ple for the  $t\bar{t}t\bar{t}$  signal, scale down the  $t\bar{t}H$  background normalization by  $y_t^2 = (1 - \hat{H})^2$  ( $\approx 0.7$   
1232 at  $\hat{H} = 0.16$ ), and run the nominal limit-setting procedure.

1233 Note that the upper limit becomes more stringent for higher values of  $\hat{H}$ , consistent with the  
1234 increase in acceptance shown in Figure 74.

1235 We have produced official gridpacks and asked the TOP MC contacts to generate official CMS  
1236 samples using the configurations for 2016, 2017 and 2018 respectively. These samples are ex-  
1237 pected to be identical to the ones we generated privately, since the same process will be used.

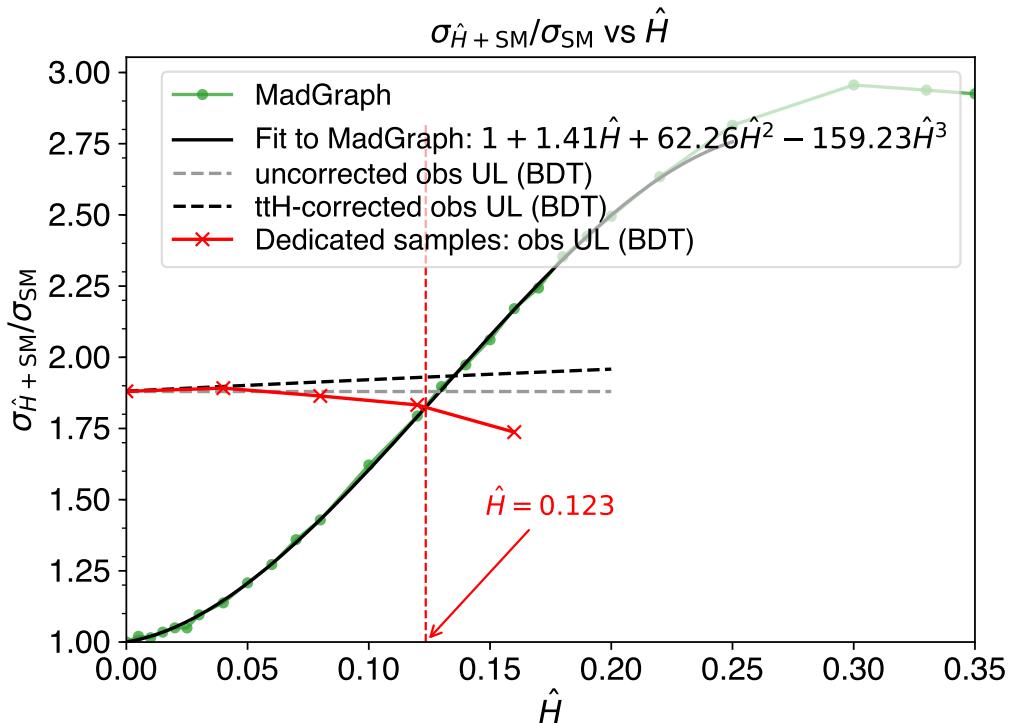


Figure 76: See caption for Figure 75. An additional red curve was added connecting the five upper limits calculated with dedicated samples. The intersection of a linear interpolation between these points (in red) with the black curve defines the 95% CL upper limit on  $\hat{H}$ .

## 12.5 Dark matter

### 12.5.1 Introduction

Upper limits on additional  $t\bar{t}t\bar{t}$  production can be translated into exclusions/upper limits on simplified dark matter models which have a scalar or pseudoscalar mediator decaying into a pair of dark matter or standard model particles. The production of the mediator (decaying into invisible dark matter) in association with a pair of top quarks was performed by CMS with the 2016 dataset in Ref. [33]. Reference [34] found that production of the mediator in association with a single top quark contributes meaningfully to the total cross-section. Thus, a search that also included these single top associated production modes was carried out with the same dataset in Ref. [28], which we aim to complement in this analysis with an orthogonal final state. A brief overview is included below, and further theoretical details and motivation can be found in the Analysis Note for EXO-18-010 ([35]).

In the framework of a simplified dark matter model, where the scalar ( $\phi$ ) or pseudoscalar ( $a$ ) mediator couples dark matter and SM particles, the relevant terms of the interaction lagrangian are of the form

$$\mathcal{L}_\phi = g_\chi \phi \bar{\chi} \chi + \frac{g_q \phi}{\sqrt{2}} \sum_f y_f \bar{f} f \quad \mathcal{L}_a = i g_\chi a \bar{\chi} \gamma^5 \chi + \frac{i g_q a}{\sqrt{2}} \sum_f y_f \bar{f} \gamma^5 f$$

where  $y_f$  are the fermionic yukawa couplings. The coupling constants  $g_\chi$  and  $g_q$  give the relative strengths of the mediator coupling to dark matter and SM particles, and are used interchangeably with  $g_{\text{DM}}$  and  $g_{\text{SM}}$ , respectively. The model has four free parameters ( $g_\chi$ ,  $g_q$ ,  $m_\chi$ , and  $m_a$ ) which is reduced to two with the assumption of  $g_\chi = g_q = 1$ .

The relevant production diagrams are shown in Figure 77. When the mediator mass is above  $2m_{\text{top}} \approx 350\text{GeV}$ , on-shell decay to  $t\bar{t}$  becomes kinematically accessible, resulting in 3 or 4 top quark final states, so we instead consider a version of the diagrams with a decay of the mediator into  $t\bar{t}$  rather than a pair of (invisible) dark matter particles. Consequently, the production diagrams and kinematics are identical to that of the 2HDM interpretation which will allow us to use existing simulation samples for this interpretation, and we only need to calculate branching ratios into  $t\bar{t}$ .

In this way, our final state is complementary to the  $t\bar{t} + \text{MET}$  final state used in Ref. [28] when the DM mass becomes large.

Using MadGraph 2.6.5, and the models from the official gridpacks used for the EXO-18-010 result we first verified consistency with the corresponding AN cross-sections before modifying the  $\phi/a \rightarrow \chi\bar{\chi}$  decay to be  $\phi/a \rightarrow t\bar{t}$ . We use MadGraph to calculate cross-section (times branching ratio) for the  $t\bar{t}$  associated process and two single-top processes for mediator masses matching the available 2HDM simulation samples at varying dark-matter masses. We use the NNPDF30\_lo\_as\_0130 PDF.

While we use MadGraph-calculated values in the end to account for all effects, we independently verified that relative branching ratios for scalar and pseudoscalar, respectively, are given by

$$\Gamma_{\text{ratio}} \equiv \frac{\Gamma(\phi \rightarrow t\bar{t})}{\Gamma(\phi \rightarrow \chi\bar{\chi})} = \frac{3g_{\text{SM}}^2 y_t^2}{2g_{\text{DM}}^2} \left( \frac{M^2 - 4m_t^2}{M^2 - 4m_\chi^2} \right)^{3/2} \quad (7)$$

$$\Gamma_{\text{ratio}} \equiv \frac{\Gamma(\phi \rightarrow t\bar{t})}{\Gamma(\phi \rightarrow \chi\bar{\chi})} = \frac{3g_{\text{SM}}^2 y_t^2}{2g_{\text{DM}}^2} \left( \frac{M^2 - 4m_t^2}{M^2 - 4m_\chi^2} \right)^{1/2} \quad (8)$$

where  $y_t$  is the top-quark yukawa coupling,  $M$  is the mediator mass, and  $\phi$  represents both the scalar and pseudoscalar. Note that the only difference is in the exponent of the mass term. In practice, these formulae would be used after calculating the production cross-section with  $\text{BR}(\phi \rightarrow t\bar{t})=1$  with MadGraph for each mediator mass value (e.g.,  $M = 450\text{GeV}$  with  $m_\chi > 450\text{GeV}$  so that the DM decay is kinematically suppressed.) Then that production cross-section is multiplied by  $\text{BR}(t\bar{t}) = \frac{\Gamma(t\bar{t})}{\Gamma(t\bar{t}) + \Gamma(\text{DM})} = \left(1 - \Gamma_{\text{ratio}}^{-1}\right)^{-1}$  in order to get the  $\sigma \cdot \text{BR}$  at an arbitrary  $m_\chi$ . The formulae assume that there are only two accessible decay modes of the mediator.

### 12.5.2 Exclusions

As noted previously, the processes are identical to those of the 2HDM interpretation with a different labeling of the mediator, so we can calculate a single cross-section upper limit for a given mediator mass using pre-existing samples. For this interpretation the SM  $t\bar{t}t\bar{t}$  process is treated as a background and assigned its SM cross section and uncertainty  $12.0^{+2.2}_{-2.5} \text{ fb}$  [6]. The single value is used to exclude along the vertical axis ( $m_{\text{DM}}$ ) since kinematics remain identical and only the branching ratio to  $t\bar{t}$  changes. This can be seen in Figure 78, which overlays exclusion contours on cross-sections calculated at discrete points. The diagonal line corresponds to the kinematic boundary between on- and off-shell mediator decays into dark matter,  $m_{\text{mediator}} = 2m_{\text{DM}}$ , above which  $t\bar{t}$  is the dominant decay. For couplings set to unity, the region above the diagonal is excluded for both mediator types between approximately 350 and 500 GeV. When decreasing the relative branching ratio to DM ( $g_{\text{DM}} = 0.5$ ), DM masses down to 1 GeV are excluded up to a pseudoscalar mediator mass of 470GeV. Using Eqs. 7-8, Figure 79 shows that the gain is larger in the pseudoscalar case when compared to the scalar mediator due to a steeper decrease in  $t\bar{t}$  branching ratio for the scalar mediator.

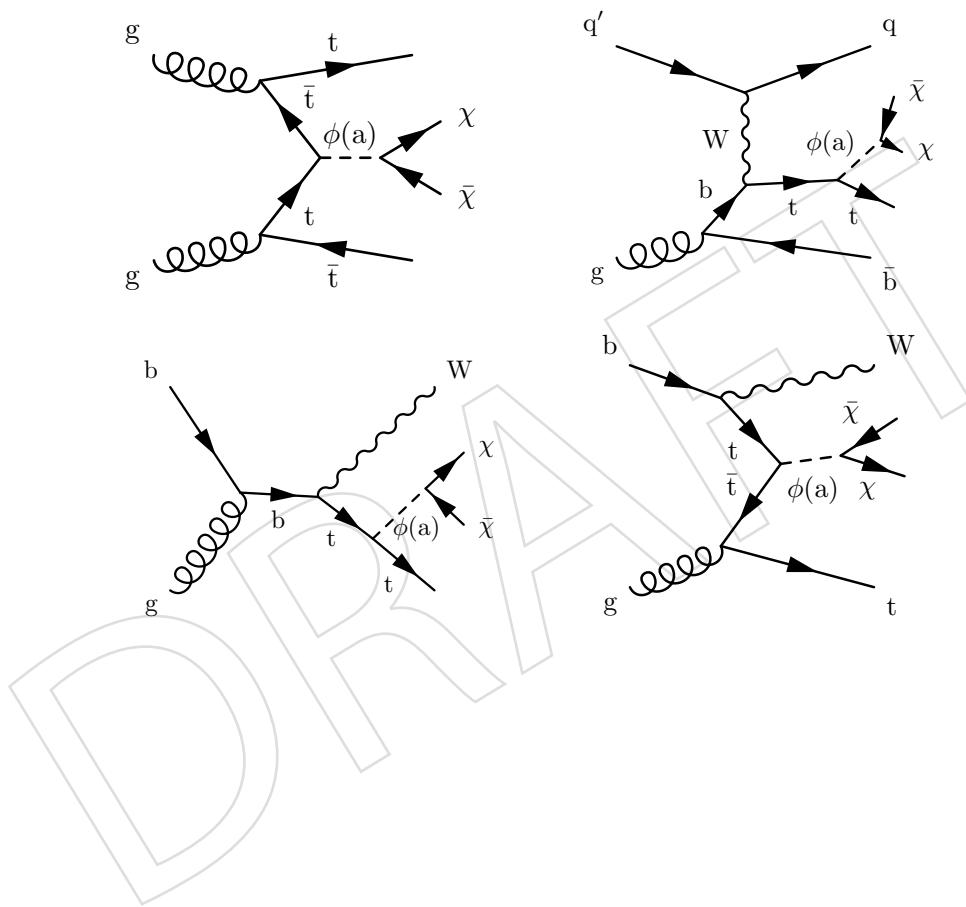


Figure 77: Diagrams for scalar (pseudoscalar) mediator production in association with a  $t\bar{t}$  pair (top left), associated t-channel single top (top right), associated  $tW$  (bottom row). In these diagrams the mediator decays into a pair of invisible particles, while for the  $t\bar{t}t\bar{t}$  analysis we are probing the equivalent diagrams where the mediator decays back to  $t\bar{t}$ . As the s-channel single top production cross-section is relatively negligible (by at least an order of magnitude), it is not included in this result.

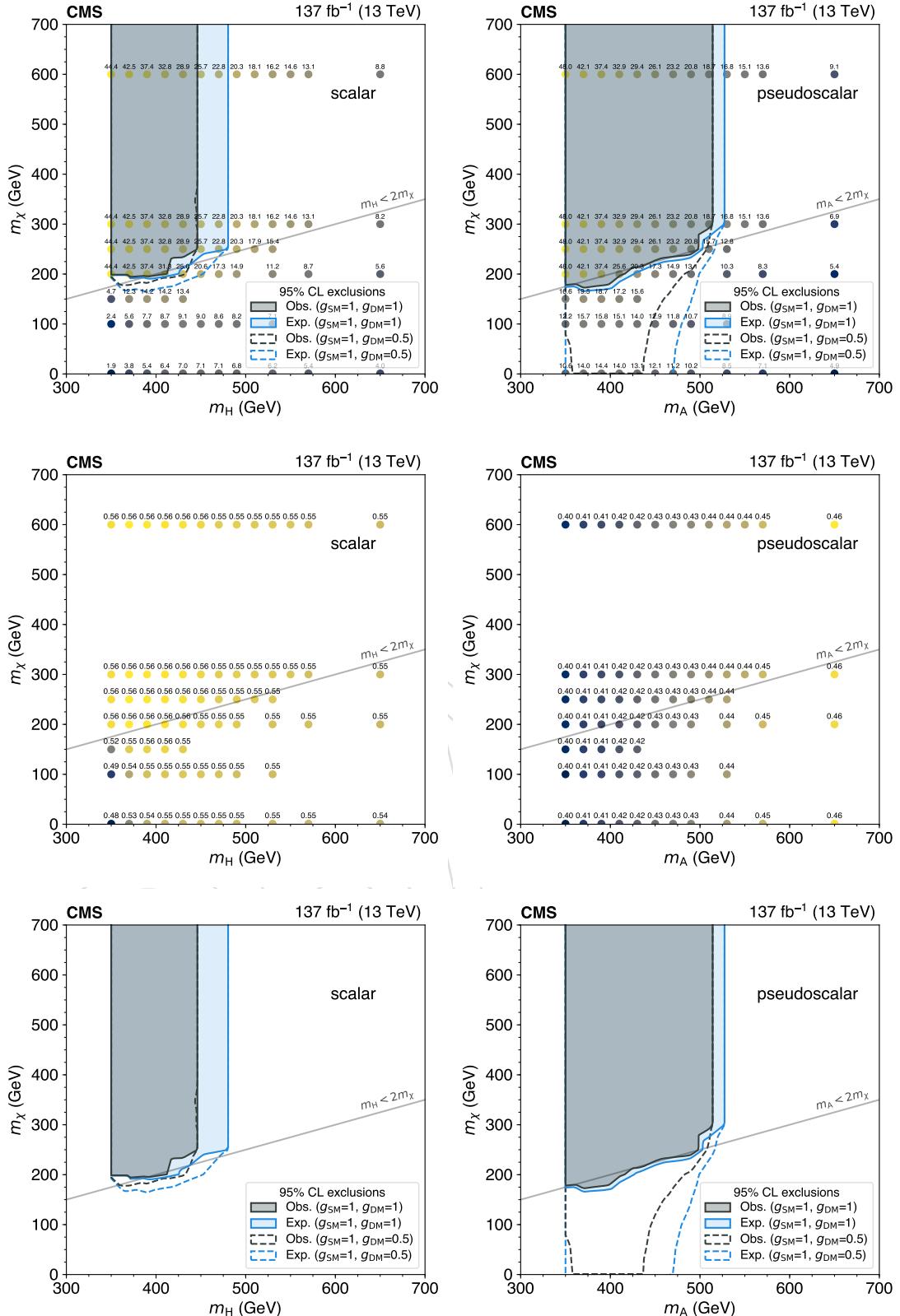


Figure 78: Expected and observed 95% CL exclusions in the plane of  $m_{\text{DM}}\text{-}m_{\text{mediator}}$  for scalar (left) and pseudoscalar (right) mediators. The top row shows two exclusions: the nominal assumption ( $g_{\text{DM}} = g_{\text{SM}} = 1$ ), and an alternate assumption ( $g_{\text{DM}} = 0.5, g_{\text{SM}} = 1$ ). Calculated cross-section times branching ratio values for the first assumption are shown as markers with accompanying values in units of fb. The middle row shows the fraction of the total production cross-section attributed to single top processes. The the bottom row is identical to the top row but without cross-section value/markers for presentational purposes.

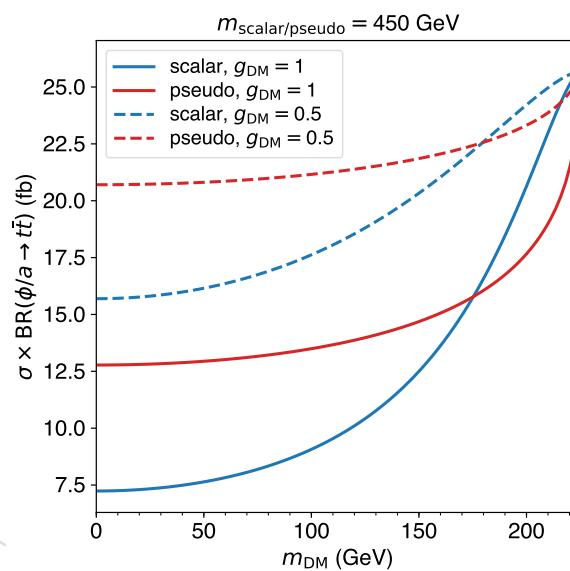


Figure 79: Cross-section times branching ratio to  $t\bar{t}$  for scalar and pseudoscalar mediators with mass 450 GeV, shown at DM coupling values of  $g_{\text{DM}} = 1$  and  $g_{\text{DM}} = 0.5$ , calculated with Eqs. 7-8.

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## 1386 A Statistical checks

### 1387 A.1 Impacts

1388 The leading 30 nuisance impacts for two sets of impacts, expected and observed for cut-based  
 1389 and BDT analyses, are shown in Figure 117 (expected cut-based analysis), Figure 118 (observed  
 1390 cut-based analysis), Figure 119 (expected BDT analysis), and Figure 120 (observed BDT anal-  
 1391 ysis). The leading expected nuisance in both cases corresponds to the  $\sigma(\text{ttbb})/\sigma(\text{ttjj})$  scaling.  
 1392 Note that the “prop binSS” nuisances for MC statistics include (and are dominated by) tight-  
 1393 loose sideband statistics.

1394 The obseved pulls show the most constrained/pulled nuisances correspond to normalization  
 1395 parameters for ttW and ttZ, as we would expect from the control regions. “TTWSF” is moved  
 1396 by approximately  $1\sigma$  ( $0.8\sigma$ ) with respect to the input nuisance sizes for the cut-based (BDT)  
 1397 analysis. “TTZSF” is moved up by approximately  $0.6\sigma$  ( $0.7\sigma$ ) for the cut-based (BDT) analysis.

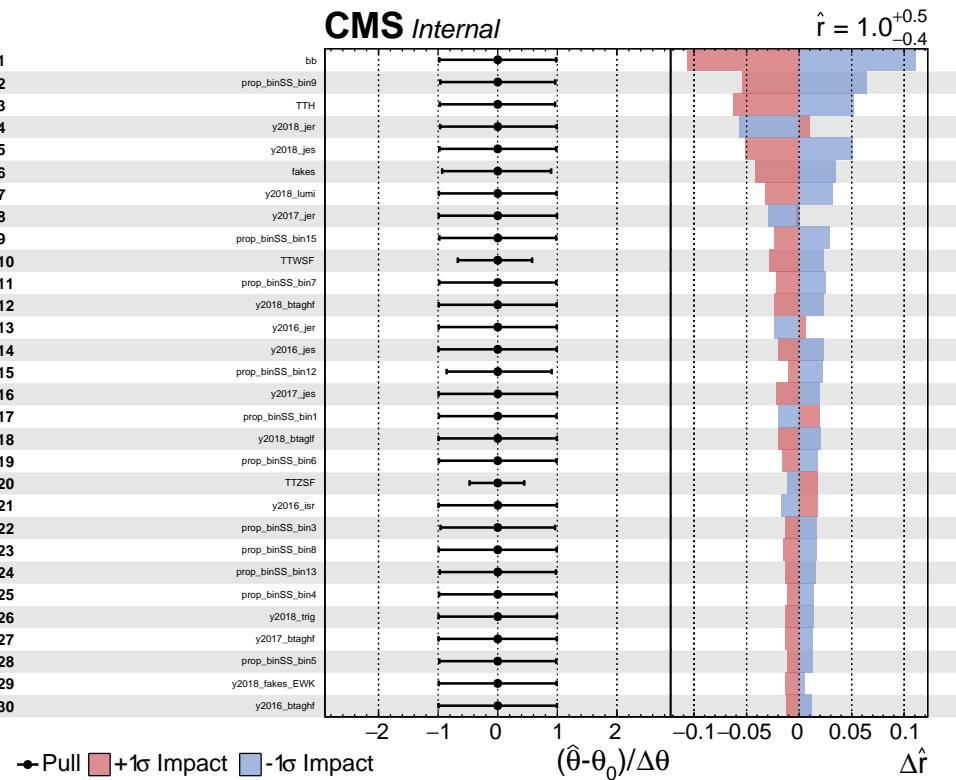


Figure 80: Expected nuisance impacts for the cut-based analysis.

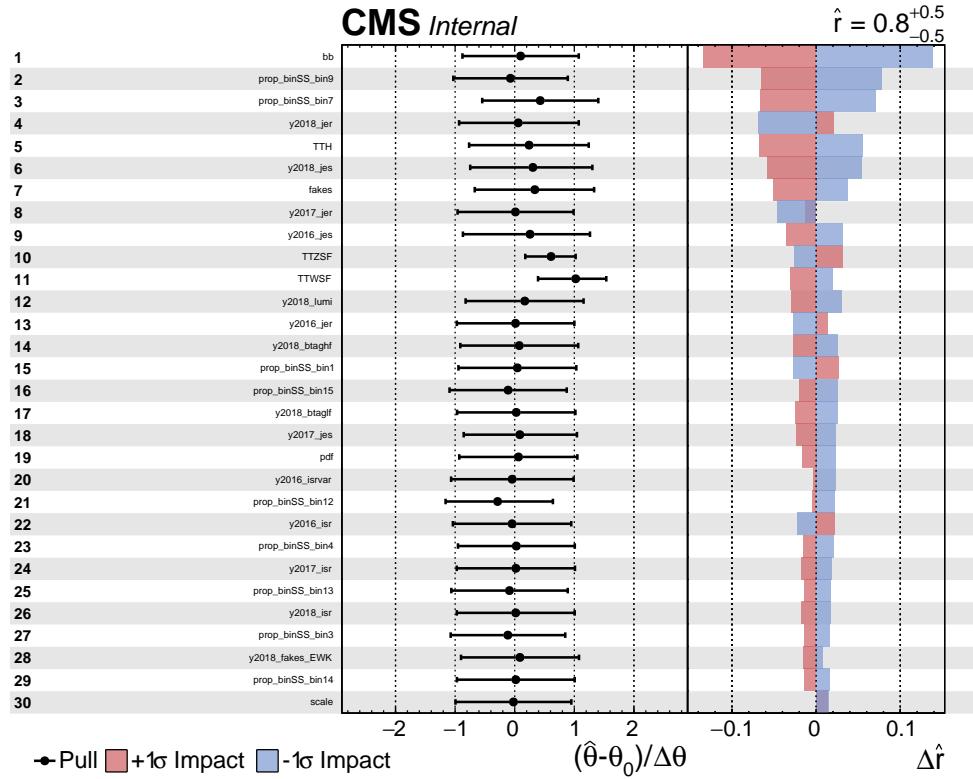


Figure 81: Observed nuisance impacts for the cut-based analysis.

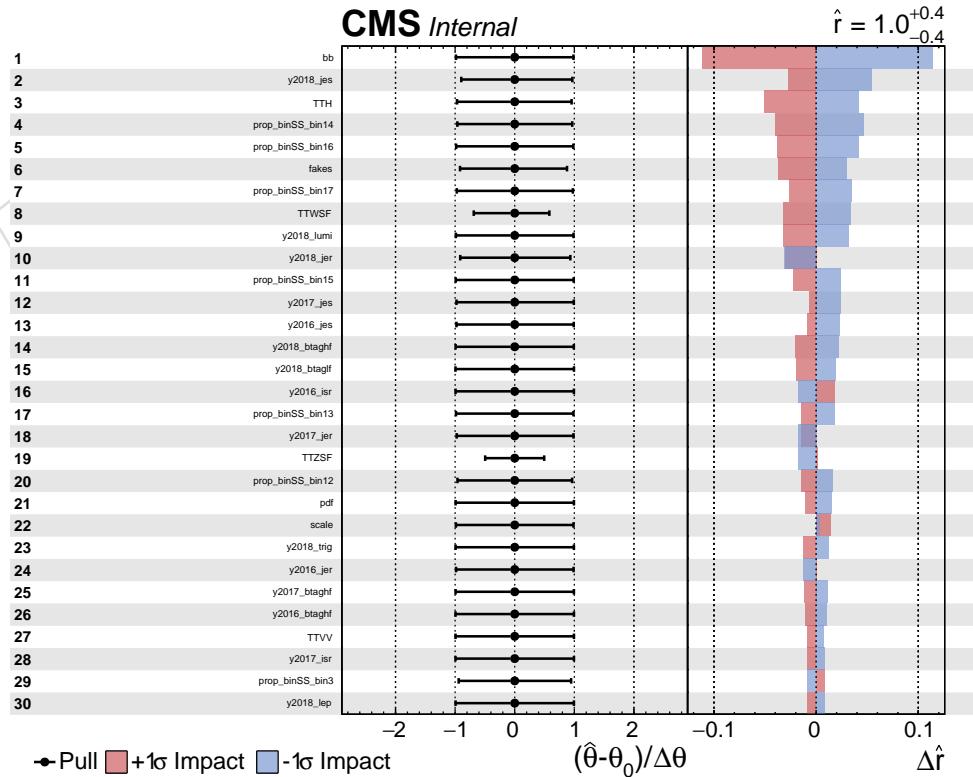


Figure 82: Expected nuisance impacts for the BDT-based analysis.

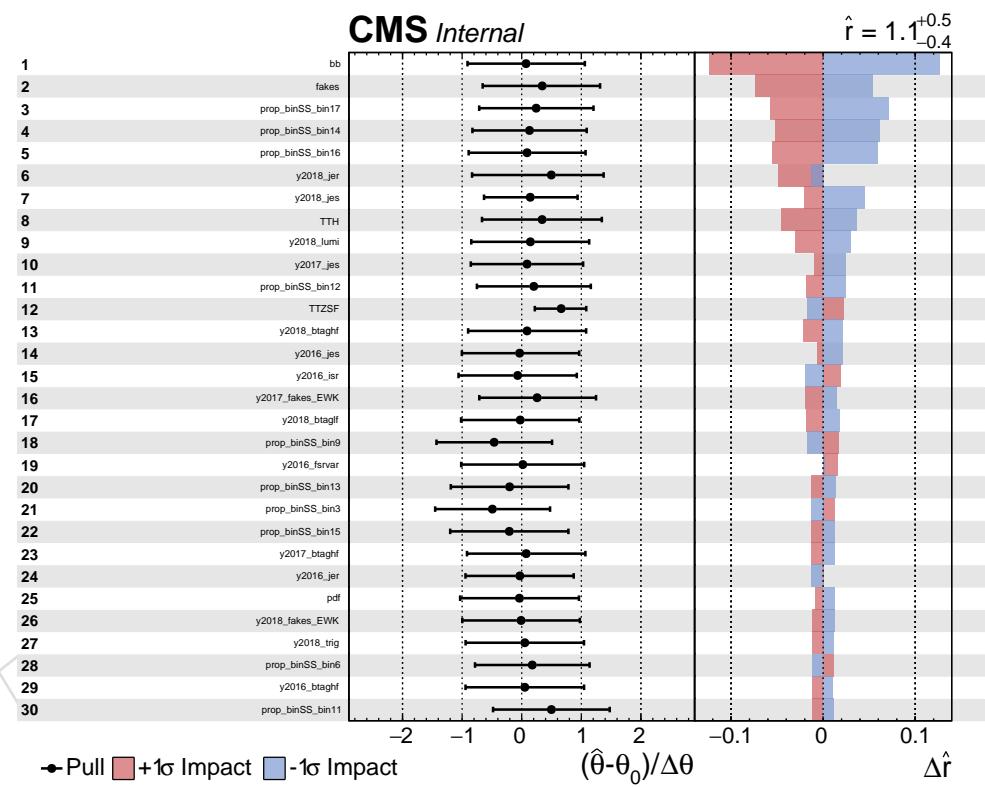


Figure 83: Observed nuisance impacts for the BDT-based analysis.

1398 **A.2 Nuisance forms**

1399 Nuisance functional forms and widths are tabulated in Table 30.

Table 30: Rows preceded by (s) apply to signal as well. For “shape” uncertainties, width is not applicable as up and down variation envelopes are taken as the systematic uncertainties.

name	function	width
lumi	lnN	1.021-1.05
(s) jes	shape	-
(s) jer	shape	-
isr	shape	-
bb	shape	-
(s) lep	shape	-
(s) trig	shape	-
(s) btaghf	shape	-
(s) btaglf	shape	-
(s) isrvvar	shape	-
(s) fsrvvar	shape	-
(s) scale	shape	-
(s) pdf	shape	-
(s) alphas	shape	-
(s) pu	shape	-
TTWSF	lnN	1.4
TTZSF	lnN	1.4
TTH	lnN	1.25
TTVV	lnN	1.11
XG	lnN	1.11
rares	lnN	1.2
fakes	lnN	1.3
fakes_EWK	shape	-
flips	lnN	1.2

**A.3 Nuisances**

- Two sets of nuisance pull values, expected and observed for cut-based and BDT analyses, are tabulated in Table 62 (expected cut-based analysis), Table 63 (observed cut-based analysis), Table 64 (expected BDT analysis), and Table 65 (observed BDT analysis).
- The most constrained nuisances correspond to normalization parameters for ttW and ttZ (“TTWSF” and “TTZSF”) due to high statistics in control regions and in the bulk (BDT). Their input normalization uncertainty is 40%.

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Table 31: Expected nuisance pulls for the cut-based analysis. The final column indicates the correlation between the nuisance and the signal strength

name	<i>b</i> -only fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	<i>s + b</i> fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	<b>+0.38, 0.99</b>	+0.00, 0.97	-0.12
TTVV	+0.05, 1.00	+0.00, 0.99	-0.02
TTWSF	<b>+0.09, 0.61</b>	<b>+0.00, 0.61</b>	-0.06
TTZSF	<b>-0.02, 0.45</b>	<b>+0.00, 0.45</b>	+0.02
XG	-0.01, 0.99	+0.00, 0.99	+0.00
alphas	+0.00, 0.99	-0.00, 0.99	+0.00
bb	<b>+0.74, 0.94</b>	+0.00, 0.98	-0.24
fakes	+0.24, 0.92	+0.00, 0.92	-0.08
pdf	+0.07, 0.99	-0.00, 0.99	-0.02
rares	-0.01, 0.99	+0.00, 0.99	+0.00
scale	-0.19, 0.96	-0.00, 0.99	+0.00
y2016_btaghf	+0.04, 0.99	+0.00, 0.99	-0.03
y2016_btagnf	+0.03, 0.99	+0.00, 0.99	-0.01
y2016_fakes_EWK	+0.04, 1.00	+0.00, 0.99	-0.01
y2016_flips	-0.02, 0.99	+0.00, 0.99	+0.00
y2016_fsrvvar	+0.00, 0.99	+0.00, 0.99	+0.02
y2016_isr	-0.10, 0.99	-0.00, 0.99	+0.04
y2016_isrvvar	+0.00, 0.99	+0.00, 0.99	-0.02
y2016_jer	<b>-0.13, 1.12</b>	-0.00, 0.99	+0.03
y2016_jes	+0.07, 0.97	+0.00, 0.99	-0.05
y2016_lep	+0.01, 0.99	+0.00, 0.99	-0.01
y2016_lumi	+0.02, 0.99	+0.00, 0.99	-0.02
y2016_prefire	-0.00, 0.99	-0.00, 0.99	+0.00
y2016_pu	+0.01, 0.99	-0.00, 0.99	+0.00
y2016_trig	+0.01, 0.99	+0.00, 0.99	-0.02
y2017_btaghf	+0.04, 0.99	+0.00, 0.99	-0.03
y2017_btagnf	+0.03, 0.99	+0.00, 0.99	-0.01
y2017_fakes_EWK	+0.00, 0.99	-0.00, 0.98	-0.01
y2017_flips	-0.02, 0.99	+0.00, 0.99	+0.00
y2017_fsrvvar	+0.00, 0.99	+0.00, 0.99	+0.00
y2017_isr	+0.07, 0.99	+0.00, 0.99	-0.03
y2017_isrvvar	+0.00, 0.99	-0.00, 0.99	-0.01
y2017_jer	<b>-0.12, 1.20</b>	-0.00, 0.99	+0.03
y2017_jes	+0.07, 0.99	+0.00, 0.99	-0.05
y2017_lep	+0.01, 0.99	+0.00, 0.99	-0.01
y2017_lumi	+0.01, 0.99	+0.00, 0.99	-0.02
y2017_prefire	-0.01, 0.99	-0.00, 0.99	+0.01
y2017_pu	-0.01, 1.00	+0.00, 0.99	+0.00
y2017_trig	+0.01, 0.99	+0.00, 0.99	-0.02
y2018_btaghf	+0.09, 0.99	+0.00, 0.99	-0.05
y2018_btagnf	+0.12, 0.99	-0.00, 0.99	-0.04
y2018_fakes_EWK	+0.08, 1.03	+0.00, 0.99	-0.02
y2018_flips	-0.03, 0.99	+0.00, 0.99	+0.01
y2018_fsrvvar	+0.00, 0.99	+0.00, 0.99	+0.00
y2018_isr	+0.08, 0.99	+0.00, 0.99	-0.03
y2018_isrvvar	+0.00, 0.99	+0.00, 0.99	+0.00
y2018_jer	<b>-0.33, 1.08</b>	-0.00, 0.99	+0.08
y2018_jes	+0.18, 0.96	-0.00, 0.98	-0.11
y2018_lep	+0.02, 0.99	+0.00, 0.99	-0.02
y2018_lumi	+0.03, 0.99	+0.00, 0.99	-0.04
y2018_pu	+0.00, 1.00	-0.00, 0.99	+0.00
y2018_trig	+0.02, 0.99	+0.00, 0.99	-0.03

Table 32: Observed nuisance pulls for the cut-based analysis. The final column indicates the correlation between the nuisance and the signal strength

name	$b$ -only fit $\Delta x / \sigma_{\text{in}}, \sigma_{\text{out}} / \sigma_{\text{in}}$	$s + b$ fit $\Delta x / \sigma_{\text{in}}, \sigma_{\text{out}} / \sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	<b>+0.48, 1.02</b>	+0.24, 1.00	-0.12
TTVV	+0.06, 1.00	+0.03, 1.00	-0.02
TTWSF	<b>+1.03, 0.57</b>	<b>+0.99, 0.56</b>	-0.06
TTZSF	<b>+0.57, 0.41</b>	<b>+0.60, 0.42</b>	+0.05
XG	+0.04, 1.00	+0.04, 1.00	+0.00
alphas	+0.00, 0.99	+0.00, 0.99	+0.00
bb	<b>+0.62, 0.92</b>	+0.10, 0.97	-0.27
fakes	<b>+0.54, 1.02</b>	<b>+0.33, 1.00</b>	-0.09
pdf	+0.15, 0.98	+0.06, 1.00	-0.04
rares	+0.06, 1.00	+0.07, 1.00	+0.01
scale	-0.08, 0.91	-0.02, 0.96	-0.00
y2016_btaghf	+0.07, 0.99	+0.04, 0.99	-0.03
y2016_btaglf	+0.03, 0.99	+0.00, 0.99	-0.01
y2016_fakes_EWK	+0.08, 1.00	+0.05, 0.99	-0.01
y2016_flips	+0.00, 0.99	+0.01, 1.00	+0.01
y2016_fsrvvar	+0.00, 0.99	-0.01, 1.04	+0.01
y2016_isr	-0.12, 0.99	-0.04, 0.99	+0.04
y2016_isrvvar	+0.00, 0.99	-0.04, 1.06	-0.03
y2016_jer	-0.07, 1.04	+0.01, 0.99	+0.04
y2016_jes	<b>+0.38, 1.08</b>	<b>+0.25, 1.12</b>	-0.08
y2016_lep	+0.04, 0.99	+0.03, 0.99	-0.01
y2016_lumi	+0.05, 1.00	+0.04, 1.00	-0.02
y2016_prefire	-0.01, 0.99	-0.00, 0.99	+0.00
y2016_pu	+0.03, 0.99	+0.01, 0.99	-0.01
y2016_trig	+0.04, 0.99	+0.03, 0.99	-0.01
y2017_btaghf	+0.10, 1.00	+0.06, 0.99	-0.03
y2017_btaglf	+0.03, 0.99	+0.00, 0.99	-0.01
y2017_fakes_EWK	+0.23, 0.98	+0.22, 0.97	-0.01
y2017_flips	+0.00, 0.99	+0.01, 1.00	+0.01
y2017_isrvvar	+0.00, 0.99	-0.01, 0.99	+0.00
y2017_isr	+0.08, 0.99	+0.02, 0.99	-0.04
y2017_isrvvar	+0.00, 0.99	-0.01, 0.99	-0.01
y2017_jer	<b>-0.07, 1.24</b>	+0.01, 0.96	+0.03
y2017_jes	+0.15, 0.93	+0.09, 0.93	-0.05
y2017_lep	+0.04, 0.99	+0.03, 0.99	-0.01
y2017_lumi	+0.05, 1.00	+0.04, 1.00	-0.01
y2017_prefire	-0.02, 0.99	-0.02, 0.99	+0.01
y2017_pu	+0.03, 1.00	+0.01, 1.00	-0.00
y2017_trig	+0.04, 0.99	+0.04, 0.99	-0.01
y2018_btaghf	+0.15, 0.99	+0.07, 0.99	-0.05
y2018_btaglf	+0.12, 0.99	+0.02, 0.99	-0.05
y2018_fakes_EWK	+0.15, 1.01	+0.09, 0.99	-0.02
y2018_flips	+0.01, 0.99	+0.02, 1.00	+0.01
y2018_isrvvar	+0.00, 0.99	+0.00, 0.99	+0.00
y2018_isr	+0.08, 0.99	+0.02, 0.99	-0.04
y2018_isrvvar	+0.00, 0.99	+0.00, 0.99	+0.00
y2018_jer	<b>-0.19, 1.17</b>	+0.06, 1.05	+0.09
y2018_jes	<b>+0.48, 1.02</b>	<b>+0.31, 1.08</b>	-0.13
y2018_lep	+0.06, 0.99	+0.05, 0.99	-0.02
y2018_lumi	+0.10, 0.99	+0.08, 0.99	-0.03
y2018_pu	+0.00, 1.00	-0.00, 0.99	+0.00
y2018_trig	+0.08, 0.99	+0.07, 0.99	-0.02

Table 33: Expected nuisance pulls for the BDT analysis. The final column indicates the correlation between the nuisance and the signal strength

name	<i>b</i> -only fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	<i>s + b</i> fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	<b>+0.32, 0.98</b>	+0.00, 0.96	-0.11
TTVV	+0.06, 1.00	+0.00, 0.99	-0.02
TTWSF	<b>+0.08, 0.63</b>	<b>+0.00, 0.62</b>	-0.08
TTZSF	<b>-0.04, 0.51</b>	<b>+0.00, 0.48</b>	+0.03
XG	-0.01, 0.99	+0.00, 0.99	+0.00
alphas	-0.00, 0.99	-0.00, 0.99	+0.00
bb	<b>+0.91, 0.93</b>	+0.00, 0.99	-0.27
fakes	+0.28, 0.91	<b>+0.00, 0.90</b>	-0.08
pdf	+0.12, 0.97	-0.00, 0.99	-0.03
rares	-0.02, 0.99	+0.00, 0.99	+0.00
scale	-0.25, 0.97	+0.00, 0.99	+0.01
y2016_btaghf	+0.04, 0.99	+0.00, 0.99	-0.03
y2016_btagnf	+0.03, 0.99	+0.00, 0.99	-0.01
y2016_fakes_EWK	+0.03, 0.99	+0.00, 0.99	-0.01
y2016_flips	-0.02, 0.99	+0.00, 0.99	+0.00
y2016_fsrvvar	-0.00, 0.99	+0.00, 0.99	+0.01
y2016_isr	-0.13, 0.99	+0.00, 0.99	+0.04
y2016_isrvvar	-0.00, 0.99	-0.00, 0.99	-0.01
y2016_jer	-0.05, 1.03	+0.00, 0.99	+0.01
y2016_jes	+0.07, 0.90	+0.00, 0.99	-0.04
y2016_lep	+0.01, 0.99	+0.00, 0.99	-0.01
y2016_lumi	+0.02, 0.99	+0.00, 0.99	-0.02
y2016_prefire	-0.00, 0.99	-0.00, 0.99	+0.00
y2016_pu	+0.04, 0.99	-0.00, 0.99	-0.01
y2016_trig	+0.01, 0.99	+0.00, 0.99	-0.02
y2017_btaghf	+0.03, 0.99	+0.00, 0.99	-0.03
y2017_btagnf	+0.03, 0.99	+0.00, 0.99	-0.01
y2017_fakes_EWK	+0.03, 0.98	+0.00, 0.97	-0.02
y2017_flips	-0.02, 0.99	+0.00, 0.99	+0.01
y2017_fsrvvar	-0.00, 0.99	-0.00, 0.99	+0.00
y2017_isr	+0.07, 0.99	-0.00, 0.99	-0.02
y2017_isrvvar	-0.00, 0.99	+0.00, 0.99	-0.01
y2017_jer	<b>-0.01, 1.19</b>	+0.00, 0.99	+0.00
y2017_jes	+0.07, 0.90	+0.00, 0.99	-0.04
y2017_lep	+0.00, 0.99	+0.00, 0.99	-0.02
y2017_lumi	+0.01, 0.99	+0.00, 0.99	-0.02
y2017_prefire	-0.00, 0.99	-0.00, 0.99	+0.01
y2017_pu	+0.03, 1.00	-0.00, 0.99	-0.01
y2017_trig	+0.00, 0.99	+0.00, 0.99	-0.02
y2018_btaghf	+0.08, 0.99	+0.00, 0.99	-0.05
y2018_btagnf	+0.12, 0.99	-0.00, 0.99	-0.05
y2018_fakes_EWK	+0.06, 0.98	+0.00, 0.98	-0.01
y2018_flips	-0.03, 0.99	+0.00, 0.99	+0.01
y2018_fsrvvar	-0.00, 0.99	+0.00, 0.99	+0.00
y2018_isr	+0.04, 0.99	-0.00, 0.99	-0.02
y2018_isrvvar	-0.00, 0.99	+0.00, 0.99	+0.00
y2018_jer	<b>-0.01, 1.56</b>	+0.00, 0.95	+0.00
y2018_jes	<b>+0.15, 0.80</b>	-0.00, 0.95	-0.10
y2018_lep	+0.01, 0.99	+0.00, 0.99	-0.02
y2018_lumi	+0.03, 0.99	+0.00, 0.99	-0.04
y2018_pu	+0.02, 1.01	-0.00, 0.99	-0.01
y2018_trig	+0.02, 0.99	+0.00, 0.99	-0.03

Table 34: Observed nuisance pulls for the BDT analysis. The final column indicates the correlation between the nuisance and the signal strength

name	$b$ -only fit $\Delta x / \sigma_{\text{in}}, \sigma_{\text{out}} / \sigma_{\text{in}}$	$s + b$ fit $\Delta x / \sigma_{\text{in}}, \sigma_{\text{out}} / \sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	<b>+0.57, 1.02</b>	<b>+0.33, 1.00</b>	-0.09
TTVV	+0.08, 1.00	+0.04, 1.00	-0.02
TTWSF	<b>+0.64, 0.66</b>	<b>+0.75, 0.60</b>	-0.01
TTZSF	<b>+0.61, 0.42</b>	<b>+0.65, 0.43</b>	+0.04
XG	+0.02, 0.99	+0.04, 1.00	+0.00
alphas	+0.00, 0.99	+0.00, 1.00	+0.00
bb	<b>+0.87, 0.95</b>	+0.07, 0.98	-0.27
fakes	<b>+0.93, 0.99</b>	<b>+0.32, 0.99</b>	-0.14
pdf	+0.05, 0.99	-0.03, 1.01	-0.02
rares	+0.03, 0.99	+0.04, 0.99	+0.00
scale	-0.24, 0.94	-0.04, 0.94	+0.00
y2016_btaghf	+0.08, 0.99	+0.05, 0.99	-0.02
y2016_btaglf	+0.04, 1.00	+0.01, 0.99	-0.01
y2016_fakes_EWK	+0.16, 0.99	+0.08, 0.99	-0.02
y2016_flips	-0.01, 0.99	+0.00, 1.00	+0.00
y2016_fsrvar	+0.00, 0.99	+0.02, 1.07	+0.02
y2016_isr	-0.18, 0.99	-0.06, 0.99	+0.04
y2016_isrvar	+0.00, 0.99	+0.01, 1.00	-0.01
y2016_jer	<b>-0.06, 0.89</b>	<b>-0.03, 0.87</b>	+0.01
y2016_jes	+0.03, 1.00	<b>-0.04, 1.11</b>	-0.03
y2016_lep	+0.04, 0.99	+0.03, 0.99	-0.01
y2016_lumi	+0.05, 1.00	+0.04, 1.00	-0.02
y2016_prefire	-0.01, 0.99	-0.00, 0.99	+0.00
y2016_pu	-0.01, 0.99	-0.05, 1.00	-0.01
y2016_trig	+0.05, 0.99	+0.04, 0.99	-0.02
y2017_btaghf	+0.09, 0.99	+0.07, 0.99	-0.03
y2017_btaglf	+0.05, 0.99	+0.02, 0.99	-0.01
y2017_fakes_EWK	<b>+0.43, 1.04</b>	+0.25, 0.97	-0.04
y2017_flips	-0.01, 0.99	+0.01, 1.00	+0.00
y2017_fsrvar	+0.00, 0.99	+0.02, 1.00	+0.01
y2017_isr	+0.08, 0.99	+0.02, 0.99	-0.02
y2017_isrvar	+0.00, 0.99	+0.00, 1.00	-0.01
y2017_jer	<b>+0.09, 0.87</b>	<b>+0.07, 0.83</b>	-0.01
y2017_jes	<b>+0.14, 0.89</b>	+0.09, 0.99	-0.04
y2017_lep	+0.05, 0.99	+0.04, 0.99	-0.01
y2017_lumi	+0.05, 1.00	+0.05, 1.00	-0.02
y2017_prefire	-0.01, 0.99	-0.01, 0.99	+0.01
y2017_pu	+0.05, 1.00	+0.03, 0.99	-0.01
y2017_trig	+0.04, 0.99	+0.04, 0.99	-0.01
y2018_btaghf	+0.13, 0.99	+0.09, 0.99	-0.05
y2018_btaglf	+0.05, 0.99	-0.03, 0.99	-0.04
y2018_fakes_EWK	+0.10, 0.98	-0.01, 0.99	-0.03
y2018_flips	-0.02, 0.99	+0.00, 0.99	+0.01
y2018_fsrvar	+0.00, 0.99	+0.00, 0.99	-0.00
y2018_isr	+0.04, 0.99	+0.01, 0.99	-0.02
y2018_isrvar	+0.00, 0.99	+0.00, 0.99	-0.00
y2018_jer	<b>+0.74, 0.77</b>	<b>+0.50, 0.97</b>	-0.14
y2018_jes	<b>+0.21, 0.67</b>	<b>+0.14, 0.81</b>	-0.07
y2018_lep	+0.05, 0.99	+0.05, 0.99	-0.02
y2018_lumi	+0.08, 0.99	+0.07, 0.99	-0.03
y2018_pu	+0.10, 1.01	+0.06, 0.99	-0.02
y2018_trig	+0.05, 0.99	+0.05, 0.99	-0.03

#### 1407 A.4 Nuisance correlation matrix

1408 Using the combine tool (via `combine -M FitDiagnostics card.txt --robustFit=1`  
 1409 `--saveShapes --saveWithUncertainties --saveOverallShapes --numToysForShapes`  
 1410 `200` ), we can show the correlations between different nuisance parameters in the fit, in Figs. 84-  
 1411 85.

Correlation matrix of fit parameters (cut-based)

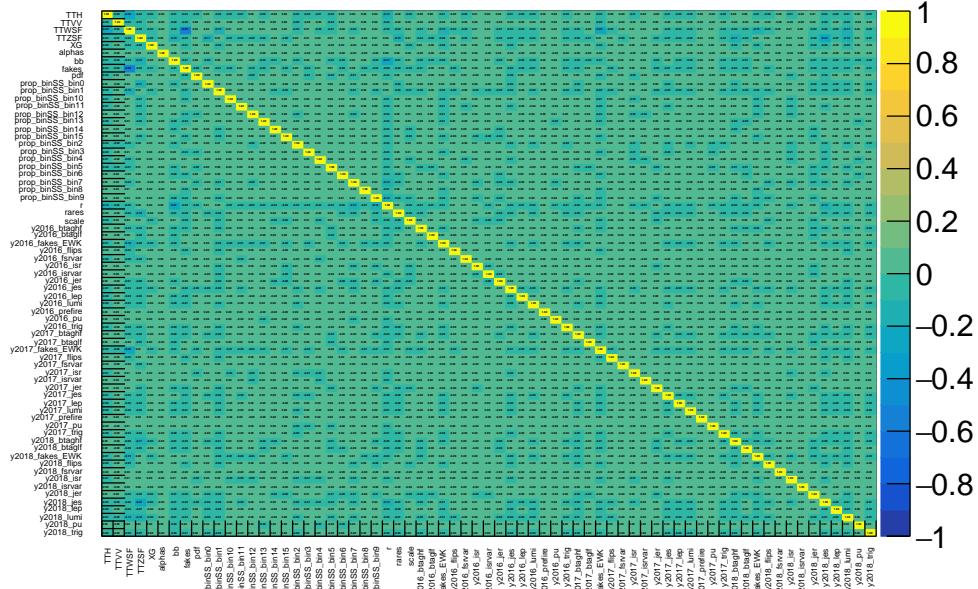


Figure 84: Correlation plot of nuisance-vs-nuisance for the s+b fit for the cut-based regions

Correlation matrix of fit parameters (BDT)

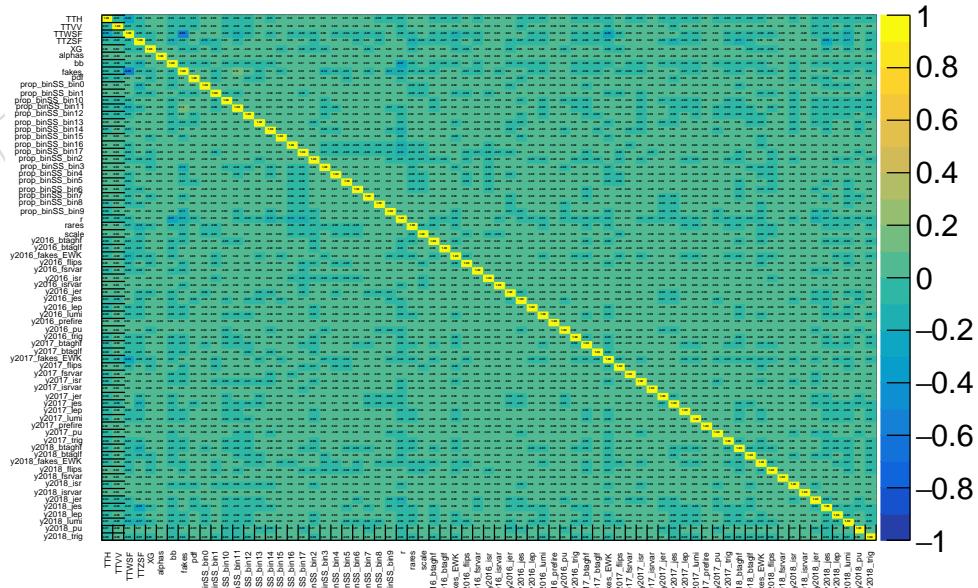


Figure 85: Correlation plot of nuisance-vs-nuisance for the s+b fit for the BDT regions

1412 **A.5 Goodness of fits**

1413 The goodness of fit distributions (using the saturated, Kolmogorov-Smirnov, and Anderson-  
 1414 Darling test statistics) with the signal+background fit to data for the cut-based and BDT anal-  
 1415 yses are shown in Figure 121. We note that the observation is generally within the bulk of the  
 expected distributions for both analyses and all three metrics.

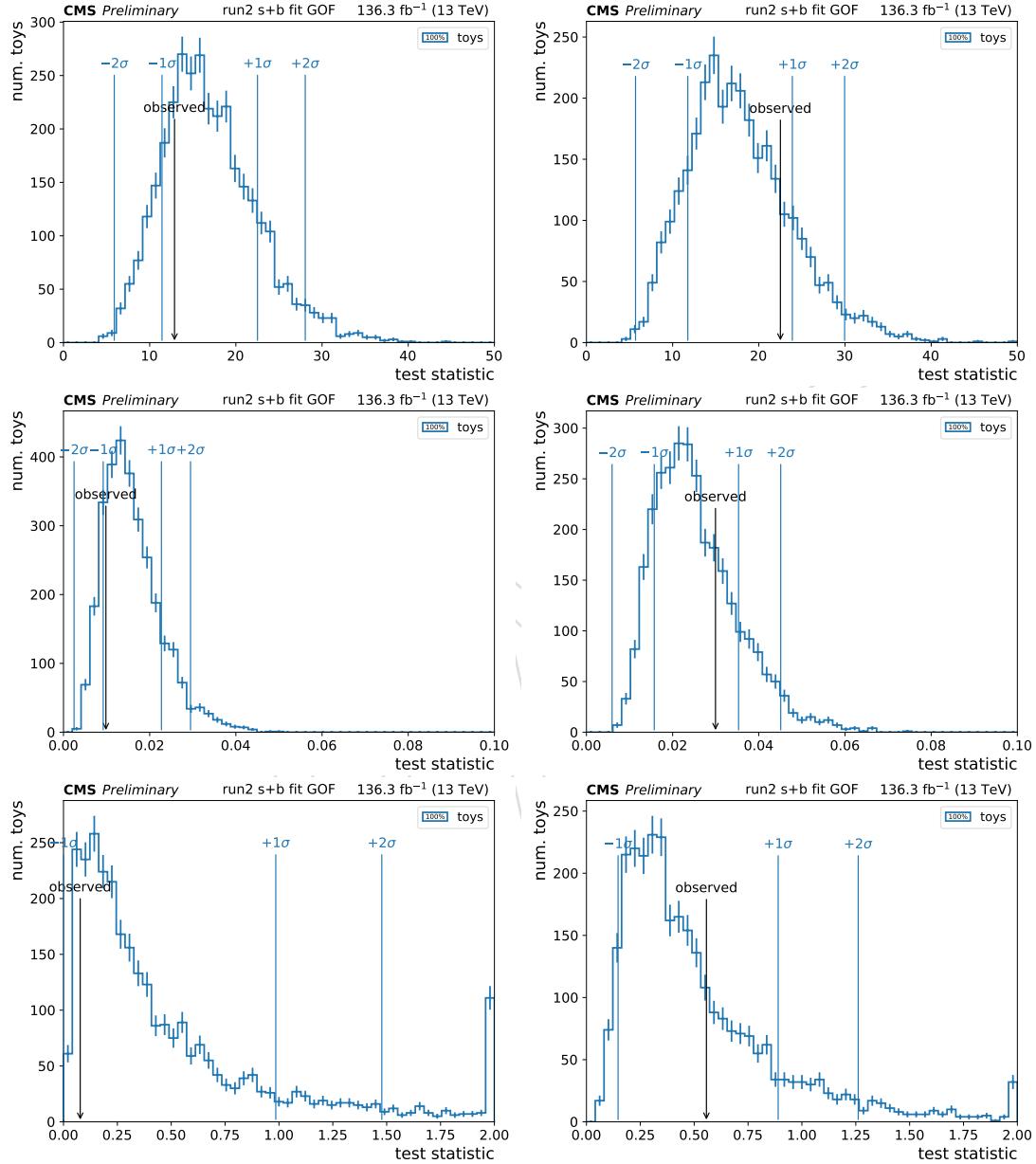


Figure 86: GOF test for the cut-based analysis (left) and BDT-based analysis (right) using the saturated (top), KS (middle), AD (bottom) test statistics.

## 1417 B Shape variations

### 1418 B.1 BDT discriminator

1419 Shape variations for the BDT discriminator for the largest backgrounds and signal, for b-  
 1420 tagging, JES, JER, PU, prefire, ISR/FSR parton shower, and ISR jet reweighting systematics  
 1421 are shown in Figures 87, 88, 89, and 90. Note that complete MC is used with events normalized  
 1422 to  $136.3 \text{ fb}^{-1}$ .

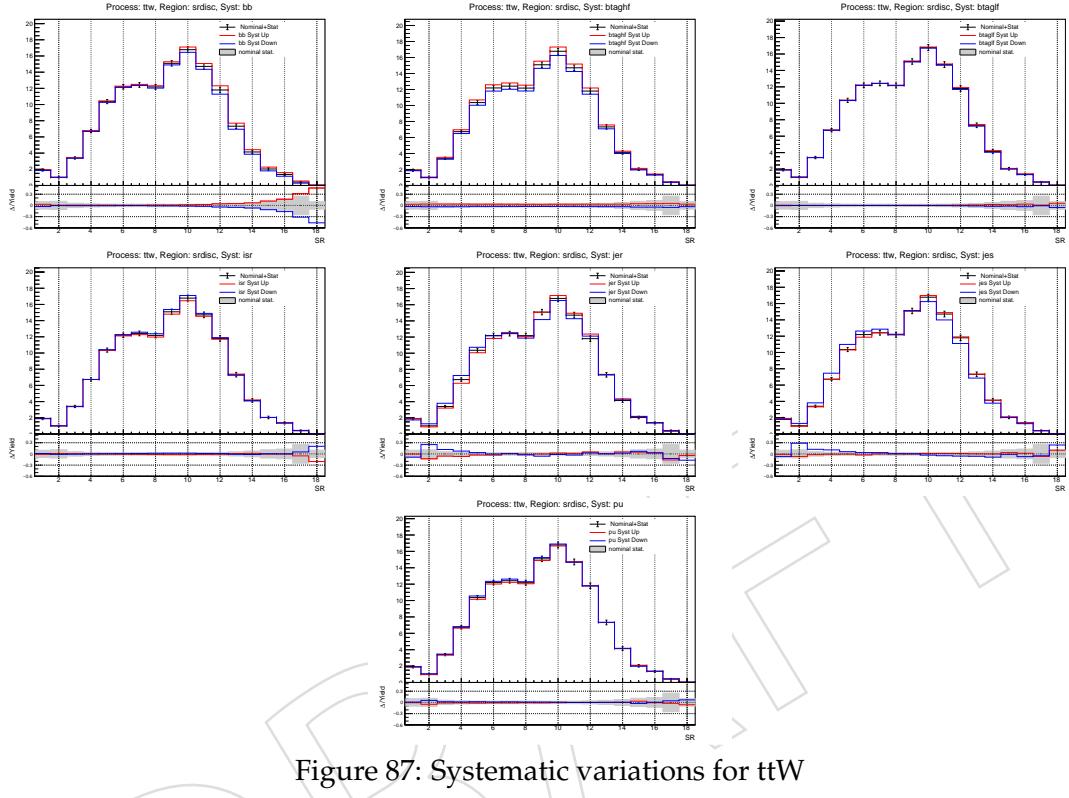
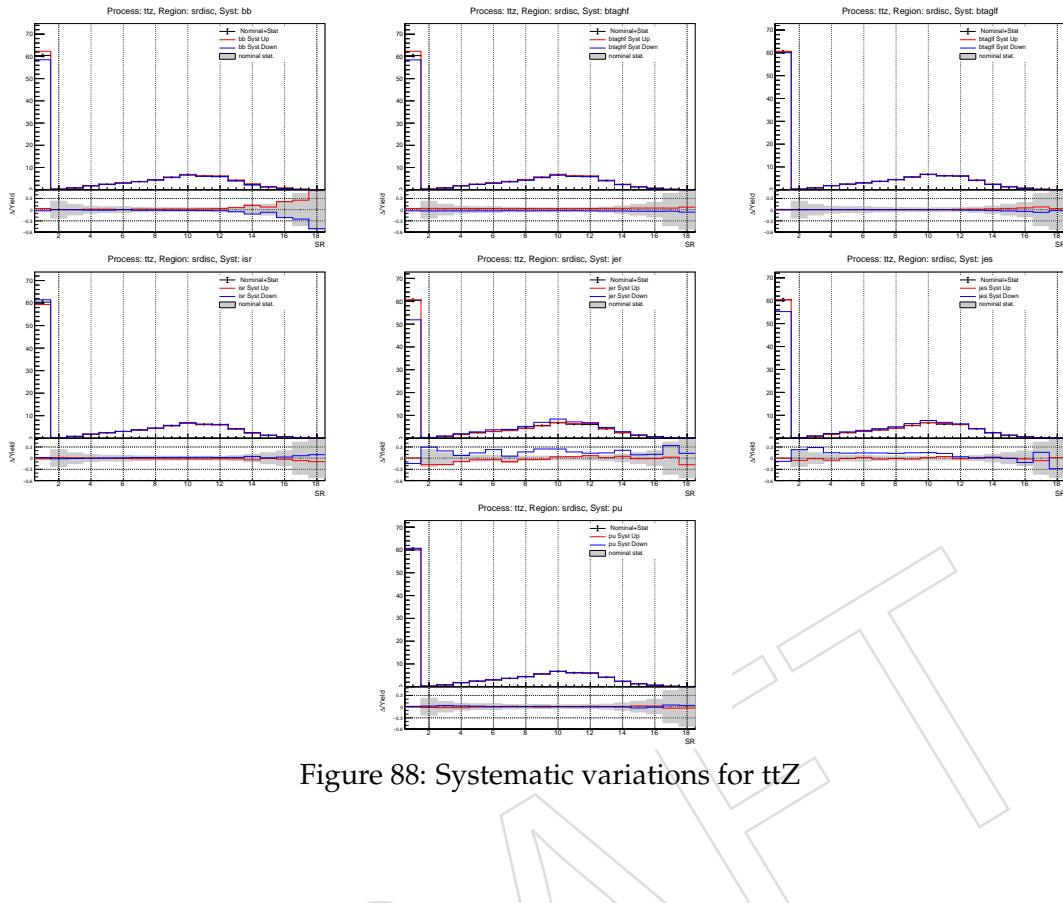
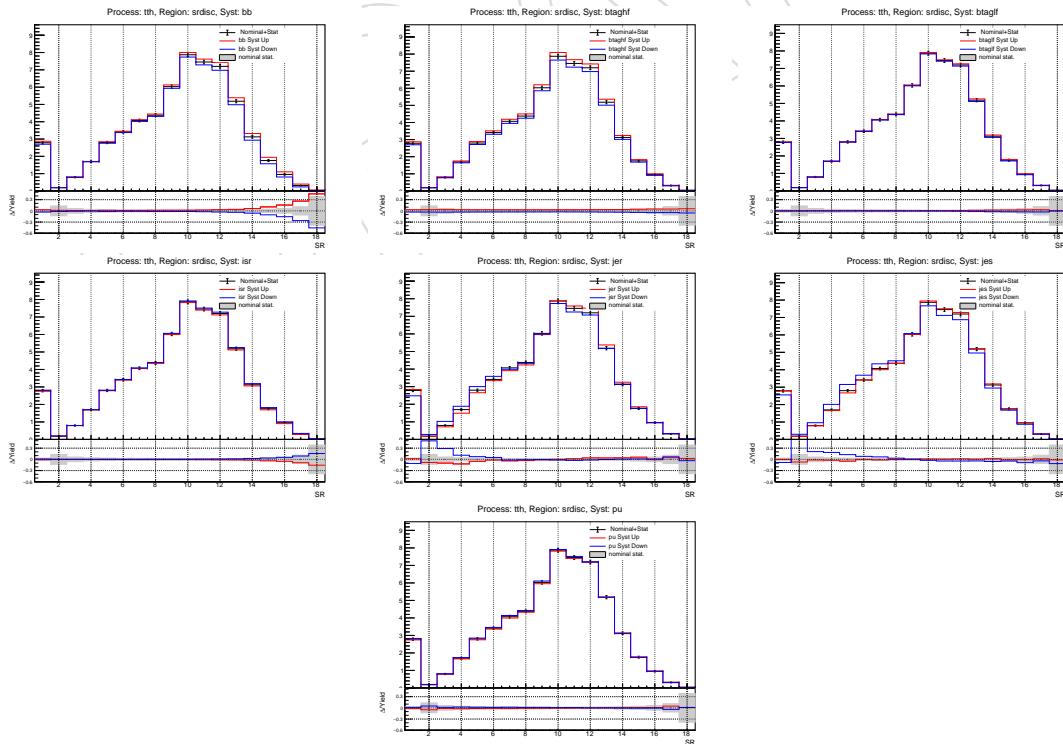
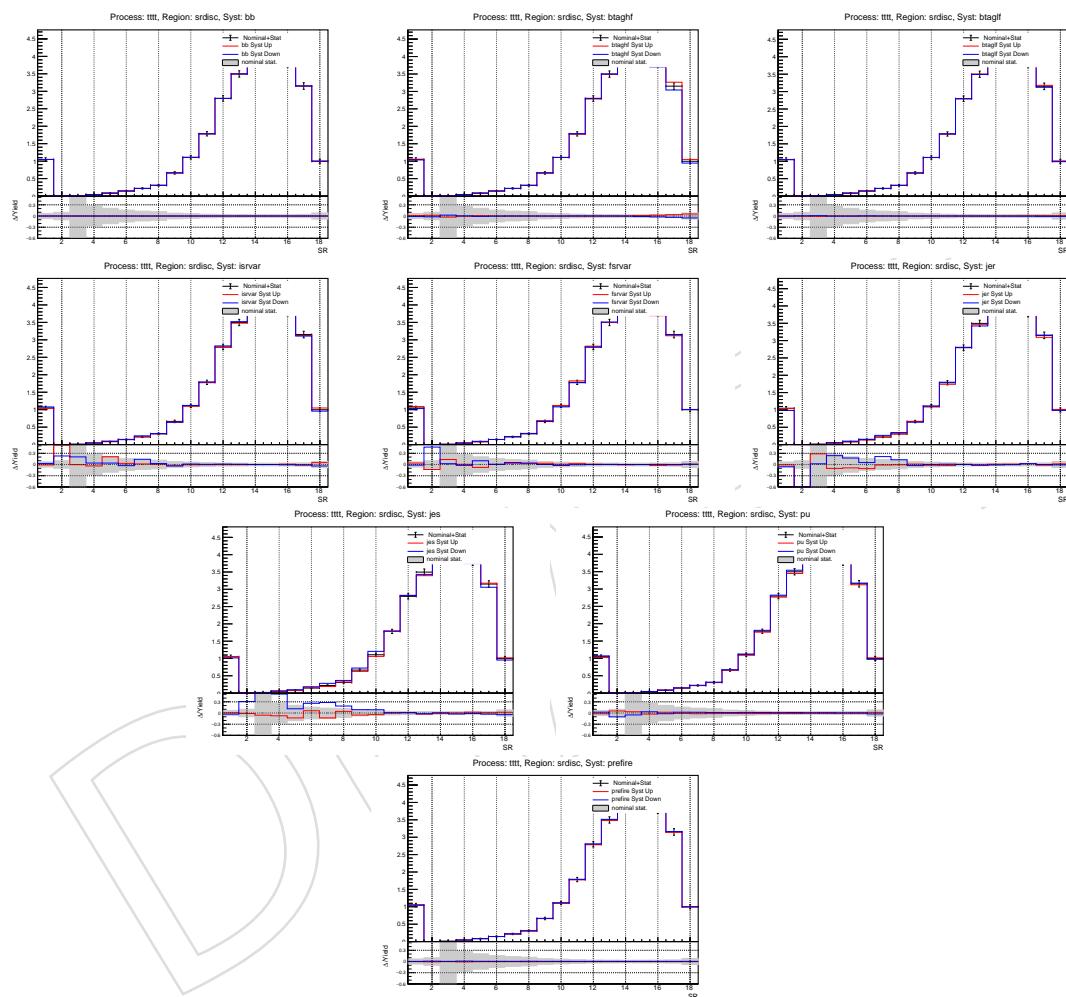


Figure 87: Systematic variations for ttW

Figure 88: Systematic variations for  $t\bar{t}Z$ Figure 89: Systematic variations for  $t\bar{t}H$

Figure 90: Systematic variations for  $t\bar{t}t\bar{t}$  signal

## 1423 C Studies which did not result in analysis changes

### 1424 C.1 Jet and b-jet thresholds

1425 We studied whether the  $t\bar{t}t\bar{t}$  signal significance would improve by lowering the  $N_{\text{jets}}$  and  $N_{\text{b jets}}$   
 1426 counting thresholds from the 2016 ones which were 40 and 25 GeV, respectively. To do this, we  
 1427 used 2016 MC to estimate the expected significance of the analysis with an expected luminosity  
 1428 of  $75 \text{ fb}^{-1}$ , using the 2016+2017 cut based signal region binning. The result, in Figure 91, shows  
 1429 that the current configuration is close to optimal, with only a 3% increase in significance which  
 1430 could be obtained from reducing the  $N_{\text{b jets}}$  threshold from 25 to 20 GeV. Given the minor  
 1431 improvement, and the potential of larger b-tagging scale factors at low  $p_T$  in the 2017 data, we  
 1432 decided to keep the 2016 thresholds.

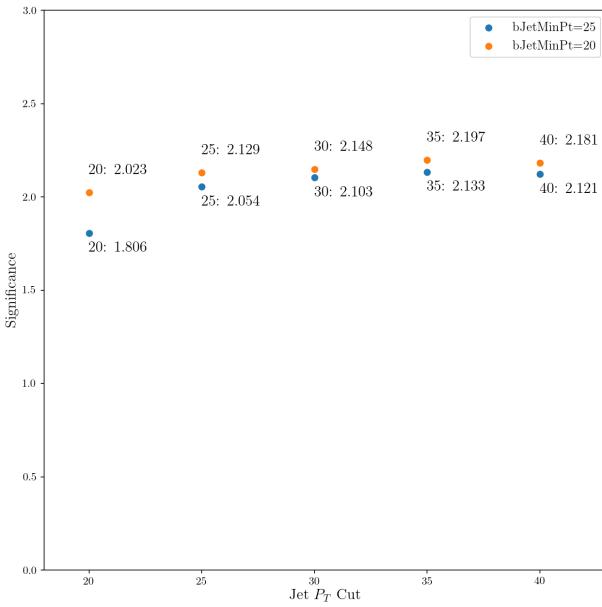


Figure 91: Signal significance for different transverse momentum requirements used in counting  $N_{\text{jets}}$  and  $N_{\text{b jets}}$ .

### 1433 C.2 Taus

1434 We studied whether the  $t\bar{t}t\bar{t}$  signal significance would improve by including taus in addition to  
 1435 electrons and muons. For this study, we used an expected luminosity of  $75 \text{ fb}^{-1}$ , and we plotted  
 1436 the expected yields for events with reconstructed tau leptons, separating the truth-matched and  
 1437 non-truth-matched ones.

1438 We focused on events with only one tau, and the other 1 or 2 leptons being  $e/\mu$ , to avoid  
 1439 large yields of fake taus. Taus were selected with the following identification requirements:  
 1440 decayModeFinding, againstMuonTight3, againstElectronTightMVA6,  
 1441 byTightIsolationMVArun2v1DBdR03oldDMwLT. The last requirement was developed by  
 1442 the ttH analysis to be used in environments with high jet multiplicity [36]. The performance of  
 1443 tau identification is shown in Figure 92 for the  $t\bar{t}t\bar{t}$  signal and the ttW/Z/H background sam-  
 1444 ples, where it is clear that there are large off-diagonal terms due to inefficiency and impurity of

the tau selection. Figure 93 shows the expected yields for each sample as a function of signal region for 1-tau events, split between  $t\bar{t}t\bar{t}$  events where the tau is truth-matched, and the sum of  $t\bar{t}W/Z/H$  events and  $t\bar{t}t\bar{t}$  events with fake taus. The plot shows some potential in the 3-lepton regions ( $eet/\mu\tau/\mu\mu\tau$ ), particularly 14-15-16, which integrate to about 0.5  $t\bar{t}t\bar{t}$  events and 1 background event, but this doesn't account for fake taus from non- $t\bar{t}W/Z/H$  backgrounds. Given these small yields and large backgrounds, we decided not to include signal regions with taus for the 2016+2017 analysis.

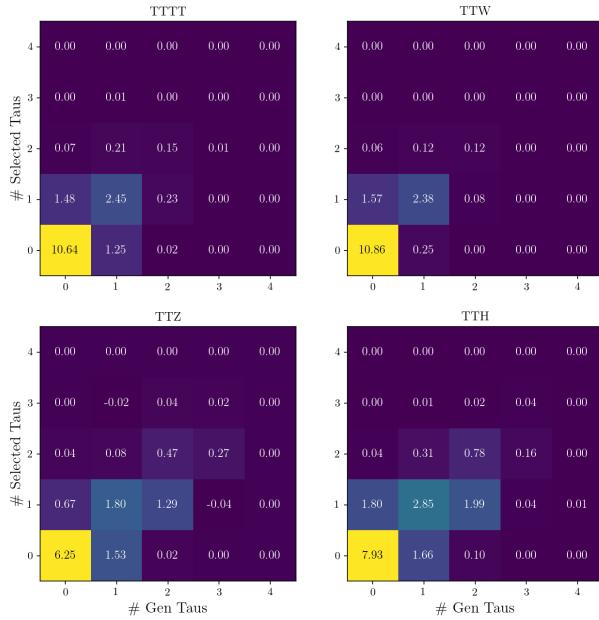


Figure 92: Number of events as a function of the number of true and reconstructed tau leptons, for different samples.

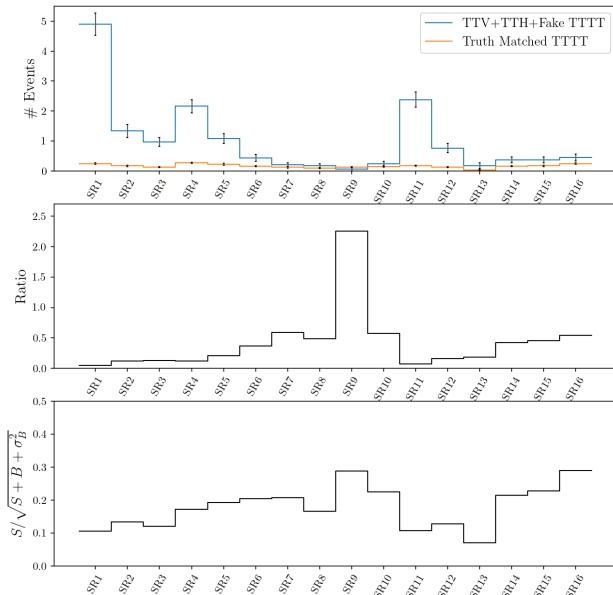


Figure 93: Top: number of events with 1 tau lepton and 1 or 2 e/ $\mu$  leptons, as a function of the signal regions defined in Table 10, with SR11 and higher including events with 3 leptons (ee $\tau$ /e $\mu\tau$ / $\mu\mu\tau$ ), showing truth-matched t $\bar{t}$ t $\bar{t}$  signal (orange) and backgrounds from ttW/Z/H and non-truth-matched t $\bar{t}$ t $\bar{t}$  (blue). Middle and bottom: S/B and and significance per bin, considering only the ttW/Z/H and non-truth-matched t $\bar{t}$ t $\bar{t}$  events as background.

### 1452 C.3 Top-tagging

1453 We explored resolved and merged top tagging as additional discrimination handles.

#### 1454 C.3.1 Resolved

1455 Pairs of 3 jets, where one jet is a b-jet candidate, are fed into a pre-trained BDT from the single  
 1456 lepton and opposite-sign analysis, AN2017-146-v17. The BDT uses 6 inputs (b-tag discriminant,  
 1457 top candidate mass, W candidate mass, the ratio of the top  $p_T$  to the trijet  $p_T$ , and the  $\Delta\phi$   
 1458 between (top,W) and (top,b)). We can find the highest and second highest trijet discriminant in  
 1459 the event. For this analysis, the leading backgrounds can all have a hadronic top, so we focus  
 1460 on the subleading trijet discriminant.

1461 Distributions of this quantity for two jet threshold schemes are shown in Figure 94. The nom-  
 1462 inal analysis uses jet thresholds of 40 GeV and b-tagged jet thresholds of 25 GeV. To explore  
 1463 the possibility that the high jet threshold is reducing the inputs to the BDT, we compare with  
 1464 looser jet thresholds of 20 GeV, 20 GeV, respectively.

1465 In both cases, the ratio of signal to background shows only a slight trend. To test this more  
 1466 quantitatively, the subleading discriminator for the 20,20 scheme was put into a 19+1 variable  
 1467 BDT and was only ranked 9th. Compared to without the variable, the maximum significance  
 1468 only increased by 1%, so we do not pursue such tagging in the analysis. This is not unexpected  
 1469 since the same-sign and multilepton final state has one or two less hadronic tops to tag.

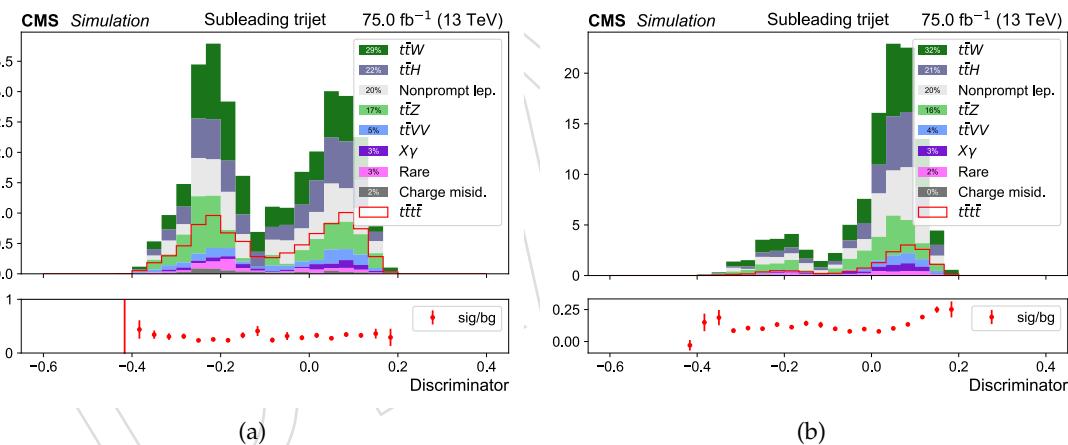


Figure 94: Distributions of subleading trijet discriminator in the signal regions for (jet,b-jet)  $p_T$  thresholds of (40,25) [left] and (20,20) [right].

#### 1470 C.3.2 Merged

1471 Top decay products may become merged through accidental overlaps or boosts. The latter  
 1472 is not a large fraction in this final state due to lack of high MET, however, we explored the  
 1473 DeepAK8 tagger developed by the SUSY Heavy Object Tagging group, which provides nu-  
 1474 mercial probabilities for objects such as top, W, Z, H, etc. for AK8 jets through a deep neural  
 1475 network acting on constituent particle flow candidates.

1476 Following the same strategy as for the resolved top tagger, we look at the leading and sublead-  
 1477 ing top discriminants for events passing the baseline selection for a few leading backgrounds  
 1478 in Figure 95.

Again, while there is a small trend toward higher values for the signal sample, the quantity of these events is minuscule (on the order of a few percent), so this tagging method was not pursued further.

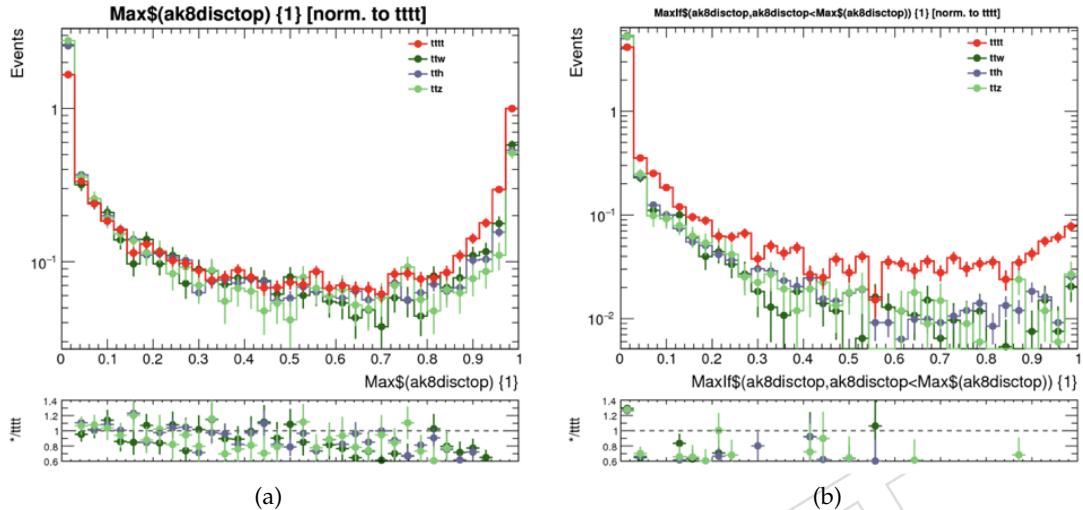


Figure 95: Maximum AK8 top discriminant in the event (left). Subleading AK8 top discriminant in the event (right). Histograms are normalized to the signal (red) cross-section.

1482 **D Impact of the L1 ECAL prefiring(2016/2017) and HEM15/16 loss  
1483 (2018) on the results**

1484 **D.1 L1 prefiring(2016/2017)**

1485 The L1 prefiring issue impacts 2016 and 2017 data collection periods. Due to mistiming and  
1486 trigger rules, events with high  $\eta$  energy deposits can be preferentially lost. The impact of this  
1487 is checked on 80X fastsim T1tttt signal samples. Inefficiency maps to be applied to simulation  
1488 are taken from [https://ncsmith.web.cern.ch/ncsmith/PrefireEfficiencyMaps/Preliminary/Jet\\_L1FinOR\\_eff\\_bx\\_m1\\_looseJet\\_SingleMuon\\_Run2016B-H.pdf](https://ncsmith.web.cern.ch/ncsmith/PrefireEfficiencyMaps/Preliminary/Jet_L1FinOR_eff_bx_m1_looseJet_SingleMuon_Run2016B-H.pdf) (Jet map  
1489 2016B-H) and <https://lathomas.web.cern.ch/lathomas/TSGStuff/L1Prefiring/PrefiringMaps/> (Jet map 2017B-F) and shown in Figure 96. For each event, consider all jets  
1490 and electrons that pick up a non-zero scale factor from the chosen inefficiency map, parameter-  
1491 ized by jet  $p_T$  and  $\eta$  and obtain a multiplicative scale factor (< 1) to apply to MC.  
1492

1493 The procedure matches the tool given in <https://twiki.cern.ch/twiki/bin/view/CMS/L1ECALPrefiringWeightRecipe>.

1494 Figures 97 and 98 show the average SFs as a function of (selected) electron and jet multiplicities per event, separately for 2016 and 2017. The average scale factors are about 2.5% (4%)  
1495 below unity for 2016 (2017). Compared to the central value of the scale factors, the variation  
1496 for increasing jet or electron multiplicity is small. Note that as a function of raw (not analy-  
1497 sis/selected) jets, the trend may be larger, as the analysis jet selection criteria requires  $|\eta| < 2.5$ .  
1498

1499 Scale factors for the prefiring inefficiency will be applied to the 2016 and 2017 MC samples.

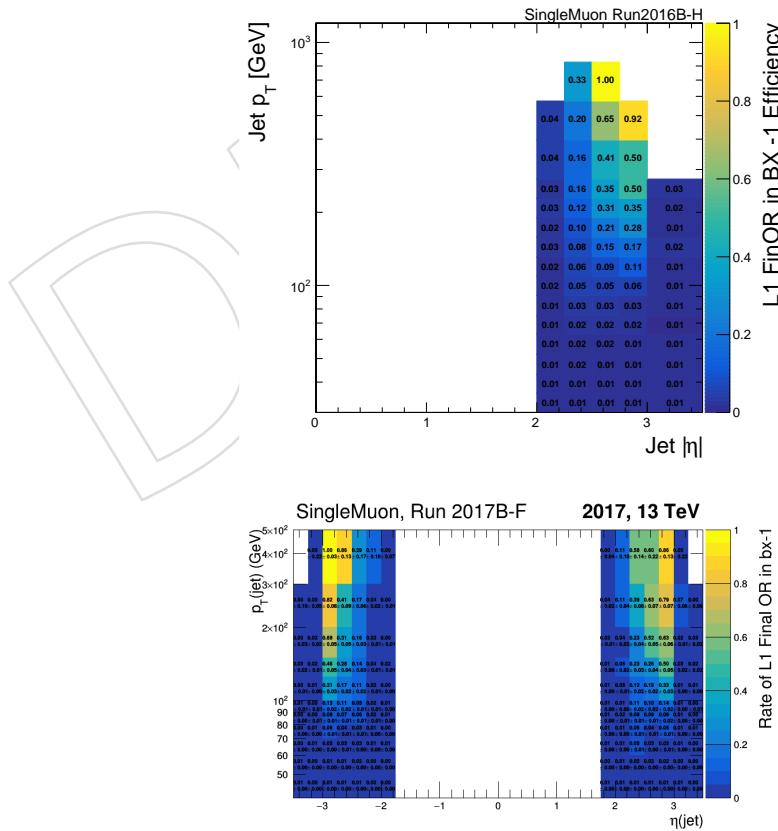


Figure 96: Prefiring inefficiency maps for 2016 (left) and 2017 (right)

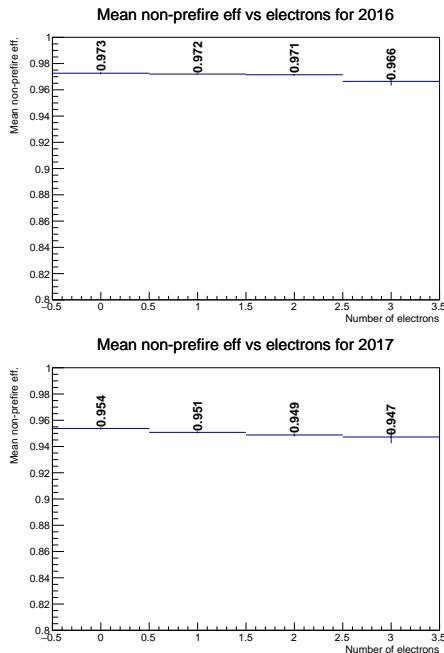


Figure 97: Average prefiring inefficiency as a function of electron multiplicity for 2016 (left) and 2017 (right)

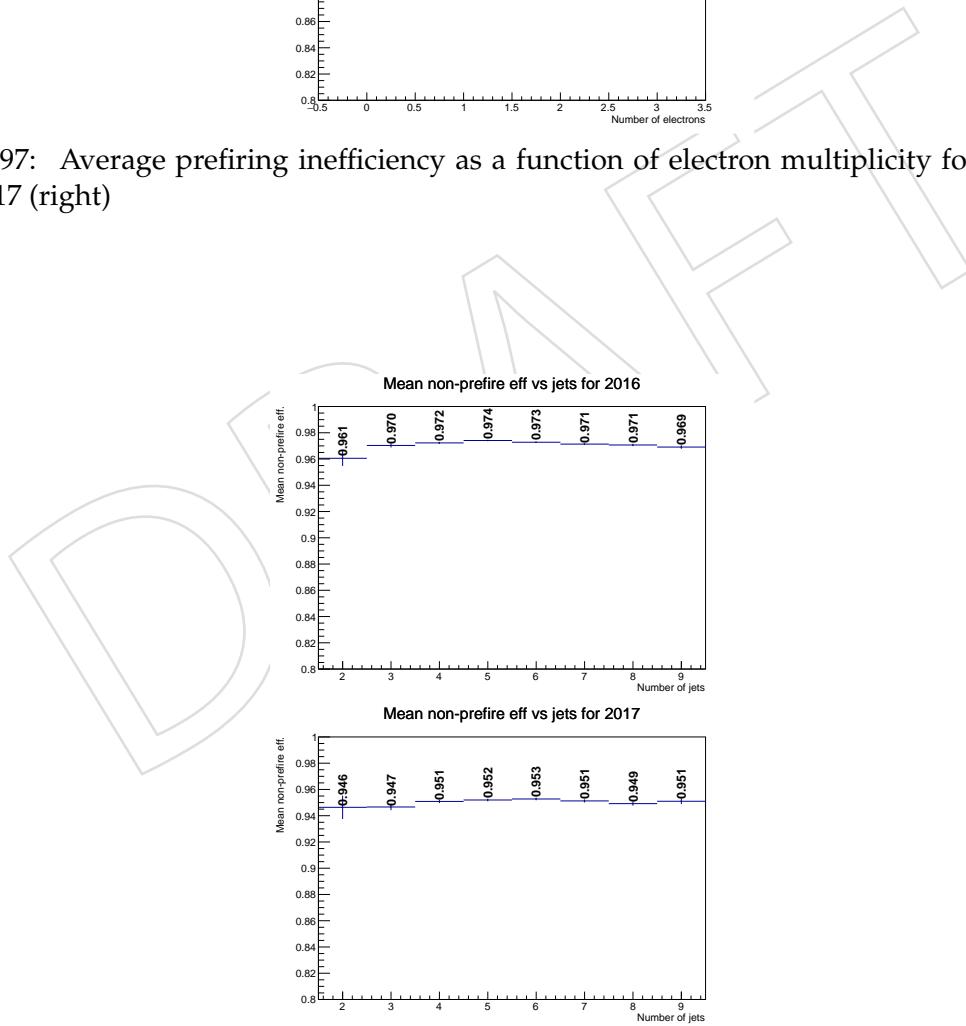


Figure 98: Average prefiring inefficiency as a function of jet multiplicity for 2016 (left) and 2017 (right)

## 1502 D.2 HEM15/16 loss in 2018

1503 Beginning in Run 319077 for the 2018 data collection period, HEM 15 and 16 sectors are switched  
 1504 off due to a power cut issue. This results in a gap of HCAL info for  $\eta \in [-3, 1.3]$  and  $\phi \in$   
 1505  $[-1.57, -0.87]$ . This has the potential to increase electron fakes due to lower H/E and isolation  
 1506 values. The impact was checked in MC and in data.

1507 The relative increase of fake electrons passing tight selection requirements was found to be  
 1508 approximately 12% in MC, while muons showed no increase. Since 23% of the Run2 dataset is  
 1509 impacted by the HEM issue, and considering fakes inclusively ( $e+\mu$ ), this increase corresponds  
 1510 to a 1% overall increase in fakes across the full dataset. Data checks showed similar increases  
 1511 to MC.

### 1512 D.2.1 HEM impact in MC

1513 Dedicated samples were produced to perform an apples-to-apples comparison of the HEM15/16  
 1514 issue. In particular, the dataset

```
1515 /TTToSemiLeptonic_TuneCP5_13TeV-powheg-pythia8
1516 /RunIIISpring18MiniAOD-HEMPremix_100X_upgrade2018_realistic_v10-v3
1517 /MINIAODSIM
```

1518 was produced with HEM sectors disabled. A sample without “HEM” in the dataset name corre-  
 1519 sponds to the nominal HEM “in” case. For both samples, truth-matched fake single lepton effi-  
 1520 ciencies were calculated with nominal analysis tight ID+Iso as the numerator and reco-leptons  
 1521 from miniAOD as the denominator. Leptons were required to have  $p_T > 25\text{ GeV}$ . Labeling  
 1522 HEM “in” and “out” samples as “good” and “bad”, both good and bad samples were checked  
 1523 inclusively and also restricting to the  $\eta\text{-}\phi$  region impacted by the HEM loss. The efficiencies are  
 1524 listed in Table 35. Muons are not impacted, while electrons show a  $(12 \pm 5)\%$  relative increase  
 1525 in fakes inclusively, or a factor of 2 increase in the HEM region specifically.

Table 35: Single lepton efficiencies for good and bad HEM samples.

region	scenario	flavor	sig. eff. (%)	bkg. eff. (%)
all	good	e	$65.8 \pm 0.1$	$0.15 \pm 0.01$
		$\mu$	$84.6 \pm 0.1$	$0.31 \pm 0.01$
	bad	e	$65.9 \pm 0.1$	$0.17 \pm 0.0$
		$\mu$	$84.7 \pm 0.1$	$0.29 \pm 0.01$
hemregion	good	e	$53.9 \pm 0.8$	$0.21 \pm 0.04$
		$\mu$	$80.7 \pm 1.0$	$0.21 \pm 0.06$
	bad	e	$59.7 \pm 0.7$	$0.47 \pm 0.03$
		$\mu$	$82.3 \pm 0.9$	$0.33 \pm 0.07$

### 1526 D.2.2 HEM impact in data

1527 To assess the impact in data, the 2018 dataset was split into two subsections: “before” (runs  
 1528 before 319077, corresponding to  $20.25\text{ fb}^{-1}$ ) and “after” (runs including and after 319077, cor-  
 1529 responding to  $30.73\text{ fb}^{-1}$ ). After normalizing both sections to the same luminosity, plots of  
 1530 lepton flavor ( $\mu\mu, e\mu, ee$ ) for 3 kinematic regions ( $Z \rightarrow \ell\ell$ -dominated,  $t\bar{t} \rightarrow \ell\ell$ -dominated, and  
 1531 the (fake-dominated) tight-loose control region) are shown in Figure 99. While the first two

prompt-enriched regions show a relative increase afterwards, this could be attributed to a variety of data-taking condition differences between the two periods. The tight-loose region shows similar counts for dimuon events, but after/before ratios of  $1.13 \pm 0.08$  for  $e\mu$  and  $1.37 \pm 0.20$  for dielectron events. Such an increase is consistent with the 12% increase found from simulation.  
 Figure 100 shows the lepton  $\phi$  distribution for Z-dominated and tight-loose regions. Muons show no relative changes. When focusing on  $\phi \in [-1.57, -0.87]$ , no increase is observed for prompt electrons, while a  $(59 \pm 23)\%$  increase is seen for the tight-loose electron counts.

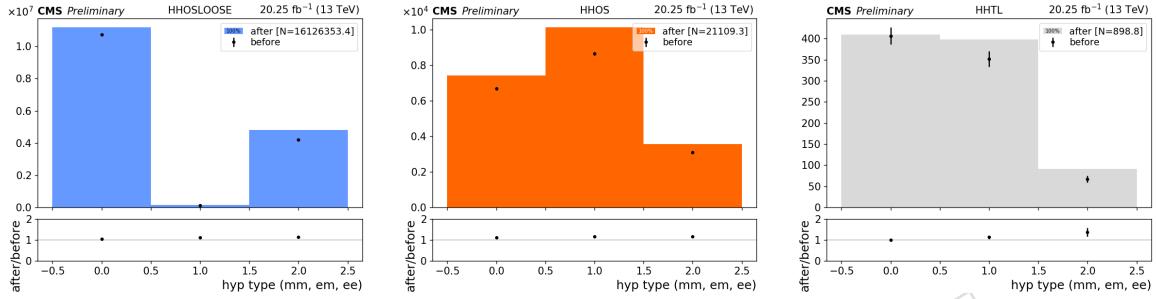


Figure 99: Z-dominated (left),  $t\bar{t}$ -dominated (center), tight-loose (right).

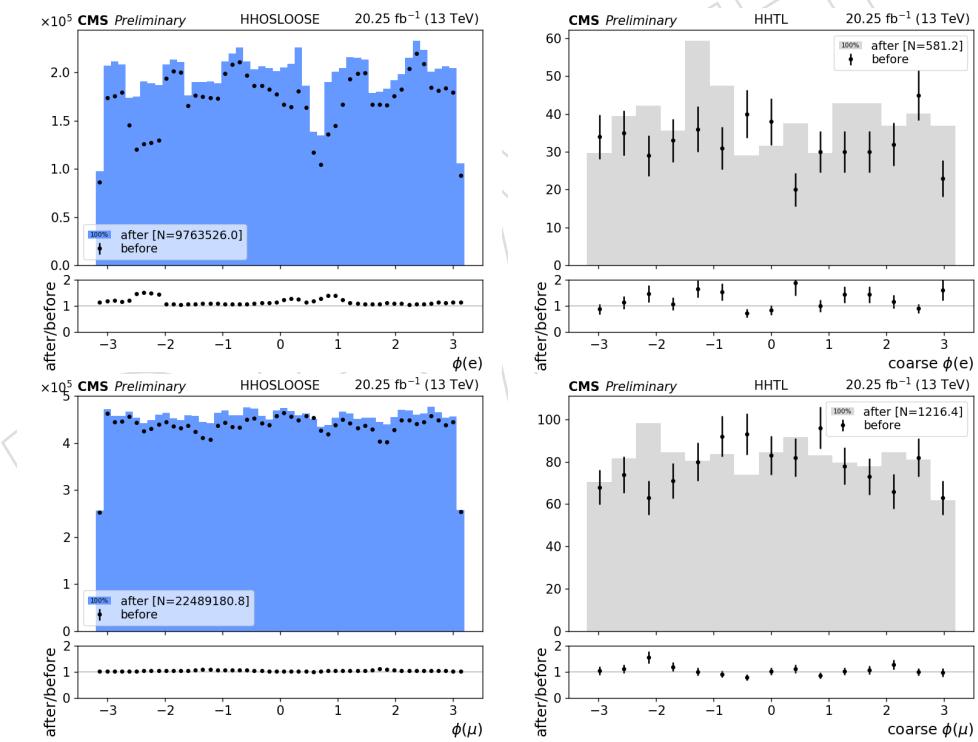


Figure 100: Z-dominated region (left), tight-loose region (right), for electrons (top) and muons (bottom)

## 1539 E Analysis changes

### 1540 E.1 Changes from ANv7

1541 This section details some changes and updates with respect to V7 of this note. Different items  
 1542 are accompanied by the relative change in expected significance from the ANV7 number (cut-  
 1543 based gave  $3.069\sigma$ , and BDT gave  $3.328\sigma$ ).

1544 However, it is worth noting that JECs for 2018 data and MC are not yet available, so 2017 JECs  
 1545 are being used instead. Several of the changes enumerated below and their conclusions are  
 1546 sensitive to JECs.

1547 The compounded effect of these changes brings the cut-based significance to  $2.371\sigma$  and the  
 1548 BDT-based significance to  $2.591\sigma$ .

- 1549 • V32 JEC

- 1550   • With respect to the previous version of JECs (V6), V32 reparameterizes the  
 1551 L1 correction for better stability. While the product of L1,L2L3 remained  
 1552 the same from V6 to V32, the relative ratios did not. Our lepton isolation  
 1553 quantities,  $p_T^{\text{ratio}}$  and  $p_T^{\text{rel}}$  depend on these two corrections separately, not  
 1554 only on their product. Upgrading to V32 caused a per lepton efficiency  
 1555 loss of a few percent, so we re-tuned our isolation working points for 2017  
 1556 and 2018 to approximately match the signal and background efficiencies  
 1557 from V6.
- 1558   • The updated working point values can be found in Table 7.
- 1559   • Relevant discussion in <https://hypernews.cern.ch/HyperNews/CMS/get/JetMET/1891.html>
- 1560   •  $\sigma$  decreases by 4%.

- 1562 • Prefiring SF

- 1563   • The latest recommended prefiring inefficiency maps (from [https://twiki.cern.ch/twiki/bin/view/CMS/SUSRecommendations18#Prefire\\_Issue](https://twiki.cern.ch/twiki/bin/view/CMS/SUSRecommendations18#Prefire_Issue)) were incorporated for 2016 and 2017 MC. Application of the maps  
 1564 scales down four top signal by 2.5% (4.7%) in 2016 (2017) and the ttW  
 1565 background by 1.5% (3.3%) in 2016 (2017).
- 1566   • After application of the inefficiency maps, disagreements in high- $|\eta|$  tails  
 1567 of the opposite-sign control region are largely cured.
- 1568   • For reference, the affected datasets (2016+2017) are only 58% of the Run2  
 1569 dataset by integrated luminosity.
- 1570   •  $\sigma$  decreases by 1%.

- 1573 • Merging 4b signal region bins

- 1574   • Signal regions 8, 9, and 10 (now just 8), corresponding to at least 4 b-  
 1575 tagged jets and at least 5 jets, were merged (removing the jet multiplicity  
 1576 boundaries) to be more conservative when dealing with the tight-loose  
 1577 control region yields for the fake background prediction
- 1578   • The cut-based signal region definition now has 14 signal regions (+2 ttW/ttZ  
 1579 dominated regions, as before).
- 1580   •  $\sigma$  decreases by 4%.

- 1581 • Single card

- 1582   • Card inputs to the HiggsCombine tool are no longer treated as 3 sepa-

1583           rate channels in a simultaneous fit. Instead, yields (including control re-  
 1584           gion yields) are summed into single histograms (2016+2017+2018). Un-  
 1585           correlated nuisances, for example, prefiring uncertainties affecting only  
 1586           2016/2017, are considered as shape variations only on the 2016/2017 com-  
 1587           ponent of the summed nominal histogram. This has the benefit of not be-  
 1588           ing as susceptible to MC statistics issues for each of the 3 years. It also is  
 1589           more conservative and takes advantage of larger control region statistics  
 1590           for the fake prediction.

- 1591           •  $\sigma$  decreases by 5%.

1592           • 2017/2018 MC

- 1593           • The Autumn18 MC campaign is nearly complete, so we have switched to  
 1594           taking relevant 2018 predictions from the correct MC rather than using  
 1595           2017 as a stand-in.
- 1596           • Explicitly, we have all the desired samples in the 2018 campaign except  
 1597           the signal sample is still lacking 25% MC statistics. This does not matter  
 1598           as we don't suffer MC statistics issues for  $t\bar{t}t\bar{t}$ .
- 1599           • In 2017, we are still lacking a very small  $Z+\gamma$  background sample, which  
 1600           we take from the 2018 campaign instead.
- 1601           •  $\sigma$  decreases by 3%.

1602           • 2017/2018 ISR reweighting

- 1603           • For the 2016 dataset, we reweight the  $t\bar{t}W$  and  $t\bar{t}Z$  background predic-  
 1604           tions in order to match the data  $N_J^{\text{ISR/FSR}}$  distribution, as detailed in Sec-  
 1605           tion 2. This had the effect of scaling down high jet multiplicity events,  
 1606           decreasing the background prediction. In 2017 and 2018, however, the  
 1607           Pythia tune was changed to CP5, so we requested dedicated samples with  
 1608           configurations matching  $t\bar{t}W$  and  $t\bar{t}Z$  in order to derive new weights.  
 1609           This results in the opposite trend to 2016, i.e., high jet multiplicity events  
 1610           can get scaled up by almost 30%, increasing the background prediction.
- 1611           • In 2017/2018, the samples for  $t\bar{t}Z$  and  $t\bar{t}W$  with 0 and 1 extra parton,  
 1612           respectively, are

- 1613           • /TT\_DiLept.TuneCP5\_13TeV-amcatnlo-pythia8/
- 1614           • /TTPlus1Jet.DiLept\_TuneCP5\_13TeV-amcatnloFXFX-pythia8/

1615           which are completed by the appropriate processing string for each cam-  
 1616           paign. In the case of 2017, full event statistics are split between a nominal  
 1617           sample and an extension (“ext1”) sample.

- 1618           •  $\sigma$  decreases by 2%.

1619           • 2018 b-tag WPs

- 1620           • The 2018 b-tag medium WP for DeepCSV preserves the b tagging and  
 1621           light mis-tagging efficiencies compared to 2017. However, it appears the  
 1622           charm mis-tagging efficiency relatively increased by about 15%. In the 3  
 1623           and 4 b-tag bins,  $t\bar{t}W$  can enter if there is additional heavy flavor contri-  
 1624           bution (e.g.,  $t\bar{t}W + bb$  via gluon splitting), or through mistags of W decays  
 1625           containing a charm quark.
- 1626           • Relevant discussion in <https://hypernews.cern.ch/HyperNews/CMS/get/btag/1637.html>
- 1627           •  $\sigma$  decreases by 3%.

1629 • ttW+bb, ttZ+bb, ttH+bb scaling

- 1630   • As a response to a 7 Dec 2018 question from Otto about Figs 12, 13 (<https://twiki.cern.ch/twiki/bin/viewauth/CMS/TOP18003>), we showed  
1631   that scaling the gluon splitting (events with extra  $b\bar{b}$ ) components up by  
1632   the systematic uncertainty of 1.35 brings better agreement in the b-tag  
1633   multiplicity distribution.
- 1634   • We now revert to the procedure of the previous iteration of this analysis  
1635   for 2016 by scaling  $(t\bar{t}W/t\bar{t}Z/t\bar{t}H)+b\bar{b}$  up by  $1.7 \pm 0.6$  with a systematic  
1636   uncertainty corresponding to the specified error, from the result of TOP-  
1637   16-010, which found a Data/MC discrepancy in the  $ttbb/ttjj$  ratio.
- 1638   •  $\sigma$  decreases by 6%.
- 1639

DRAFT

## 1640 F Unblinding of 2016 dataset

1641 Following the pre-approval talk for the analysis on Feb. 15, 2019, the analysis was unblinded  
 1642 using the 2016 dataset. The yields and results are shown in Section F.1 with consistency checks  
 1643 of cut-based and BDT results in Section F.2. We include the results of various statistical checks,  
 1644 including nuisance pulls (Section F.3), and nuisance impacts (Section F.4).

### 1645 F.1 Yields and results

1646 Plots for the unblinded prefit and postfit event yields for the 2016 data, with a total luminosity  
 1647 of  $35.9 \text{ fb}^{-1}$ , are shown for both the cut-based and BDT based analysis in Figure 101. Numerical  
 1648 yields are also tabulated in Table 36 (prefit cut-based analysis), Table 37 (postfit cut-based  
 1649 analysis), Table 38 (prefit BDT analysis), Table 39 (postfit BDT analysis).

1650 A binned likelihood fit is performed over signal regions using the Higgs Combine tool [22];  
 1651 exclusion limits at 95% CL are calculated with the Asymptotic CLs method.

1652 With the cut-based analysis, an observed (expected) upper limit on the production cross section  
 1653 of  $33.67 \text{ fb}$  ( $20.48^{+11.18}_{-6.83} \text{ fb}$ ), assuming the signal process does not exist. The observed (expected)  
 1654 significance is  $1.012$  ( $1.374$ ) standard deviations, corresponding to a measured signal strength  
 1655 parameter of  $0.812^{+1.054}_{-0.800}$ .

1656 With the BDT analysis, an observed (expected) upper limit on the production cross section of  
 1657  $42.59 \text{ fb}$  ( $19.36^{+10.11}_{-6.46} \text{ fb}$ ), assuming the signal process does not exist. The observed (expected)  
 1658 significance is  $1.812$  ( $1.424$ ) standard deviations, corresponding to a measured signal strength  
 1659 parameter of  $1.538^{+1.137}_{-0.935}$ .

### 1660 F.2 Cut-based and BDT consistency checks

1661 One noteworthy excess for the cut-based is SR4, which has 8 observed events with about  $2.2 \pm$   
 1662 0.5 predicted background. This is a 2-lepton, 5 jet, 3 b-tagged jet bin. In the 2016 result, this  
 1663 corresponded to part of oldSR4 (2-lepton, 5 or 6 jets, 3 b-tagged jets), and had 8 observed events.  
 1664 With the new binning which splits oldSR4 into SR4 and SR5, the data events remain in the 5 jet  
 1665 bin, SR4, leaving the slightly higher s/b bin SR5 with 0 data.

1666 It is also worth pointing out the mildness of the disagreement between cut-based and BDT  
 1667 significances (both have differing directions for observed significance with respect to their ex-  
 1668 pected significance). We have checked with toy pseudo-datasets that the probability of this  
 1669 happening is close to 50%.

1670 The asimov-like toys were constructed in a way that would preserve the correlation between  
 1671 the cut-based and BDT significance results. The procedure for creating one toy is to consider all  
 1672 background events for the analysis, storing the cut-based and BDT SR bin number and event  
 1673 weight. For each background/signal process, calculate the total yield  $y$  in the cut-based/BDT  
 1674 signal regions and sample a number  $N$  from a Poisson distribution with parameter  $y$ . Next,  
 1675 draw  $N$  random events from the set of events for a particular process with relative probability  
 1676 given by the absolute value of the event weight (so there is a slight bias due to negatively  
 1677 weighted events). These  $N$  sampled events then are used to fill a histogram for the cut-based  
 1678 and BDT shapes. In this way, we create  $\mu \cdot s + b$  toys with  $\mu = 1$ , which is then an input to  
 1679 HiggsCombine to produce a pair of significances (cut-based  $\sigma$  and BDT  $\sigma$ ).

1680 Several thousand pairs of values from the toy pseudo-datasets are binned into 2D histogram  
 1681 with the 1,2, and  $3\sigma$  contours overlaid. The threshold for the contours is calculated with a  
 1682 2-sample  $\chi^2$  metric. Overlaying the observed significances from above, we find it lies within

1683 the  $0.7\sigma$  contour with a p-value of 0.5, as shown in Figure 108. As described in the caption, the  
 1684 observed point is “corrected” (only within the scope of this toy study) before plotting it.

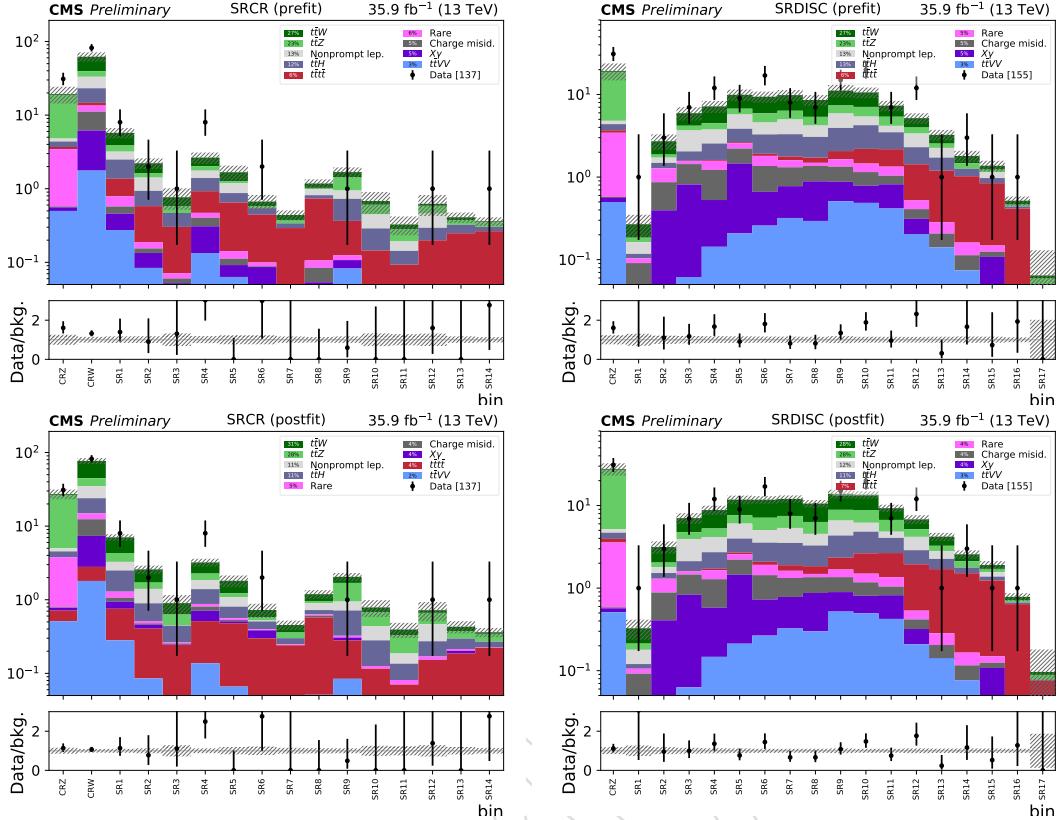


Figure 101: 2016: data yields compared to prefit (top) and postfit (bottom) SM predictions in CRZ, CRW, and the cut-based signal regions (left), and CRZ with the BDT signal regions (right).  
**Note that we plan to retrain the BDT with latest JECs/corrections for 2018 (with identical settings)**

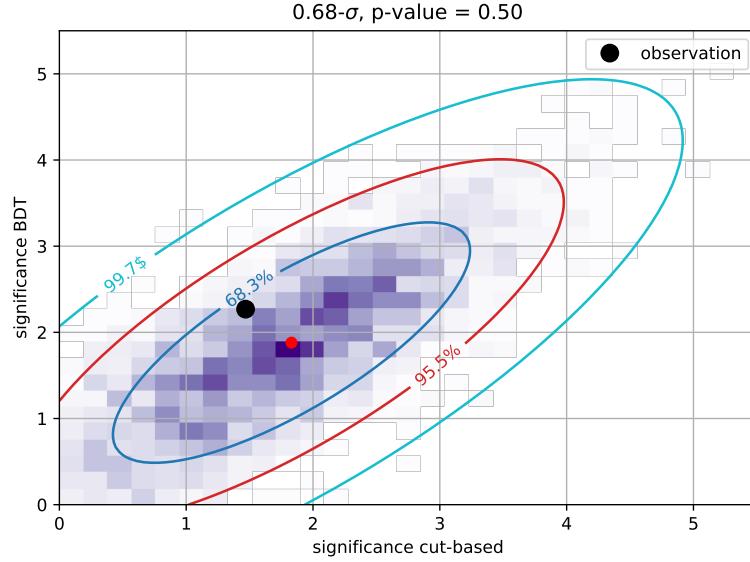


Figure 102: Comparison of correlated expected significances between cut-based and BDT fits with observed point overlaid. Because significances from toy datasets are used, we get a correction factor for expected BDT/cut-based significances with the full fit and the mean of the significances from the toy datasets. This correction factor is then applied to the observed point such that the expected significance values from the full analysis would match the mean values from the toy pseudo-datasets.

Table 36: Prefit event yields in cut-based regions for 2016.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	0.50± 0.18	14.01± 4.84	0.69± 0.17	0.50± 0.06	0.06± 0.01	2.87± 0.58	0.02± 0.00	0.42± 0.18	19.06± 4.95	31	0.25± 0.03
CRW	22.56± 7.66	6.31± 2.14	8.44± 2.07	1.76± 0.21	4.41± 0.63	2.33± 0.48	4.96± 0.95	10.05± 3.79	60.82± 9.06	82	1.23± 0.11
SR1	1.87± 0.68	0.68± 0.24	1.12± 0.30	0.27± 0.04	0.18± 0.08	0.21± 0.05	0.12± 0.02	0.70± 0.32	5.16± 0.93	8	0.58± 0.06
SR2	0.61± 0.24	0.20± 0.09	0.35± 0.09	0.08± 0.02	0.05± 0.02	0.03± 0.01	0.02± 0.00	0.48± 0.23	1.83± 0.37	2	0.39± 0.03
SR3	0.18± 0.10	0.11± 0.06	0.17± 0.06	0.05± 0.01	0.00± 0.01	0.01± 0.00	0.01± 0.00	0.00± 0.17	0.53± 0.25	1	0.23± 0.05
SR4	0.63± 0.24	0.26± 0.11	0.47± 0.12	0.13± 0.02	0.17± 0.07	0.07± 0.01	0.09± 0.02	0.36± 0.21	2.19± 0.46	8	0.46± 0.06
SR5	0.40± 0.18	0.09± 0.03	0.22± 0.07	0.06± 0.02	0.03± 0.01	0.03± 0.01	0.02± 0.00	0.31± 0.30	1.17± 0.39	0	0.51± 0.05
SR6	0.09± 0.06	0.03± 0.04	0.10± 0.04	0.01± 0.01	0.07± 0.05	0.01± 0.00	0.01± 0.00	0.00± 0.07	0.32± 0.16	2	0.34± 0.04
SR7	0.06± 0.04	0.04± 0.02	0.05± 0.02	0.01± 0.00	0.00± 0.00	0.01± 0.00	0.00± 0.00	0.00± 0.01	0.16± 0.07	0	0.28± 0.08
SR8	0.15± 0.08	0.04± 0.02	0.08± 0.03	0.05± 0.01	0.00± 0.00	0.02± 0.00	0.03± 0.01	0.17± 0.11	0.55± 0.17	0	0.63± 0.06
SR9	0.25± 0.09	0.49± 0.17	0.36± 0.09	0.08± 0.01	0.02± 0.00	0.02± 0.01	0.00± 0.00	0.22± 0.11	1.44± 0.27	1	0.24± 0.07
SR10	0.08± 0.03	0.16± 0.07	0.14± 0.04	0.02± 0.00	0.00± 0.00	0.01± 0.01	0.00± 0.00	0.16± 0.19	0.57± 0.22	0	0.11± 0.02
SR11	0.04± 0.02	0.09± 0.05	0.05± 0.02	0.01± 0.00	0.00± 0.00	0.01± 0.01	0.00± 0.00	0.05± 0.04	0.26± 0.09	0	0.07± 0.02
SR12	0.04± 0.02	0.13± 0.05	0.09± 0.03	0.02± 0.00	0.00± 0.00	0.02± 0.00	0.00± 0.00	0.17± 0.18	0.46± 0.20	1	0.17± 0.02
SR13	0.04± 0.02	0.05± 0.03	0.08± 0.02	0.01± 0.00	0.02± 0.00	0.01± 0.00	0.00± 0.00	0.00± 0.02	0.20± 0.07	0	0.21± 0.02
SR14	0.01± 0.02	0.05± 0.03	0.04± 0.02	0.01± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.01	0.11± 0.05	1	0.25± 0.04

Table 37: Postfit event yields in cut-based regions for 2016.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	0.72± 0.18	21.58± 4.13	0.75± 0.17	0.51± 0.06	0.06± 0.01	3.00± 0.59	0.02± 0.00	0.47± 0.21	27.09± 4.12	31	0.20± 0.22
CRW	32.12± 7.40	9.64± 1.87	9.08± 2.03	1.80± 0.22	4.57± 0.66	2.45± 0.50	5.03± 0.84	11.13± 4.44	75.82± 7.26	82	0.99± 1.03
SR1	2.72± 0.68	1.04± 0.21	1.21± 0.27	0.28± 0.03	0.19± 0.06	0.22± 0.05	0.12± 0.02	0.77± 0.35	6.56± 0.71	8	0.47± 0.52
SR2	0.87± 0.23	0.31± 0.07	0.38± 0.09	0.09± 0.01	0.05± 0.01	0.04± 0.01	0.02± 0.00	0.51± 0.23	2.26± 0.33	2	0.32± 0.30
SR3	0.27± 0.10	0.18± 0.05	0.18± 0.05	0.05± 0.01	0.00± 0.00	0.01± 0.00	0.01± 0.00	0.00± 0.18	0.71± 0.26	1	0.19± 0.21
SR4	0.94± 0.25	0.42± 0.09	0.52± 0.12	0.14± 0.02	0.20± 0.05	0.07± 0.01	0.10± 0.02	0.46± 0.24	2.83± 0.40	8	0.37± 0.35
SR5	0.59± 0.16	0.13± 0.03	0.24± 0.07	0.07± 0.01	0.03± 0.00	0.03± 0.01	0.02± 0.00	0.28± 0.26	1.39± 0.33	0	0.41± 0.38
SR6	0.14± 0.06	0.06± 0.03	0.11± 0.03	0.01± 0.01	0.08± 0.05	0.01± 0.00	0.00± 0.00	0.01± 0.08	0.44± 0.16	2	0.29± 0.26
SR7	0.09± 0.05	0.06± 0.02	0.05± 0.02	0.01± 0.00	0.00± 0.00	0.01± 0.00	0.00± 0.00	0.00± 0.01	0.22± 0.07	0	0.23± 0.19
SR8	0.23± 0.08	0.06± 0.02	0.09± 0.03	0.05± 0.01	0.00± 0.00	0.02± 0.00	0.03± 0.01	0.18± 0.11	0.67± 0.16	0	0.52± 0.50
SR9	0.35± 0.09	0.75± 0.16	0.39± 0.09	0.08± 0.01	0.03± 0.00	0.02± 0.01	0.00± 0.00	0.23± 0.13	1.85± 0.26	1	0.20± 0.16
SR10	0.11± 0.03	0.24± 0.06	0.15± 0.04	0.02± 0.00	0.00± 0.00	0.01± 0.01	0.00± 0.00	0.15± 0.17	0.69± 0.20	0	0.09± 0.11
SR11	0.06± 0.02	0.15± 0.04	0.05± 0.02	0.01± 0.00	0.00± 0.00	0.01± 0.00	0.00± 0.00	0.05± 0.05	0.34± 0.09	0	0.06± 0.07
SR12	0.06± 0.02	0.20± 0.05	0.10± 0.02	0.02± 0.00	0.00± 0.00	0.02± 0.00	0.00± 0.00	0.19± 0.19	0.58± 0.21	1	0.14± 0.14
SR13	0.05± 0.03	0.08± 0.04	0.08± 0.02	0.01± 0.00	0.02± 0.00	0.01± 0.00	0.00± 0.00	0.00± 0.02	0.26± 0.08	0	0.17± 0.16
SR14	0.02± 0.02	0.07± 0.03	0.04± 0.01	0.01± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.01	0.15± 0.05	1	0.21± 0.21

Table 38: Prefit event yields in BDT regions for 2016.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	0.50± 0.17	14.01± 4.41	0.69± 0.17	0.50± 0.06	0.06± 0.01	2.87± 0.56	0.02± 0.00	0.42± 0.17	19.06± 4.49	31	0.25± 0.03
SR1	0.08± 0.04	0.02± 0.01	0.01± 0.01	0.00± 0.00	0.02± 0.01	0.01± 0.01	0.07± 0.01	0.05± 0.05	0.27± 0.08	1	0.00± 0.00
SR2	0.84± 0.30	0.16± 0.06	0.21± 0.05	0.02± 0.00	0.38± 0.19	0.41± 0.18	0.47± 0.10	0.23± 0.15	2.71± 0.58	3	0.00± 0.01
SR3	1.75± 0.65	0.52± 0.19	0.47± 0.15	0.06± 0.01	0.75± 0.21	0.13± 0.05	0.62± 0.13	1.62± 0.69	5.93± 1.09	7	0.01± 0.01
SR4	2.70± 0.89	0.72± 0.32	0.94± 0.24	0.14± 0.02	0.38± 0.19	0.35± 0.10	0.69± 0.14	1.21± 0.57	7.14± 1.20	12	0.04± 0.01
SR5	3.19± 1.07	1.00± 0.45	1.20± 0.31	0.21± 0.03	1.25± 0.29	0.36± 0.20	0.75± 0.15	1.88± 0.73	9.84± 1.60	9	0.08± 0.01
SR6	3.49± 1.15	1.33± 0.45	1.37± 0.35	0.26± 0.03	0.40± 0.21	0.44± 0.12	0.70± 0.14	1.29± 0.48	9.29± 1.45	17	0.11± 0.05
SR7	3.47± 1.14	1.37± 0.48	1.53± 0.37	0.32± 0.04	0.46± 0.15	0.29± 0.06	0.54± 0.11	1.68± 0.70	9.65± 1.50	8	0.18± 0.03
SR8	3.22± 1.06	1.16± 0.43	1.35± 0.34	0.29± 0.06	0.59± 0.15	0.12± 0.10	0.48± 0.10	1.14± 0.48	8.35± 1.26	7	0.23± 0.02
SR9	3.78± 1.40	1.48± 0.51	1.87± 0.48	0.51± 0.06	0.37± 0.22	0.31± 0.07	0.46± 0.09	1.95± 0.92	10.72± 1.91	15	0.42± 0.04
SR10	3.52± 1.21	1.68± 0.56	2.05± 0.52	0.49± 0.06	0.30± 0.07	0.31± 0.06	0.36± 0.07	1.15± 0.52	9.85± 1.56	20	0.74± 0.07
SR11	1.94± 0.70	1.21± 0.41	1.31± 0.36	0.42± 0.06	0.40± 0.09	0.30± 0.08	0.22± 0.04	0.69± 0.39	6.48± 1.07	7	0.84± 0.05
SR12	1.19± 0.46	0.62± 0.22	0.87± 0.23	0.20± 0.03	0.11± 0.02	0.11± 0.03	0.10± 0.02	1.07± 0.45	4.26± 0.75	12	0.91± 0.06
SR13	0.70± 0.28	0.29± 0.11	0.52± 0.18	0.14± 0.03	0.00± 0.03	0.08± 0.02	0.07± 0.01	0.54± 0.32	2.33± 0.48	1	0.91± 0.04
SR14	0.34± 0.16	0.20± 0.08	0.24± 0.08	0.07± 0.01	0.00± 0.03	0.05± 0.01	0.04± 0.01	0.00± 0.11	0.94± 0.28	3	0.86± 0.07
SR15	0.14± 0.08	0.09± 0.04	0.13± 0.05	0.04± 0.01	0.06± 0.06	0.02± 0.01	0.02± 0.00	0.17± 0.12	0.68± 0.21	1	0.69± 0.06
SR16	0.05± 0.03	0.02± 0.02	0.03± 0.01	0.01± 0.00	0.00± 0.00	0.01± 0.00	0.00± 0.00	0.00± 0.03	0.12± 0.07	1	0.39± 0.05
SR17	0.01± 0.01	0.01± 0.01	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.06	0.02± 0.07	0	0.05± 0.01

Table 39: Postfit event yields in BDT regions for 2016.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	0.66± 0.19	21.93± 4.77	0.74± 0.17	0.51± 0.06	0.06± 0.01	2.98± 0.55	0.02± 0.00	0.47± 0.18	27.37± 4.71	31	0.38± 0.24
SR1	0.11± 0.04	0.03± 0.01	0.01± 0.01	0.00± 0.00	0.02± 0.01	0.01± 0.03	0.07± 0.01	0.06± 0.06	0.33± 0.09	1	0.00± 0.00
SR2	1.10± 0.33	0.24± 0.08	0.22± 0.05	0.02± 0.00	0.39± 0.18	0.41± 0.18	0.48± 0.09	0.27± 0.17	3.14± 0.54	3	0.00± 0.01
SR3	2.30± 0.71	0.80± 0.21	0.51± 0.14	0.06± 0.01	0.77± 0.17	0.13± 0.05	0.63± 0.11	1.82± 0.80	7.02± 1.04	7	0.02± 0.02
SR4	3.54± 0.97	1.11± 0.40	1.01± 0.22	0.15± 0.02	0.43± 0.18	0.37± 0.09	0.70± 0.13	1.44± 0.71	8.75± 1.15	12	0.06± 0.04
SR5	4.18± 1.18	1.56± 0.57	1.28± 0.28	0.21± 0.03	1.24± 0.26	0.36± 0.16	0.76± 0.14	2.06± 0.77	11.66± 1.41	9	0.12± 0.08
SR6	4.64± 1.26	2.06± 0.48	1.47± 0.32	0.27± 0.04	0.46± 0.19	0.48± 0.11	0.71± 0.13	1.50± 0.62	11.58± 1.24	17	0.17± 0.16
SR7	4.55± 1.22	2.13± 0.52	1.63± 0.35	0.32± 0.04	0.46± 0.15	0.29± 0.06	0.54± 0.10	1.82± 0.76	11.75± 1.23	8	0.27± 0.17
SR8	4.23± 1.12	1.79± 0.49	1.43± 0.31	0.30± 0.05	0.57± 0.13	0.13± 0.08	0.48± 0.09	1.26± 0.57	10.20± 1.11	7	0.36± 0.21
SR9	5.00± 1.48	2.30± 0.56	2.00± 0.45	0.52± 0.06	0.37± 0.17	0.32± 0.07	0.46± 0.08	2.24± 1.07	13.22± 1.68	15	0.66± 0.39
SR10	4.69± 1.32	2.64± 0.62	2.18± 0.49	0.49± 0.06	0.31± 0.06	0.32± 0.06	0.37± 0.07	1.36± 0.64	12.37± 1.37	20	1.14± 0.67
SR11	2.59± 0.74	1.88± 0.46	1.40± 0.34	0.42± 0.06	0.40± 0.09	0.31± 0.07	0.22± 0.04	0.76± 0.43	7.98± 0.96	7	1.30± 0.76
SR12	1.62± 0.52	1.00± 0.27	0.93± 0.22	0.21± 0.03	0.11± 0.02	0.12± 0.03	0.10± 0.02	1.30± 0.65	5.39± 0.84	12	1.41± 0.80
SR13	0.94± 0.29	0.45± 0.13	0.56± 0.17	0.14± 0.02	0.00± 0.02	0.08± 0.01	0.07± 0.01	0.54± 0.33	2.78± 0.48	1	1.41± 0.83
SR14	0.45± 0.18	0.33± 0.10	0.26± 0.08	0.08± 0.01	0.00± 0.02	0.05± 0.01	0.04± 0.01	0.00± 0.11	1.22± 0.30	3	1.35± 0.76
SR15	0.19± 0.09	0.15± 0.06	0.14± 0.05	0.05± 0.01	0.06± 0.05	0.02± 0.00	0.02± 0.00	0.18± 0.12	0.82± 0.22	1	1.09± 0.60
SR16	0.07± 0.04	0.04± 0.03	0.03± 0.01	0.01± 0.00	0.00± 0.00	0.01± 0.00	0.00± 0.00	0.00± 0.03	0.16± 0.08	1	0.62± 0.36
SR17	0.01± 0.01	0.01± 0.02	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.08	0.02± 0.08	0	0.07± 0.06

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**F.3 Nuisances**

Four sets of nuisance pull values, expected and observed for cut-based and BDT analyses, are tabulated in Table 40 (expected cut-based analysis), Table 41 (observed cut-based analysis), Table 42 (expected BDT analysis), Table 43 (observed BDT analysis).

Table 40: Expected nuisance pulls for the cut-based analysis. The final column indicates the correlation between the nuisance and the signal strength

name	<i>b</i> -only fit		<i>s + b</i> fit	$\rho(\theta, \mu)$
	$\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	
TTH	+0.14, 0.99		+0.00, 0.98	-0.08
TTVV	+0.02, 1.00		+0.00, 0.99	-0.01
TTWSF	<b>+0.16, 0.76</b>		<b>+0.00, 0.78</b>	-0.14
TTZSF	<b>+0.06, 0.69</b>		<b>+0.00, 0.69</b>	-0.05
XG	+0.01, 0.99		+0.00, 0.99	-0.00
alphas	+0.00, 0.99		-0.00, 0.99	+0.00
bb	+0.25, 0.96		+0.00, 0.99	-0.13
btaghf	+0.07, 0.99		+0.00, 0.99	-0.07
btaglf	+0.05, 0.99		+0.00, 0.99	-0.03
fakes	+0.09, 0.97		+0.00, 0.97	-0.07
fakes_EWK	+0.07, 0.96		+0.00, 0.96	-0.05
flips	-0.01, 0.99		+0.00, 0.99	+0.00
fsrvar	+0.00, 0.99		-0.00, 0.99	+0.04
isr	-0.14, 0.98		-0.00, 0.99	+0.09
isrvar	+0.00, 0.99		+0.00, 0.99	-0.03
jer	-0.09, 0.97		-0.00, 0.99	+0.05
jes	+0.10, 0.91		+0.00, 0.99	-0.10
lep	+0.02, 0.99		+0.00, 0.99	-0.03
lumi	+0.04, 0.99		+0.00, 0.99	-0.02
pdf	+0.09, 0.96		-0.00, 0.99	-0.05
prefire	-0.01, 0.99		-0.00, 0.99	+0.01
pu	+0.00, 0.99		-0.00, 0.99	+0.02
rares	-0.00, 0.99		+0.00, 0.99	-0.00
scale	-0.11, 1.00		+0.00, 0.99	+0.04
trig	+0.03, 0.99		+0.00, 0.99	-0.04

Table 41: Observed nuisance pulls for the cut-based analysis. The final column indicates the correlation between the nuisance and the signal strength

name	<i>b</i> -only fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	<i>s + b</i> fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	<b>+0.32, 1.02</b>	+0.23, 1.01	-0.08
TTVV	+0.04, 1.00	+0.03, 1.00	-0.01
TTWSF	<b>+1.09, 0.73</b>	<b>+0.97, 0.74</b>	-0.17
TTZSF	<b>+1.23, 0.61</b>	<b>+1.18, 0.62</b>	-0.06
XG	+0.09, 1.00	+0.07, 1.00	-0.01
alphas	-0.00, 0.99	+0.00, 0.99	+0.00
bb	+0.29, 0.97	+0.17, 0.98	-0.11
btaghf	<b>+0.33, 0.99</b>	+0.27, 0.99	-0.07
btaglf	+0.14, 0.99	+0.09, 0.99	-0.04
fakes	+0.25, 1.02	+0.21, 1.01	-0.04
fakes_EWK	+0.17, 0.98	+0.14, 0.95	-0.02
flips	+0.07, 1.00	+0.08, 1.00	+0.01
fsrvar	-0.00, 0.99	<b>-0.03, 1.15</b>	+0.00
isr	-0.24, 0.97	-0.15, 0.98	+0.09
isrvar	-0.00, 0.99	+0.06, 1.01	+0.02
jer	<b>-0.05, 0.52</b>	<b>-0.01, 0.50</b>	+0.04
jes	<b>+0.13, 0.66</b>	<b>+0.07, 0.62</b>	-0.09
lep	+0.16, 0.99	+0.14, 0.99	-0.03
lumi	+0.25, 0.99	+0.22, 0.99	-0.03
pdf	+0.23, 0.94	+0.10, 0.97	-0.12
prefire	-0.04, 0.99	-0.04, 0.99	+0.01
pu	-0.11, 0.98	-0.10, 0.99	+0.01
rares	+0.11, 1.00	+0.12, 1.00	+0.00
scale	+0.01, 0.96	+0.07, 0.95	+0.04
trig	+0.21, 0.99	+0.19, 0.99	-0.04

Table 42: Expected nuisance pulls for the BDT analysis. The final column indicates the correlation between the nuisance and the signal strength

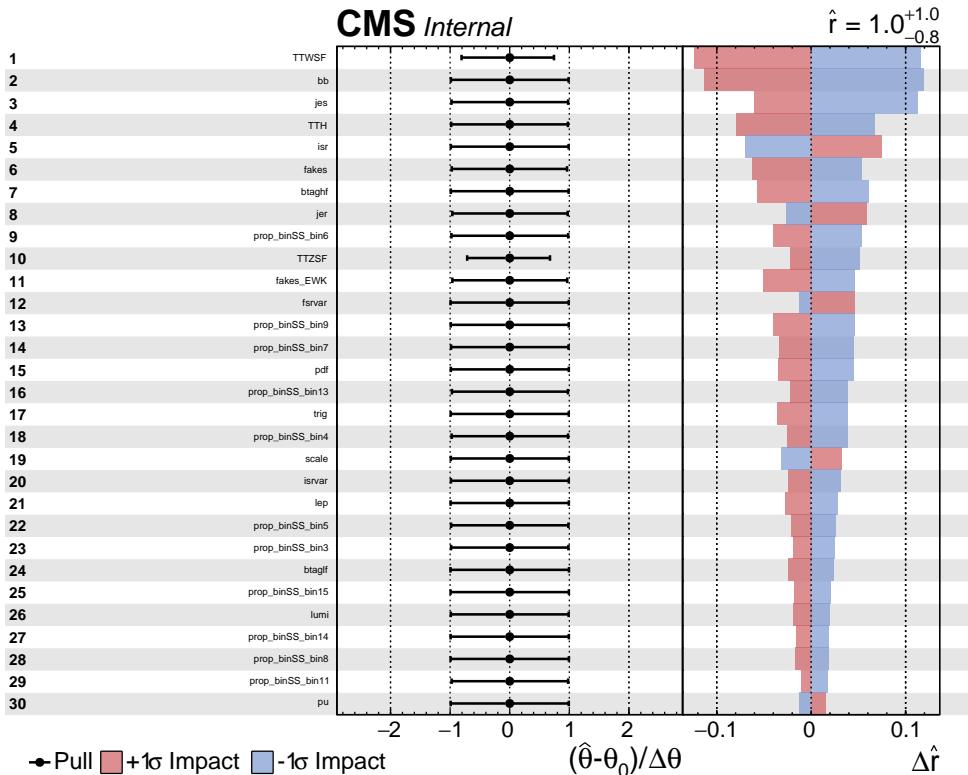
name	$b$ -only fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$s + b$ fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	+0.12, 0.99	-0.00, 0.98	-0.08
TTVV	+0.02, 1.00	+0.00, 0.99	-0.01
TTWSF	<b>+0.20, 0.76</b>	<b>-0.00, 0.78</b>	-0.16
TTZSF	<b>+0.04, 0.69</b>	<b>+0.00, 0.69</b>	-0.05
XG	-0.00, 0.99	+0.00, 0.99	+0.00
alphas	+0.00, 0.99	-0.00, 0.99	+0.00
bb	<b>+0.33, 0.95</b>	+0.00, 0.99	-0.16
btaghf	+0.07, 0.99	+0.00, 0.99	-0.07
btaglf	+0.05, 0.99	+0.00, 0.99	-0.03
fakes	+0.07, 0.97	-0.00, 0.96	-0.06
fakes_EWK	+0.06, 0.94	+0.00, 0.95	-0.06
flips	-0.00, 0.99	-0.00, 0.99	+0.00
fsrvar	+0.00, 0.99	-0.00, 0.99	+0.02
isr	-0.17, 0.97	-0.00, 0.99	+0.10
isrvar	+0.00, 0.99	-0.00, 0.99	-0.03
jer	-0.07, 1.07	+0.00, 0.99	+0.04
jes	+0.11, 0.98	+0.00, 0.97	-0.09
lep	+0.02, 0.99	+0.00, 0.99	-0.03
lumi	+0.04, 0.99	+0.00, 0.99	-0.02
pdf	+0.09, 0.97	-0.00, 0.99	-0.05
prefire	-0.01, 0.99	-0.00, 0.99	+0.01
pu	+0.04, 0.99	-0.00, 0.99	-0.01
rares	+0.00, 0.99	+0.00, 0.99	-0.00
scale	-0.14, 1.01	+0.00, 0.99	+0.04
trig	+0.03, 0.99	+0.00, 0.99	-0.05

Table 43: Observed nuisance pulls for the BDT analysis. The final column indicates the correlation between the nuisance and the signal strength

name	$b$ -only fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$s + b$ fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	<b>+0.33, 1.02</b>	+0.21, 1.01	-0.07
TTVV	+0.03, 1.00	+0.02, 1.00	-0.01
TTWSF	<b>+1.05, 0.74</b>	<b>+0.74, 0.78</b>	-0.20
TTZSF	<b>+1.28, 0.62</b>	<b>+1.24, 0.62</b>	-0.05
XG	-0.02, 0.99	-0.01, 0.99	+0.00
alphas	+0.00, 0.99	+0.00, 1.00	+0.00
bb	<b>+0.63, 0.95</b>	+0.18, 0.98	-0.16
btaghf	+0.25, 0.99	+0.16, 0.99	-0.08
btaglf	+0.12, 0.99	+0.04, 0.99	-0.04
fakes	<b>+0.30, 1.01</b>	+0.26, 1.00	-0.05
fakes_EWK	+0.24, 0.99	+0.21, 0.96	-0.04
flips	+0.05, 1.00	+0.06, 1.00	+0.01
fsrvar	+0.00, 0.99	-0.04, 1.01	+0.00
isr	<b>-0.44, 0.97</b>	-0.17, 0.98	+0.12
isrvar	+0.00, 0.99	+0.01, 1.01	-0.03
jer	-0.15, 0.93	<b>-0.03, 0.83</b>	+0.06
jes	<b>+0.08, 0.78</b>	<b>-0.02, 0.75</b>	-0.09
lep	+0.15, 0.99	+0.13, 0.99	-0.03
lumi	+0.23, 0.99	+0.18, 0.99	-0.02
pdf	+0.23, 0.95	+0.08, 0.97	-0.08
prefire	-0.05, 0.99	-0.04, 0.99	+0.02
pu	-0.18, 1.04	-0.23, 1.03	-0.00
rares	+0.10, 1.00	+0.11, 1.00	+0.00
scale	+0.00, 0.96	+0.13, 0.95	+0.02
trig	+0.19, 0.99	+0.16, 0.99	-0.05

1689 **F.4 Impacts**

1690 The leading 30 nuisance impacts for four sets of impacts, expected and observed for cut-based  
 1691 and BDT analyses, are shown in Figure 103 (expected cut-based analysis), Figure 104 (observed  
 1692 cut-based analysis), Figure 105 (expected BDT analysis), Figure 106 (observed BDT analysis).  
 1693 The most constrained nuisances correspond to normalization parameters for ttW and ttZ; as we  
 1694 would expect from the control regions, “TTWSF” and “TTZSF”, are moved by approximately  
 1 $\sigma$  with respect to the input nuisance sizes.



1695 Figure 103: Expected nuisance impacts for the cut-based analysis.

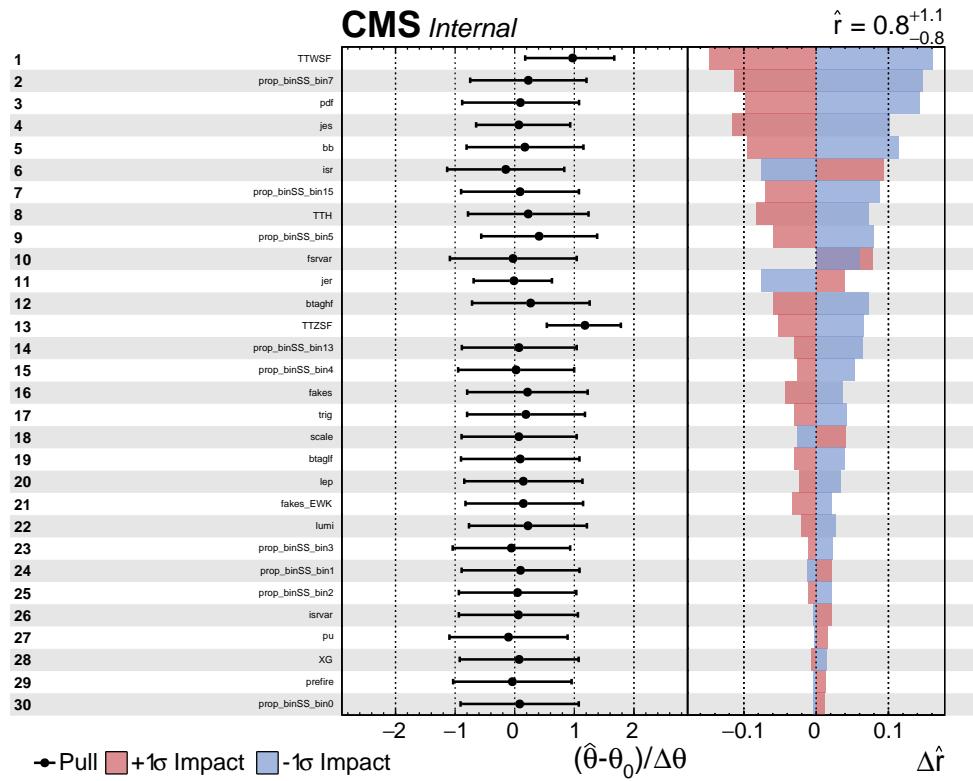


Figure 104: Observed nuisance impacts for the cut-based analysis.

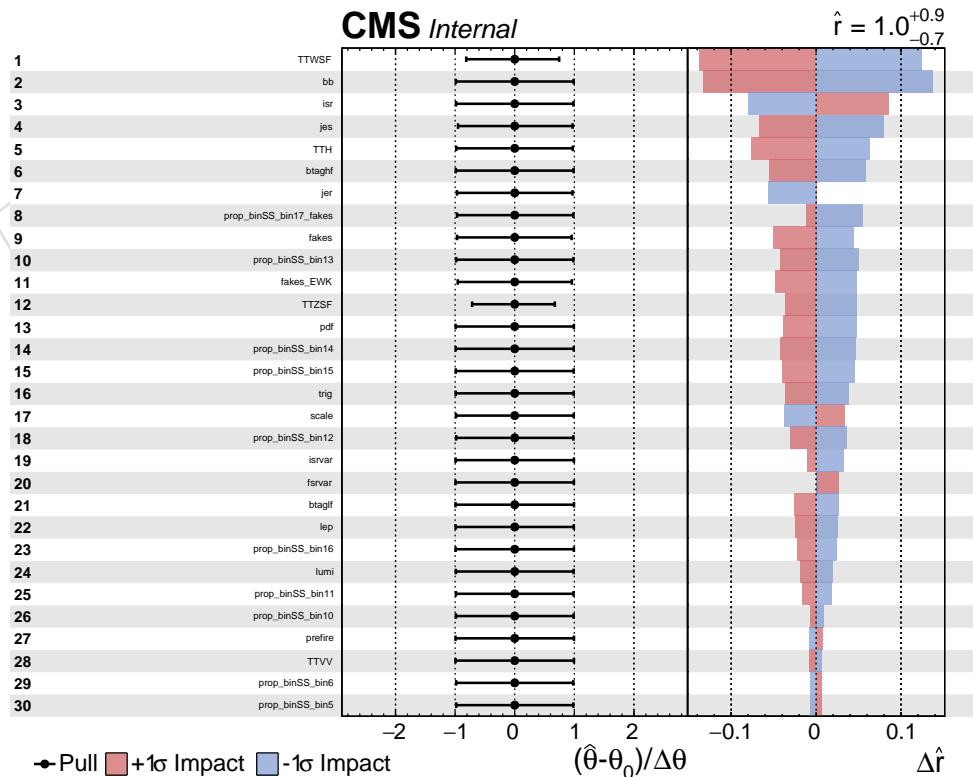


Figure 105: Expected nuisance impacts for the BDT-based analysis.

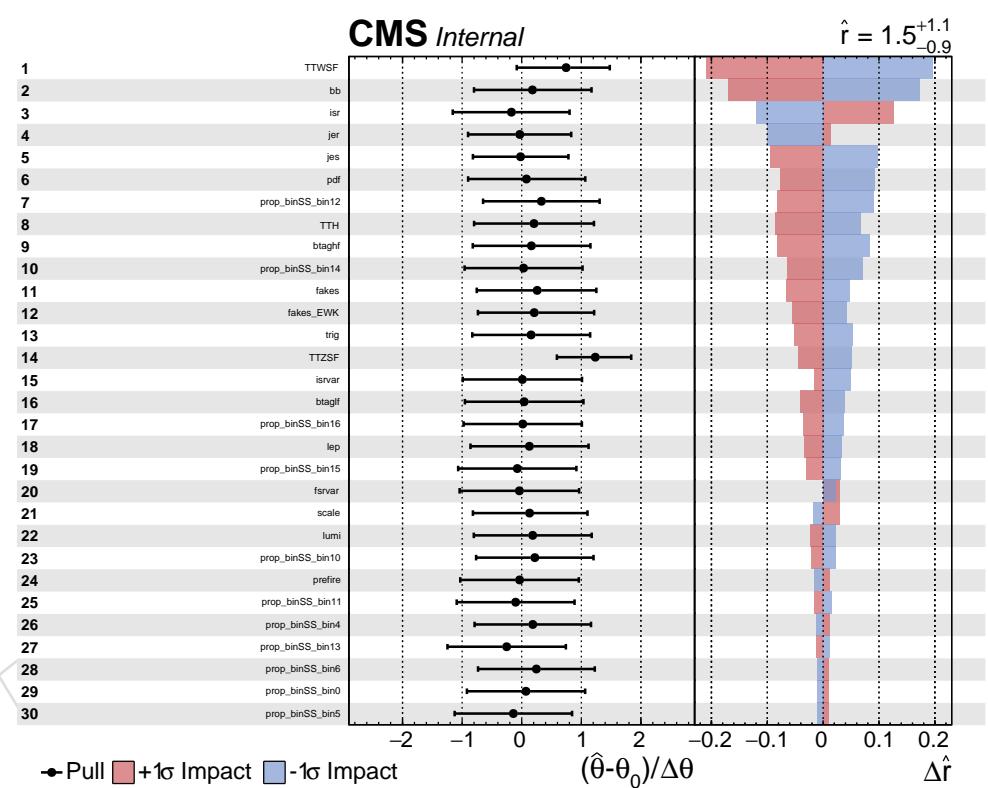


Figure 106: Observed nuisance impacts for the BDT-based analysis.

## 1696 G Unblinding of 2017 dataset

1697 Following the 2016 unblinding, the analysis was unblinded using the 2017 dataset. The yields  
 1698 and results are shown in Section G.1 with consistency checks of cut-based and BDT results  
 1699 in Section G.2. We include the results of various statistical checks, including nuisance pulls  
 1700 (Section G.3), nuisance impacts (Section G.4), and goodness of fits (Section G.5). Additionally,  
 1701 Section G.6 shows information about the combination of unblinded 2017 with unblinded 2016  
 1702 data.

### 1703 G.1 Yields and results

1704 Plots for the unblinded prefit and postfit event yields for the 2017 data, with a total luminosity  
 1705 of  $41.5 \text{ fb}^{-1}$ , are shown for both the cut-based and BDT based analysis in Figure 107. Numerical  
 1706 yields are also tabulated in Table 44 (prefit cut-based analysis), Table 45 (postfit cut-based  
 1707 analysis), Table 46 (prefit BDT analysis), Table 47 (postfit BDT analysis).

1708 A binned likelihood fit is performed over signal regions using the Higgs Combine tool [22];  
 1709 exclusion limits at 95% CL are calculated with the Asymptotic CLs method.

1710 With the cut-based analysis, an observed (expected) upper limit on the production cross section  
 1711 of  $30.18 \text{ fb}$  ( $17.77^{+9.63}_{-5.93} \text{ fb}$ ), assuming the signal process does not exist. The observed (expected)  
 1712 significance is  $0.853$  ( $1.543$ ) standard deviations, corresponding to a measured signal strength  
 1713 parameter of  $0.703^{+0.985}_{-0.703}$ .

1714 With the BDT analysis, an observed (expected) upper limit on the production cross section of  
 1715  $29.85 \text{ fb}$  ( $15.94^{+8.71}_{-5.36} \text{ fb}$ ), assuming the signal process does not exist. The observed (expected)  
 1716 significance is  $1.105$  ( $1.756$ ) standard deviations, corresponding to a measured signal strength  
 1717 parameter of  $0.806^{+0.924}_{-0.736}$ .

### 1718 G.2 Cut-based and BDT consistency checks

1719 In 2016, cut-based SR4 had a high data yield with respect to the SM prediction. In 2017, this  
 1720 region is in agreement. The situation is reversed for cut-based SR3, for example.

1721 For the 2016 unblinding, we quantified the mildness of the disagreement between cut-based  
 1722 and BDT significances (both had differing directions for observed significance with respect to  
 1723 their expected significance), using fits of toy pseudodatasets correlated between cut-based and  
 1724 BDT regions.

1725 Although 2017 shows much better agreement between the two fits, we repeat the same proce-  
 1726 dure for 2017. The observed values lie within a  $0.3\sigma$  contour, with a p-value of nearly 0.8, as  
 1727 shown in Figure 108.

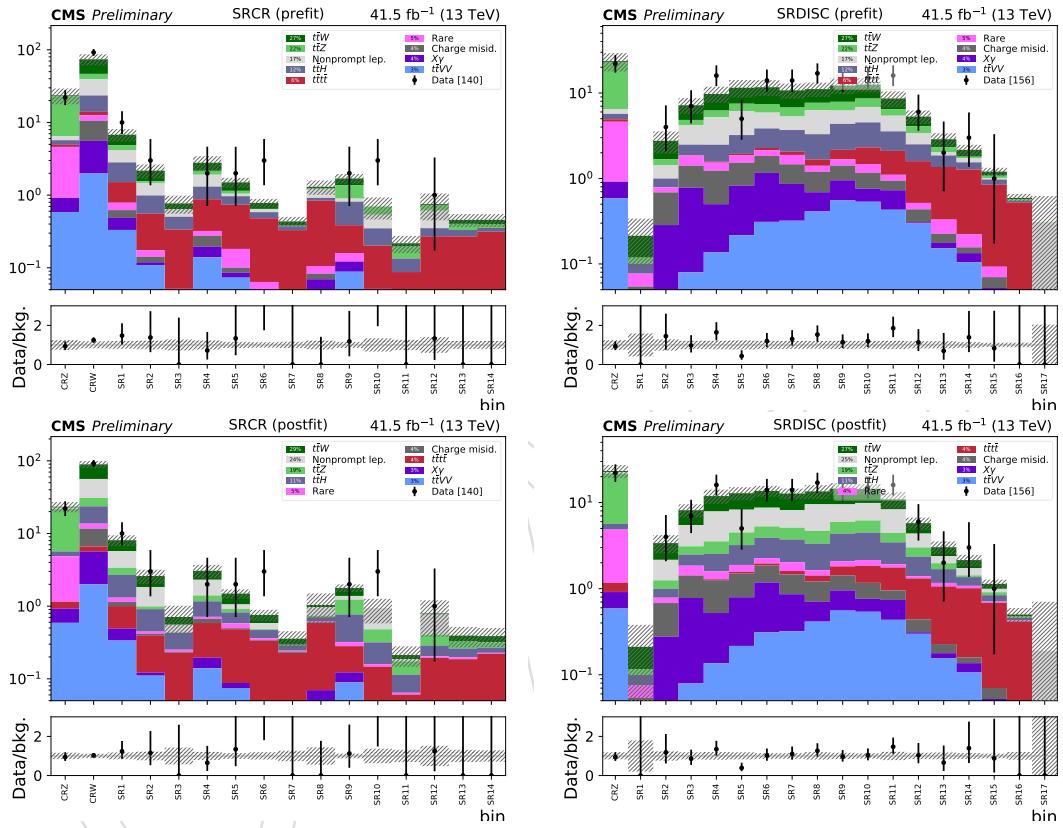


Figure 107: 2017: data yields compared to prefit (top) and postfit (bottom) SM predictions in CRZ, CRW, and the cut-based signal regions (left), and CRZ with the BDT signal regions (right). Note that we plan to retrain the BDT with latest JECs/corrections for 2018 (with identical settings)

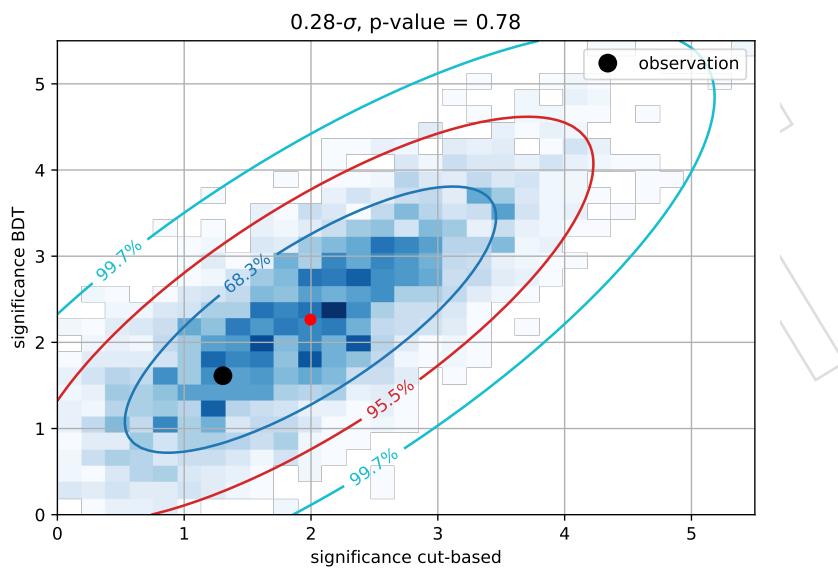


Figure 108: Comparison of correlated expected significances between cut-based and BDT fits with observed point overlaid. Because significances from toy datasets are used, we get a correction factor for expected BDT/cut-based significances with the full fit and the mean of the significances from the toy datasets. This correction factor is then applied to the observed point such that the expected significance values from the full analysis would match the mean values from the toy pseudo-datasets.

Table 44: Prefit event yields in cut-based regions for 2017.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄
CRZ	0.53±0.19	16.56±5.59	0.77±0.21	0.58±0.07	0.33±0.13	3.69±0.78	0.00±0.00	0.79±0.42	23.26±5.81	22	0.32±0.02
CRW	27.27±9.20	744±2.50	9.52±2.48	1.98±0.25	3.60±0.50	2.01±0.41	4.98±0.95	15.64±8.38	72.45±12.99	92	1.43±0.10
SR1	1.92±0.66	0.68±0.24	1.31±0.35	0.33±0.04	0.15±0.04	0.16±0.05	0.14±0.03	1.36±1.04	6.05±1.33	10	0.73±0.06
SR2	0.62±0.25	0.08±0.07	0.43±0.12	0.11±0.01	0.01±0.01	0.03±0.01	0.02±0.00	0.48±0.37	1.79±0.47	3	0.38±0.03
SR3	0.13±0.07	0.00±0.02	0.17±0.06	0.03±0.01	0.00±0.00	0.02±0.00	0.01±0.00	0.14±0.16	0.48±0.21	0	0.28±0.05
SR4	0.62±0.24	0.25±0.12	0.43±0.13	0.14±0.02	0.06±0.01	0.05±0.01	0.08±0.02	0.60±0.50	2.23±0.65	2	0.56±0.07
SR5	0.35±0.18	0.10±0.04	0.23±0.07	0.07±0.01	0.01±0.01	0.08±0.02	0.01±0.00	0.08±0.09	0.94±0.23	2	0.56±0.04
SR6	0.13±0.06	0.01±0.03	0.10±0.03	0.03±0.00	0.01±0.00	0.02±0.00	0.01±0.00	0.05±0.06	0.36±0.11	3	0.42±0.04
SR7	0.05±0.03	0.01±0.01	0.04±0.02	0.01±0.00	0.01±0.00	0.01±0.00	0.01±0.00	0.00±0.04	0.14±0.07	0	0.30±0.03
SR8	0.06±0.04	0.04±0.02	0.08±0.03	0.05±0.01	0.02±0.01	0.02±0.00	0.01±0.00	0.28±0.27	0.57±0.29	0	0.73±0.08
SR9	0.32±0.13	0.48±0.16	0.43±0.11	0.09±0.01	0.03±0.00	0.04±0.01	0.00±0.00	0.08±0.07	1.47±0.26	2	0.23±0.02
SR10	0.00±0.03	0.16±0.07	0.15±0.05	0.03±0.00	0.01±0.01	0.01±0.00	0.00±0.00	0.19±0.21	0.55±0.24	3	0.15±0.01
SR11	0.02±0.01	0.06±0.03	0.05±0.02	0.01±0.00	0.00±0.00	0.00±0.00	0.00±0.00	0.00±0.03	0.14±0.06	0	0.08±0.01
SR12	0.02±0.01	0.12±0.05	0.08±0.03	0.03±0.00	0.01±0.00	0.01±0.00	0.00±0.00	0.26±0.30	0.53±0.30	1	0.23±0.03
SR13	0.05±0.03	0.07±0.03	0.06±0.02	0.01±0.00	0.00±0.00	0.01±0.00	0.00±0.00	0.00±0.07	0.21±0.09	0	0.25±0.02
SR14	0.06±0.03	0.05±0.03	0.04±0.02	0.01±0.00	0.00±0.00	0.01±0.00	0.00±0.00	0.00±0.07	0.16±0.09	0	0.29±0.03

Table 45: Postfit event yields in cut-based regions for 2017.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄
CRZ	0.65±0.20	16.04±3.53	0.81±0.20	0.59±0.08	0.34±0.13	3.70±0.77	0.00±0.00	0.98±0.72	23.11±3.77	22	0.22±0.23
CRW	32.98±9.73	723±1.61	9.98±2.40	2.00±0.26	3.63±0.51	4.19±0.41	5.00±0.84	25.44±12.71	88.27±10.26	92	1.01±1.04
SR1	2.36±0.74	0.67±0.15	1.39±0.33	0.34±0.04	0.16±0.04	0.17±0.04	0.14±0.02	2.35±1.56	7.56±1.37	10	0.51±0.53
SR2	0.76±0.29	0.08±0.06	0.46±0.11	0.11±0.01	0.01±0.01	0.03±0.01	0.02±0.00	0.85±0.59	2.33±0.57	3	0.27±0.29
SR3	0.15±0.07	0.00±0.01	0.18±0.05	0.03±0.01	0.00±0.00	0.02±0.00	0.01±0.00	0.13±0.27	0.51±0.30	0	0.20±0.19
SR4	0.75±0.25	0.24±0.07	0.45±0.12	0.14±0.02	0.06±0.01	0.05±0.01	0.08±0.01	0.91±0.63	2.67±0.64	2	0.39±0.39
SR5	0.43±0.18	0.10±0.03	0.24±0.07	0.07±0.01	0.01±0.00	0.08±0.02	0.01±0.00	0.14±0.12	1.09±0.22	2	0.39±0.41
SR6	0.17±0.07	0.01±0.03	0.11±0.03	0.03±0.00	0.01±0.00	0.02±0.00	0.01±0.00	0.10±0.11	0.46±0.13	3	0.29±0.28
SR7	0.07±0.03	0.01±0.01	0.04±0.02	0.01±0.00	0.01±0.00	0.01±0.00	0.01±0.00	0.00±0.08	0.15±0.10	0	0.21±0.22
SR8	0.07±0.05	0.04±0.01	0.08±0.03	0.05±0.01	0.02±0.01	0.02±0.00	0.01±0.00	0.23±0.45	0.52±0.46	0	0.52±0.52
SR9	0.40±0.15	0.46±0.11	0.45±0.11	0.09±0.01	0.03±0.01	0.04±0.01	0.00±0.00	0.15±0.12	1.62±0.22	2	0.16±0.16
SR10	0.00±0.03	0.16±0.05	0.16±0.04	0.03±0.00	0.01±0.00	0.01±0.00	0.00±0.00	0.44±0.35	0.81±0.34	3	0.11±0.11
SR11	0.03±0.01	0.06±0.02	0.05±0.01	0.01±0.00	0.00±0.00	0.00±0.00	0.00±0.00	0.01±0.06	0.16±0.07	0	0.05±0.06
SR12	0.02±0.01	0.12±0.03	0.08±0.03	0.03±0.00	0.01±0.00	0.01±0.00	0.00±0.00	0.37±0.41	0.63±0.41	1	0.16±0.16
SR13	0.07±0.03	0.07±0.02	0.06±0.02	0.01±0.00	0.00±0.00	0.01±0.00	0.00±0.00	0.00±0.11	0.22±0.12	0	0.17±0.17
SR14	0.07±0.03	0.05±0.02	0.04±0.02	0.01±0.00	0.00±0.00	0.01±0.00	0.00±0.00	0.00±0.10	0.18±0.11	0	0.21±0.21

Table 46: Prefit event yields in BDT regions for 2017.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄
CRZ	0.53±0.21	16.56±5.54	0.77±0.16	0.58±0.08	0.33±0.15	3.69±0.76	0.00±0.00	0.79±0.44	23.26±5.72	22	0.32±0.02
SR1	0.09±0.05	0.02±0.01	0.02±0.01	0.00±0.00	0.01±0.01	0.02±0.03	0.05±0.01	0.00±0.09	0.21±0.13	0	0.00±0.00
SR2	1.07±0.46	0.26±0.15	0.20±0.06	0.03±0.00	0.26±0.12	0.10±0.06	0.40±0.08	0.42±0.53	2.75±0.80	4	0.00±0.00
SR3	2.35±1.01	0.65±0.31	0.62±0.14	0.08±0.02	0.70±0.24	0.44±0.17	0.63±0.12	1.71±1.01	7.18±1.52	7	0.01±0.01
SR4	3.52±1.38	1.05±0.53	0.91±0.21	0.14±0.03	0.36±0.20	0.30±0.14	0.73±0.14	2.70±1.31	9.71±2.10	16	0.04±0.01
SR5	4.02±1.69	1.34±0.60	1.28±0.29	0.22±0.03	0.61±0.19	0.33±0.17	0.70±0.13	2.93±1.88	11.44±2.67	5	0.08±0.01
SR6	4.33±1.73	1.45±0.64	1.58±0.34	0.31±0.04	0.85±0.21	0.30±0.11	0.67±0.13	2.05±1.52	11.54±2.46	14	0.14±0.02
SR7	3.97±1.63	1.44±0.65	1.63±0.35	0.32±0.05	0.54±0.08	0.42±0.15	0.60±0.11	1.63±1.41	10.56±2.17	14	0.17±0.03
SR8	3.42±1.37	1.41±0.57	1.63±0.36	0.41±0.05	0.28±0.09	0.21±0.05	0.50±0.10	2.97±1.60	10.85±2.13	17	0.26±0.02
SR9	4.33±1.74	1.65±0.73	2.19±0.46	0.55±0.06	0.40±0.09	0.26±0.06	0.46±0.09	1.88±1.83	11.73±2.57	14	0.49±0.06
SR10	3.92±1.59	1.78±0.67	2.22±0.49	0.53±0.07	0.23±0.10	0.31±0.07	0.40±0.08	2.30±1.48	11.70±2.39	15	0.83±0.05
SR11	2.18±1.00	1.01±0.45	1.45±0.33	0.43±0.05	0.30±0.08	0.16±0.04	0.21±0.04	1.87±1.45	7.62±1.82	16	1.00±0.04
SR12	1.21±0.55	0.68±0.25	0.91±0.23	0.30±0.04	0.01±0.04	0.08±0.02	0.13±0.02	0.89±0.60	4.21±0.95	6	1.08±0.05
SR13	0.56±0.30	0.30±0.14	0.49±0.13	0.15±0.02	0.02±0.01	0.10±0.02	0.05±0.01	0.16±0.28	1.84±0.51	2	1.03±0.05
SR14	0.36±0.21	0.06±0.04	0.27±0.08	0.10±0.02	0.03±0.01	0.07±0.01	0.02±0.00	0.19±0.17	1.11±0.32	3	1.05±0.05
SR15	0.12±0.07	0.07±0.04	0.12±0.04	0.05±0.01	0.00±0.00	0.02±0.01	0.02±0.00	0.04±0.06	0.45±0.13	1	0.76±0.06
SR16	0.02±0.02	0.02±0.01	0.03±0.01	0.01±0.00	0.00±0.00	0.01±0.00	0.00±0.00	0.00±0.05	0.10±0.07	0	0.49±0.03
SR17	0.00±0.00	0.00±0.00	0.00±0.00	0.00±0.00	0.00±0.00	0.00±0.00	0.00±0.00	0.25±0.32	0.26±0.32	0	0.05±0.01

Table 47: Postfit event yields in BDT regions for 2017.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄
CRZ	0.61±0.20	15.95±3.76	0.80±0.18	0.59±0.07	0.32±0.12	3.68±0.77	0.00±0.00	1.00±0.74	22.96±4.00	22	0.26±0.22
SR1	0.11±0.04	0.02±0.01	0.02±0.01	0.00±0.00	0.01±0.01	0.02±0.03	0.05±0.01	0.00±0.15	0.21±0.17	0	0.00±0.00
SR2	1.22±0.43	0.25±0.10	0.21±0.07	0.03±0.00	0.25±0.11	0.10±0.07	0.40±0.07	0.91±0.78	3.37±0.81	4	0.00±0.00
SR3	2.66±0.95	0.61±0.21	0.64±0.17	0.08±0.02	0.70±0.24	0.42±0.15	0.63±0.11	2.38±1.37	8.13±1.43	7	0.01±0.02
SR4	4.00±1.31	1.00±0.36	0.95±0.26	0.14±0.03	0.38±0.18	0.30±0.16	0.73±0.13	4.35±2.37	11.85±2.25	16	0.03±0.03
SR5	4.56±1.57	1.28±0.38	1.32±0.35	0.22±0.03	0.58±0.15	0.32±0.20	0.70±0.12	3.78±1.98	12.76±2.24	5	0.07±0.07
SR6	4.92±1.61	1.38±0.39	1.63±0.40	0.31±0.04	0.87±0.18	0.30±0.13	0.67±0.11	3.44±2.08	13.52±2.00	14	0.11±0.11
SR7	4.54±1.52	1.37±0.39	1.70±0.41	0.32±0.05	0.54±0.09	0.43±0.17	0.60±0.10	3.07±2.10	12.56±1.85	14	0.14±0.15
SR8	3.90±1.28	1.36±0.34	1.70±0.43	0.41±0.05	0.29±0.10	0.21±0.05	0.51±0.0				

1728 **G.3 Nuisances**

1729 Four sets of nuisance pull values, expected and observed for cut-based and BDT analyses, are  
 1730 tabulated in Table 48 (expected cut-based analysis), Table 49 (observed cut-based analysis),  
 1731 Table 50 (expected BDT analysis), Table 51 (observed BDT analysis).

Table 48: Expected nuisance pulls for the cut-based analysis. The final column indicates the correlation between the nuisance and the signal strength

name	<i>b</i> -only fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	<i>s + b</i> fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	+0.15, 0.99	-0.00, 0.98	-0.08
TTVV	+0.02, 1.00	-0.00, 0.99	-0.01
TTWSF	<b>+0.13, 0.80</b>	<b>-0.00, 0.81</b>	-0.08
TTZSF	<b>+0.05, 0.67</b>	<b>-0.00, 0.67</b>	-0.03
XG	-0.00, 0.99	-0.00, 0.99	+0.00
alphas	+0.00, 0.99	-0.00, 0.99	+0.00
bb	+0.25, 0.96	-0.00, 0.99	-0.12
btaghf	+0.10, 0.98	-0.00, 0.98	-0.08
btaglf	+0.05, 0.99	-0.00, 0.99	-0.03
fakes	+0.14, 0.95	-0.00, 0.96	-0.07
fakes_EWK	<b>+0.10, 0.85</b>	<b>-0.00, 0.85</b>	-0.08
flips	-0.01, 0.99	-0.00, 0.99	+0.00
fsrvar	+0.00, 0.99	-0.00, 0.99	+0.01
isr	+0.10, 0.98	+0.00, 0.99	-0.05
isrvar	+0.00, 0.99	+0.00, 0.99	-0.02
jer	-0.08, 0.95	+0.00, 0.99	+0.05
jes	+0.10, 0.92	+0.00, 0.99	-0.10
lep	+0.03, 0.99	-0.00, 0.99	-0.04
lumi	+0.03, 0.99	-0.00, 0.99	-0.02
pdf	+0.01, 1.01	+0.00, 0.99	+0.01
prefire	-0.02, 0.99	+0.00, 0.99	+0.02
pu	-0.02, 0.99	+0.00, 0.99	+0.01
rares	+0.01, 0.99	-0.00, 0.99	-0.00
scale	-0.05, 0.99	+0.00, 0.99	-0.00
trig	+0.03, 0.99	-0.00, 0.99	-0.04

Table 49: Observed nuisance pulls for the cut-based analysis. The final column indicates the correlation between the nuisance and the signal strength

name	<i>b</i> -only fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	<i>s + b</i> fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	+0.26, 1.02	+0.20, 1.01	-0.07
TTVV	+0.03, 1.00	+0.02, 1.00	-0.01
TTWSF	<b>+0.59, 0.96</b>	<b>+0.56, 0.93</b>	-0.03
TTZSF	<b>-0.08, 0.71</b>	<b>-0.10, 0.69</b>	-0.03
XG	+0.01, 0.99	+0.01, 1.00	+0.00
alphas	-0.00, 0.99	+0.00, 0.99	+0.00
bb	+0.05, 0.98	-0.04, 0.99	-0.09
btaghf	+0.13, 0.98	+0.03, 0.98	-0.11
btaglf	+0.00, 0.99	-0.02, 0.99	-0.03
fakes	<b>+0.70, 0.97</b>	<b>+0.58, 0.98</b>	-0.13
fakes_EWK	<b>+1.00, 0.95</b>	<b>+0.84, 1.01</b>	-0.17
flips	+0.02, 0.99	+0.03, 1.00	+0.01
fsrvar	-0.00, 0.99	+0.02, 1.01	+0.03
isr	+0.21, 0.99	+0.13, 0.99	-0.09
isrvar	-0.00, 0.99	-0.01, 0.99	-0.02
jer	-0.15, 1.06	<b>-0.05, 0.82</b>	+0.08
jes	<b>+0.17, 0.80</b>	<b>+0.13, 0.81</b>	-0.07
lep	+0.05, 0.99	+0.04, 0.99	-0.02
lumi	+0.07, 0.99	+0.06, 0.99	-0.01
pdf	-0.02, 1.05	-0.02, 1.05	-0.00
prefire	-0.03, 0.99	-0.02, 0.99	+0.02
pu	+0.02, 1.01	-0.01, 1.01	-0.03
rares	-0.01, 0.99	-0.01, 0.99	-0.01
scale	-0.11, 1.05	-0.10, 1.07	-0.01
trig	+0.07, 0.99	+0.06, 0.99	-0.03

Table 50: Expected nuisance pulls for the BDT analysis. The final column indicates the correlation between the nuisance and the signal strength

name	$b$ -only fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$s + b$ fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	+0.17, 0.99	+0.00, 0.98	-0.08
TTVV	+0.03, 1.00	+0.00, 0.99	-0.01
TTWSF	<b>+0.16, 0.80</b>	<b>+0.00, 0.82</b>	-0.09
TTZSF	<b>+0.04, 0.67</b>	<b>+0.00, 0.67</b>	-0.03
XG	-0.01, 0.99	+0.00, 0.99	+0.00
alphas	+0.00, 0.99	+0.00, 0.99	+0.00
bb	<b>+0.36, 0.95</b>	+0.00, 0.99	-0.14
btaghf	+0.11, 0.98	+0.00, 0.98	-0.09
btaglf	+0.05, 0.99	+0.00, 0.99	-0.03
fakes	+0.05, 0.94	+0.00, 0.94	-0.04
fakes_EWK	<b>+0.12, 0.86</b>	<b>+0.00, 0.82</b>	-0.10
flips	-0.01, 0.99	+0.00, 0.99	+0.00
fsrvar	+0.00, 0.99	+0.00, 0.99	+0.01
isr	+0.08, 0.99	+0.00, 0.99	-0.04
isrvar	+0.00, 0.99	+0.00, 0.99	-0.01
jer	<b>-0.10, 1.32</b>	+0.00, 0.99	+0.04
jes	+0.12, 0.97	+0.00, 0.97	-0.08
lep	+0.03, 0.99	+0.00, 0.99	-0.04
lumi	+0.04, 0.99	+0.00, 0.99	-0.02
pdf	+0.03, 1.01	+0.00, 0.99	-0.01
prefire	-0.02, 0.99	+0.00, 0.99	+0.02
pu	+0.05, 0.99	+0.00, 0.99	-0.02
rares	+0.02, 0.99	+0.00, 0.99	-0.00
scale	-0.02, 0.97	+0.00, 0.99	-0.01
trig	+0.04, 0.99	+0.00, 0.99	-0.05

Table 51: Observed nuisance pulls for the BDT analysis. The final column indicates the correlation between the nuisance and the signal strength

name	<i>b</i> -only fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	<i>s + b</i> fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	+0.25, 1.02	+0.17, 1.01	-0.07
TTVV	+0.04, 1.00	+0.02, 1.00	-0.01
TTWSF	<b>+0.42, 0.98</b>	<b>+0.38, 0.95</b>	-0.03
TTZSF	<b>-0.11, 0.69</b>	<b>-0.13, 0.69</b>	-0.02
XG	-0.01, 0.99	+0.00, 0.99	+0.00
alphas	+0.00, 0.99	+0.00, 0.99	+0.00
bb	+0.23, 0.97	+0.04, 0.99	-0.14
btaghf	+0.10, 0.98	+0.03, 0.98	-0.09
btaglf	+0.06, 0.99	+0.02, 0.99	-0.04
fakes	<b>+0.75, 0.97</b>	<b>+0.64, 0.97</b>	-0.09
fakes_EWK	<b>+1.01, 0.88</b>	<b>+0.83, 0.91</b>	-0.17
flips	+0.00, 0.99	+0.01, 0.99	+0.01
fsrvar	+0.00, 0.99	-0.01, 0.99	-0.00
isr	+0.14, 0.99	+0.07, 0.99	-0.05
isrvar	+0.00, 0.99	-0.01, 0.99	-0.02
jer	<b>-0.09, 1.15</b>	-0.05, 1.06	+0.03
jes	+0.24, 1.01	<b>+0.17, 1.10</b>	-0.09
lep	+0.03, 0.99	+0.02, 0.99	-0.02
lumi	+0.05, 0.99	+0.04, 0.99	-0.01
pdf	+0.01, 1.03	-0.03, 1.03	-0.01
prefire	-0.01, 0.99	-0.01, 0.99	+0.02
pu	+0.08, 1.00	+0.04, 1.00	-0.03
rares	-0.02, 0.99	-0.03, 0.99	-0.01
scale	+0.03, 0.93	+0.04, 0.94	-0.01
trig	+0.05, 0.99	+0.04, 0.99	-0.03

## 1732 G.4 Impacts

1733 The leading 30 nuisance impacts for four sets of impacts, expected and observed for cut-based  
 1734 and BDT analyses, are shown in Figure 109 (expected cut-based analysis), Figure 110 (observed  
 1735 cut-based analysis), Figure 111 (expected BDT analysis), Figure 112 (observed BDT analysis).  
 1736 Note that statistical nuisances tend to show up in the impact plots because the binning for the  
 1737 cut-based and BDT analyses were optimized for the full luminosity (3.3 times what is shown  
 1738 here for 2017). As in 2016, the most constrained nuisances correspond to normalization  
 1739 parameters for ttW and ttZ, as we would expect from the control regions. “TTWSF” is moved  
 1740 by approximately  $1\sigma$  with respect to the input nuisance sizes. Because the ttZ control region  
 1741 has better data, simulation agreement in 2017 compared to 2016, the “TTZSF” nuisance is not  
 pulled up here.

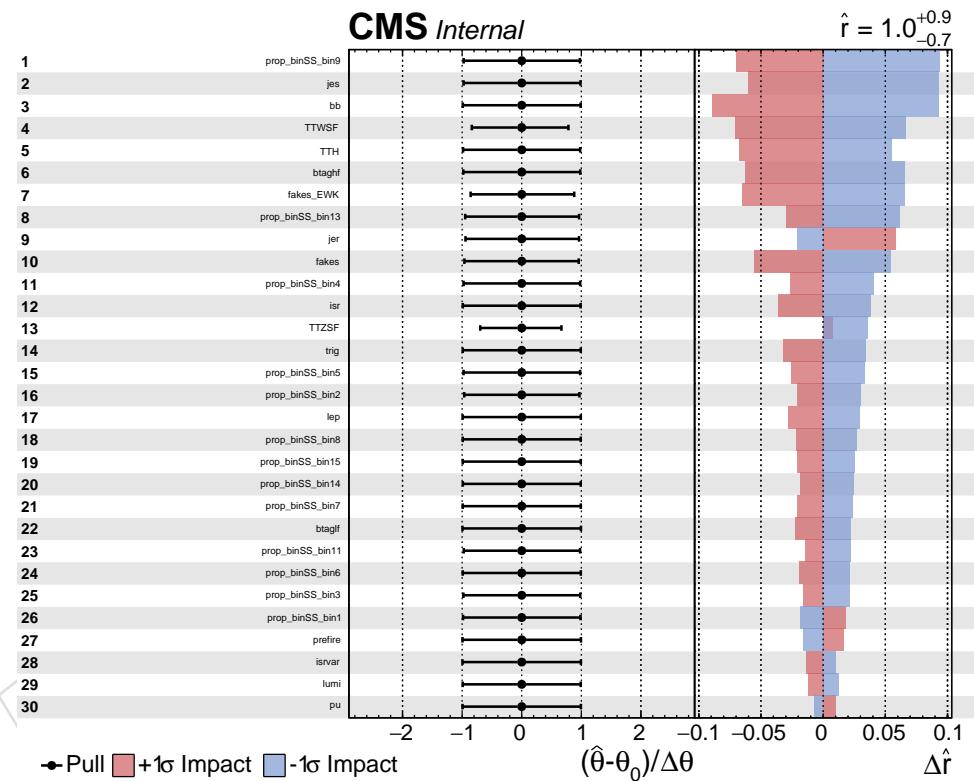


Figure 109: Expected nuisance impacts for the cut-based analysis.

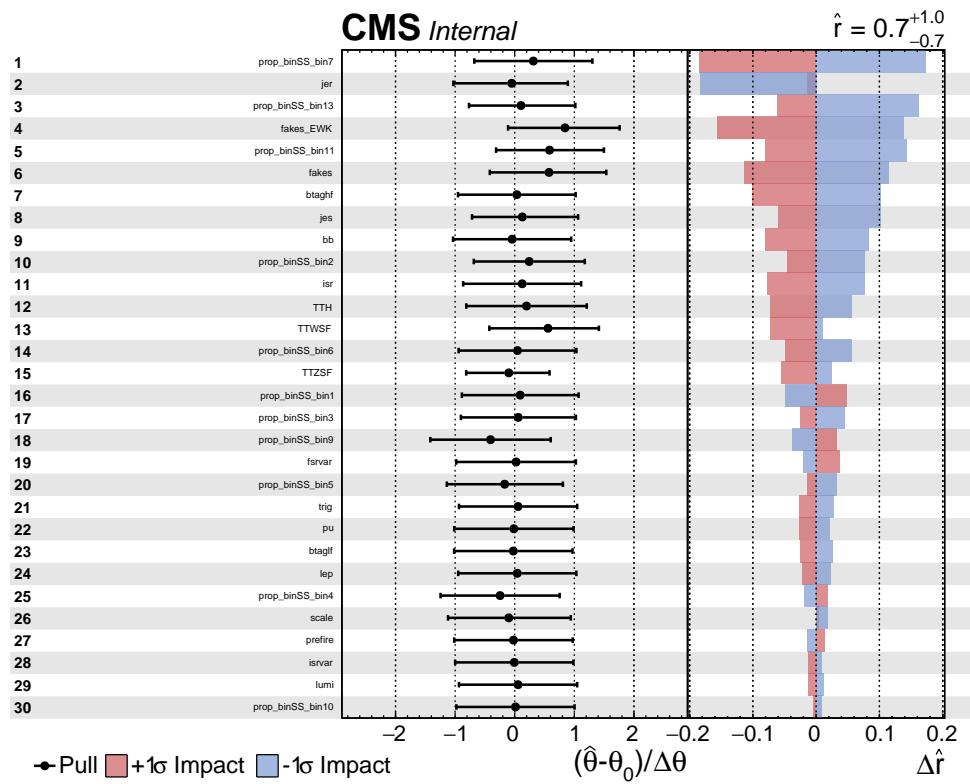


Figure 110: Observed nuisance impacts for the cut-based analysis.

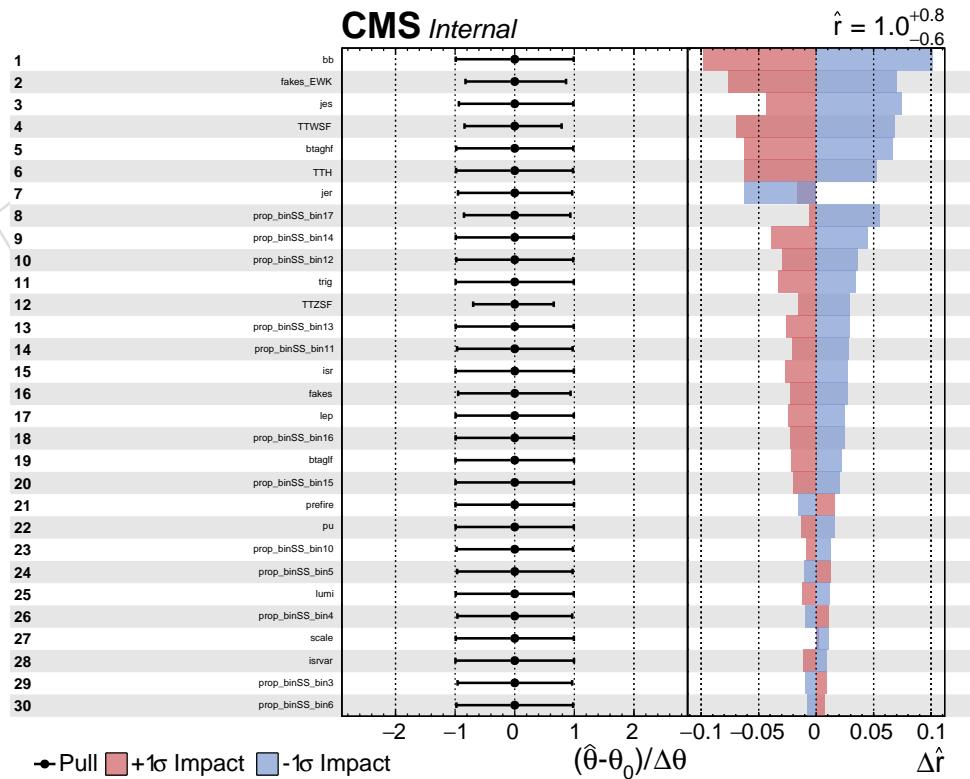


Figure 111: Expected nuisance impacts for the BDT-based analysis.

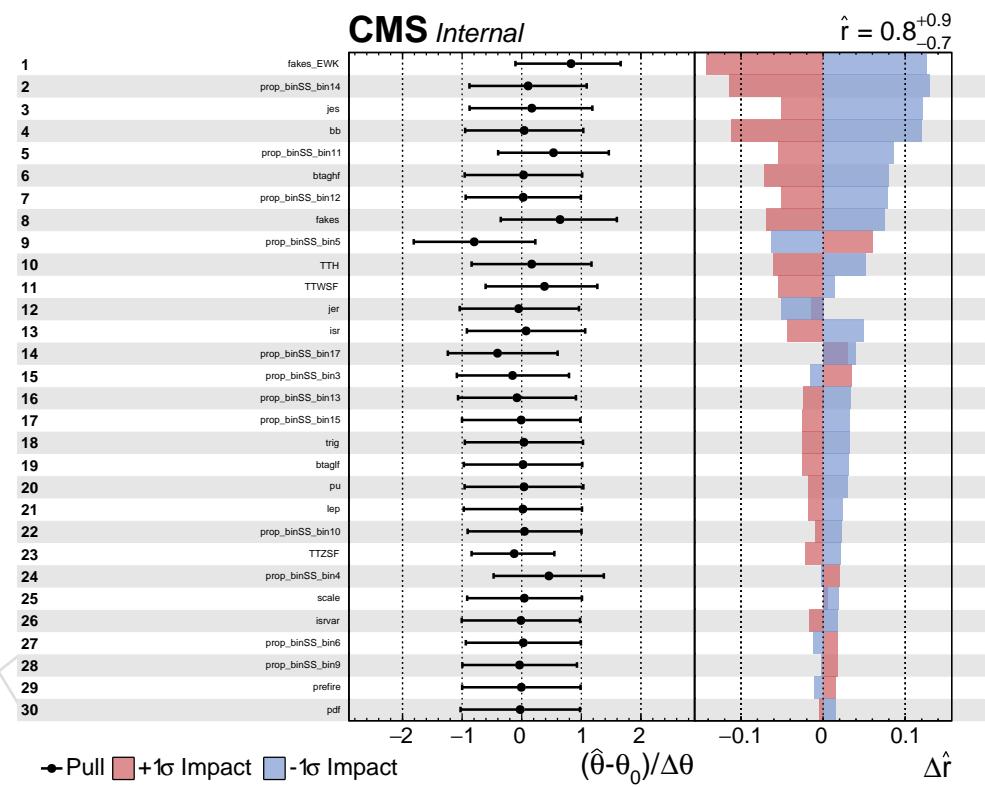


Figure 112: Observed nuisance impacts for the BDT-based analysis.

1743 **G.5 Goodness of fits**

1744 The goodness of fit distributions (using the saturated test statistic with the signal+background fit to data and asimov toys) for the cut-based and BDT analyses are shown in Figure 113.

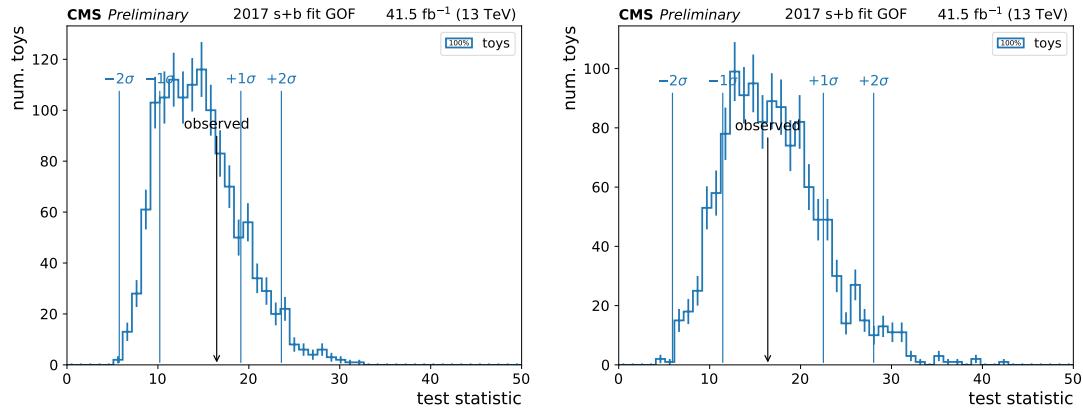


Figure 113: GOF test for the cut-based analysis (left) and BDT-based analysis (right)

1745

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## 1746 G.6 Combination with 2016

1747 This subsection presents the results of the combination of the unblinded 2017 dataset from  
 1748 above with the 2016 dataset from Appendix F following the correlation model from Section 9.

### 1749 G.6.1 Yields and results

1750 Plots for the unblinded prefit and postfit event yields for the 2016+2017 data, with a total lu-  
 1751 minosity of  $35.9 \text{ fb}^{-1} + 41.5 \text{ fb}^{-1} = 77.4 \text{ fb}^{-1}$ , are shown for both the cut-based and BDT based  
 1752 analysis in Figure 114. Numerical yields are also tabulated in Table 52 (postfit cut-based analy-  
 1753 sis), Table 53 (postfit BDT analysis).

1754 With the cut-based analysis, an observed (expected) upper limit on the production cross section  
 1755 of  $23.15 \text{ fb}$  ( $12.34^{+6.30}_{-4.00} \text{ fb}$ ), assuming the signal process does not exist. The observed (expected)  
 1756 significance is  $1.208$  ( $2.074$ ) standard deviations, corresponding to a measured signal strength  
 1757 parameter of  $0.696^{+0.690}_{-0.584}$ .

1758 With the BDT analysis, an observed (expected) upper limit on the production cross section of  
 1759  $28.47 \text{ fb}$  ( $11.36^{+5.84}_{-3.71} \text{ fb}$ ), assuming the signal process does not exist. The observed (expected)  
 1760 significance is  $2.221$  ( $2.272$ ) standard deviations, corresponding to a measured signal strength  
 1761 parameter of  $1.167^{+0.697}_{-0.593}$ .

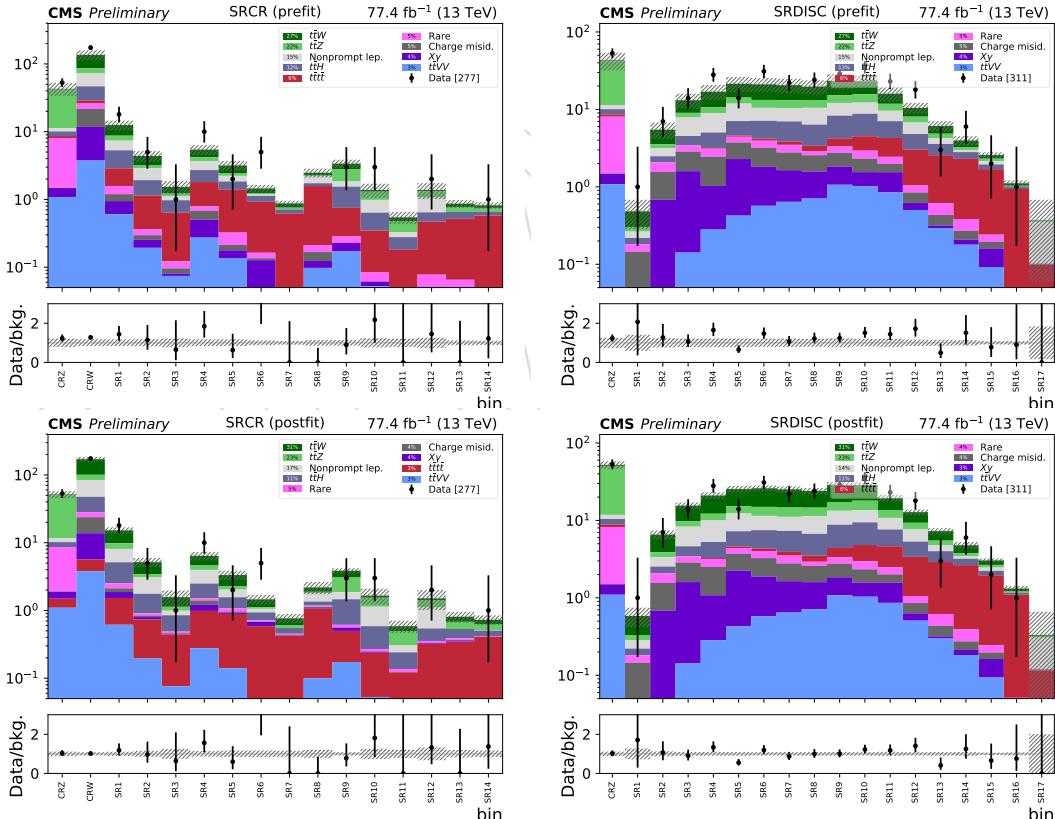


Figure 114: 2016+2017: data yields compared to prefit (top) and postfit (bottom) SM predictions in CRZ, CRW, and the cut-based signal regions (left), and CRZ with the BDT signal regions (right). Note that we plan to retrain the BDT with latest JECs/corrections for 2018 (with identical settings)

Table 52: Postfit event yields in cut-based regions for 2016+2017.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄
CRZ	1.45± 0.36	38.30± 6.11	1.57± 0.38	1.10± 0.13	0.40± 0.09	6.71± 1.25	0.02± 0.00	1.50± 0.74	51.05± 6.13	53	0.39± 0.35
CRW	70.12±16.58	17.21± 2.75	19.31± 4.64	3.79± 0.45	8.14± 0.98	4.44± 0.83	9.98± 1.16	35.72±18.98	168.72±14.28	174	1.85± 1.56
SR1	5.38± 1.34	1.71± 0.28	2.63± 0.66	0.62± 0.07	0.34± 0.07	0.39± 0.08	0.25± 0.03	2.96± 1.89	14.28± 1.65	18	0.91± 0.79
SR2	1.72± 0.45	0.35± 0.08	0.85± 0.21	0.20± 0.02	0.06± 0.02	0.07± 0.01	0.04± 0.01	1.32± 0.86	4.62± 0.79	5	0.54± 0.45
SR3	0.44± 0.15	0.14± 0.04	0.36± 0.10	0.08± 0.01	0.01± 0.00	0.03± 0.01	0.02± 0.00	0.14± 0.34	1.21± 0.40	1	0.36± 0.31
SR4	1.79± 0.44	0.64± 0.11	0.98± 0.23	0.28± 0.04	0.24± 0.05	0.12± 0.02	0.18± 0.02	1.47± 1.07	5.69± 0.92	10	0.71± 0.57
SR5	1.07± 0.31	0.23± 0.05	0.49± 0.13	0.14± 0.02	0.04± 0.01	0.11± 0.02	0.04± 0.00	0.46± 0.44	2.58± 0.53	2	0.75± 0.61
SR6	0.34± 0.11	0.06± 0.04	0.22± 0.06	0.05± 0.01	0.09± 0.03	0.03± 0.01	0.01± 0.00	0.12± 0.17	0.92± 0.24	5	0.54± 0.44
SR7	0.16± 0.06	0.06± 0.02	0.10± 0.03	0.02± 0.00	0.01± 0.00	0.01± 0.00	0.01± 0.00	0.00± 0.07	0.36± 0.12	0	0.40± 0.32
SR8	0.30± 0.09	0.10± 0.02	0.17± 0.04	0.10± 0.01	0.03± 0.01	0.05± 0.01	0.04± 0.01	0.44± 0.42	1.22± 0.45	0	0.96± 0.80
SR9	0.79± 0.20	1.21± 0.19	0.85± 0.20	0.17± 0.02	0.06± 0.01	0.06± 0.01	0.00± 0.00	0.39± 0.22	3.53± 0.34	3	0.33± 0.24
SR10	0.11± 0.04	0.41± 0.08	0.32± 0.08	0.05± 0.01	0.01± 0.00	0.02± 0.01	0.00± 0.00	0.55± 0.43	1.47± 0.44	3	0.18± 0.17
SR11	0.09± 0.03	0.19± 0.05	0.10± 0.03	0.02± 0.00	0.00± 0.00	0.01± 0.00	0.00± 0.00	0.06± 0.09	0.48± 0.12	0	0.10± 0.09
SR12	0.08± 0.03	0.31± 0.05	0.19± 0.04	0.04± 0.01	0.01± 0.00	0.03± 0.01	0.00± 0.00	0.57± 0.50	1.23± 0.49	2	0.28± 0.23
SR13	0.13± 0.05	0.15± 0.04	0.15± 0.04	0.03± 0.00	0.02± 0.00	0.02± 0.00	0.00± 0.00	0.00± 0.10	0.49± 0.15	0	0.32± 0.27
SR14	0.10± 0.03	0.12± 0.03	0.09± 0.02	0.02± 0.00	0.00± 0.00	0.01± 0.00	0.00± 0.00	0.01± 0.09	0.34± 0.12	1	0.38± 0.33

Table 53: Postfit event yields in BDT regions for 2016+2017.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄
CRZ	1.46± 0.36	38.83± 6.01	1.58± 0.36	1.10± 0.12	0.40± 0.10	6.70± 1.18	0.02± 0.00	1.34± 0.52	51.42± 6.02	53	0.66± 0.35
SR1	0.25± 0.07	0.05± 0.02	0.04± 0.02	0.00± 0.00	0.03± 0.01	0.04± 0.03	0.12± 0.02	0.06± 0.13	0.58± 0.16	1	0.00± 0.00
SR2	2.69± 0.68	0.53± 0.14	0.44± 0.12	0.05± 0.01	0.64± 0.22	0.52± 0.18	0.88± 0.12	0.81± 0.57	6.56± 0.86	7	0.00± 0.01
SR3	5.76± 1.52	1.46± 0.31	1.17± 0.31	0.14± 0.02	1.46± 0.28	0.57± 0.15	1.26± 0.17	3.67± 1.37	15.50± 1.55	14	0.03± 0.02
SR4	8.76± 2.13	2.23± 0.54	1.99± 0.49	0.28± 0.03	0.80± 0.21	0.67± 0.18	1.43± 0.20	4.72± 1.82	20.89± 2.05	28	0.09± 0.05
SR5	10.13± 2.50	2.96± 0.68	2.66± 0.65	0.43± 0.04	1.82± 0.28	0.69± 0.30	1.46± 0.20	5.06± 2.04	25.21± 2.50	14	0.18± 0.11
SR6	11.04± 2.67	3.50± 0.60	3.17± 0.77	0.58± 0.06	1.31± 0.26	0.77± 0.18	1.37± 0.19	3.96± 1.86	25.70± 2.27	31	0.29± 0.19
SR7	10.50± 2.58	3.55± 0.66	3.40± 0.81	0.65± 0.07	1.00± 0.18	0.73± 0.18	1.14± 0.15	3.80± 1.91	24.77± 2.33	22	0.41± 0.23
SR8	9.40± 2.32	3.26± 0.60	3.21± 0.80	0.72± 0.09	0.88± 0.13	0.34± 0.08	0.99± 0.13	4.66± 1.92	23.45± 2.21	24	0.58± 0.30
SR9	11.50± 2.89	3.98± 0.71	4.40± 1.08	1.08± 0.11	0.77± 0.20	0.57± 0.12	0.92± 0.13	4.52± 2.34	27.73± 2.73	29	1.07± 0.54
SR10	10.62± 2.62	4.40± 0.70	4.63± 1.15	1.03± 0.12	0.53± 0.11	0.63± 0.12	0.77± 0.10	4.13± 1.95	26.75± 2.45	35	1.84± 1.00
SR11	5.91± 1.57	2.83± 0.51	3.00± 0.76	0.86± 0.10	0.71± 0.10	0.48± 0.10	0.44± 0.06	3.15± 1.88	17.37± 2.03	23	2.16± 1.14
SR12	3.46± 0.91	1.66± 0.29	1.94± 0.48	0.51± 0.06	0.12± 0.04	0.20± 0.05	0.23± 0.03	2.33± 0.96	10.45± 1.08	18	2.33± 1.21
SR13	1.79± 0.49	0.74± 0.13	1.11± 0.29	0.30± 0.04	0.02± 0.02	0.18± 0.03	0.12± 0.02	0.75± 0.50	4.99± 0.67	3	2.28± 1.22
SR14	1.00± 0.33	0.35± 0.08	0.55± 0.14	0.18± 0.02	0.03± 0.03	0.12± 0.02	0.06± 0.01	0.23± 0.23	2.53± 0.46	6	2.25± 1.18
SR15	0.38± 0.11	0.20± 0.04	0.27± 0.07	0.09± 0.01	0.07± 0.04	0.05± 0.01	0.03± 0.00	0.23± 0.13	1.33± 0.19	2	1.71± 0.88
SR16	0.11± 0.04	0.05± 0.03	0.06± 0.02	0.02± 0.00	0.00± 0.00	0.02± 0.00	0.01± 0.00	0.00± 0.06	0.27± 0.10	1	1.05± 0.56
SR17	0.01± 0.00	0.02± 0.01	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.00± 0.00	0.19± 0.33	0.22± 0.33	0	0.11± 0.05

## 1762 H Unblinding of 2018 dataset, BDT retraining

1763 For reference, a table with “historical” numbers can be found here: <https://docs.google.com/spreadsheets/d/140yqrUwEsJtOJ8OmDdOun-J8iUrV49aAdUkvfNjpn8k/edit#gid=0>  
 1764 The last row of the table corresponds to the results presented below.

### 1766 H.1 BDT retraining

1767 With the availability of latest recipes for 2018, the BDT was retrained, keeping identical settings  
 1768 to what was used previously, in Section 6.3. In particular, the finalized retraining included two  
 1769 new key components: Autumn18V8 MC JECs and 2018 b-tag scale factors.

### 1770 H.2 Yields and results

1771 Plots for the prefit and postfit event yields for the 2018 data, with a total luminosity of  $59.6 \text{ fb}^{-1}$ ,  
 1772 are shown for both the cut-based and BDT based analysis in Figure 115. Numerical yields  
 1773 are also tabulated in Table 54 (prefit cut-based analysis), Table 55 (postfit cut-based analysis),  
 1774 Table 56 (prefit BDT analysis), and Table 57 (postfit BDT analysis).

1775 Similarly, plots of yields for combination of the full Run2 dataset, corresponding to  $137.2 \text{ fb}^{-1}$ ,  
 1776 are shown for both the cut-based and BDT-based analyses in Figure 116. Full Run2 dataset  
 1777 numerical yields are also tabulated in Table 58 (prefit cut-based analysis), Table 59 (postfit cut-  
 1778 based analysis), Table 60 (prefit BDT analysis), and Table 61 (postfit BDT analysis).

1779 A binned likelihood fit is performed over signal regions using the Higgs Combine tool [22]  
 1780 to extract signal strength and significance; exclusion limits at 95% CL are calculated with the  
 1781 Asymptotic CLs method.

1782 Using the full run2 data, the *cut-based* analysis sets an observed (expected) upper limit on the  
 1783 production cross section of  $20.31 \text{ fb}$  ( $9.40^{+4.38}_{-2.87} \text{ fb}$ ), assuming the signal process does not exist.  
 1784 The observed (expected) significance is  $1.718$  ( $2.461$ ) standard deviations, corresponding to a  
 1785 measured observed (expected) signal strength parameter of  $0.795^{+0.520}_{-0.475}$  ( $1.000^{+0.485}_{-0.436}$ )

1786 The *BDT* analysis sets an observed (expected) upper limit on the production cross section of  
 1787  $22.86 \text{ fb}$  ( $8.56^{+3.96}_{-2.63} \text{ fb}$ ), assuming the signal process does not exist. The observed (expected)  
 1788 significance is  $2.568$  ( $2.680$ ) standard deviations, corresponding to a measured observed (ex-  
 1789 pected) signal strength parameter of  $1.069^{+0.490}_{-0.443}$  ( $1.000^{+0.447}_{-0.403}$ )

1790 To isolate the effect of the updates to 2018 data and MC, we recalculated the expected limit and  
 1791 significance with the latest corrections used above, but with the old BDT function, which gave  
 1792 an expected upper limit on the production cross section of  $8.95^{+4.25}_{-2.80}$  (central value 4.6% worse  
 1793 than new training, relative), an expected significance of 2.636 (1.6% worse than new training,  
 1794 relative), and an expected signal strength of  $1.000^{+0.470}_{-0.422}$  (measurement precision 5% worse than  
 1795 new training, relative).

1796 Due to the retraining and updated SFs, updated numbers for the BDT analysis observed (ex-  
 1797 pected) significances are as follows: 1.342 (1.496) for 2016, 1.182 (1.740) for 2017, and 2.950  
 1798 (1.924) for 2018.

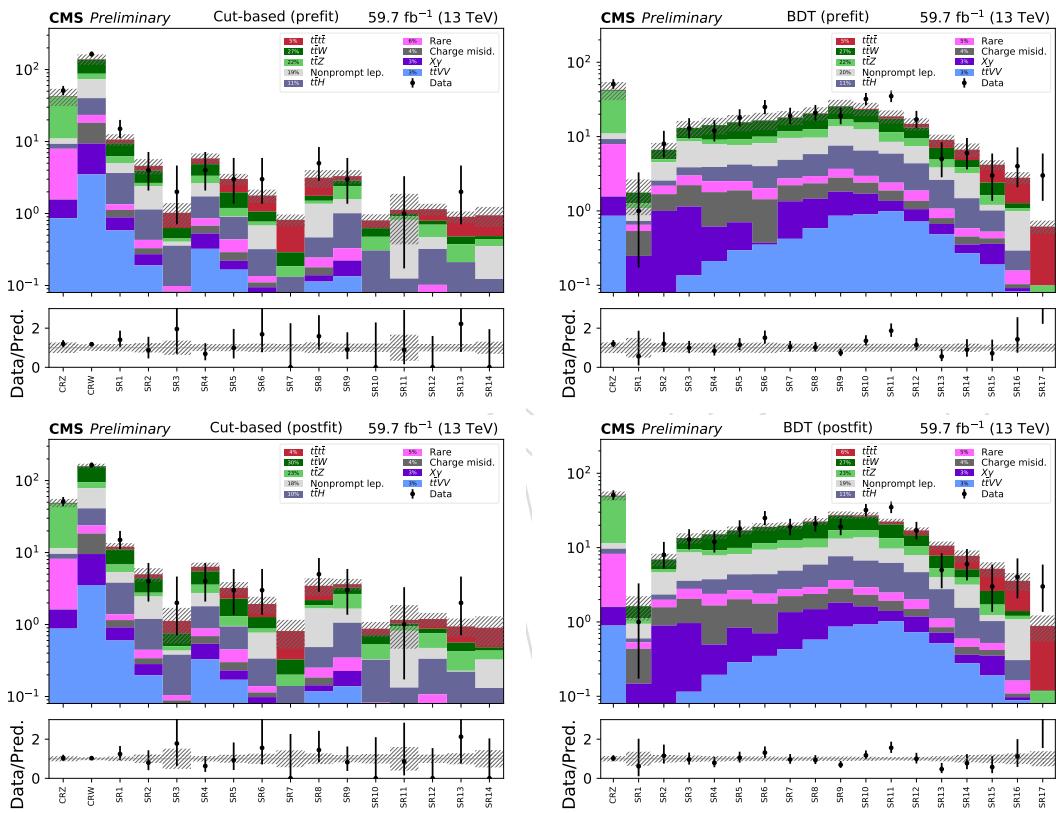


Figure 115: 2018: data yields compared to prefit (top) and postfit (bottom) SM predictions in CRZ, CRW, and the cut-based signal regions (left), and CRZ with the BDT signal regions (right).

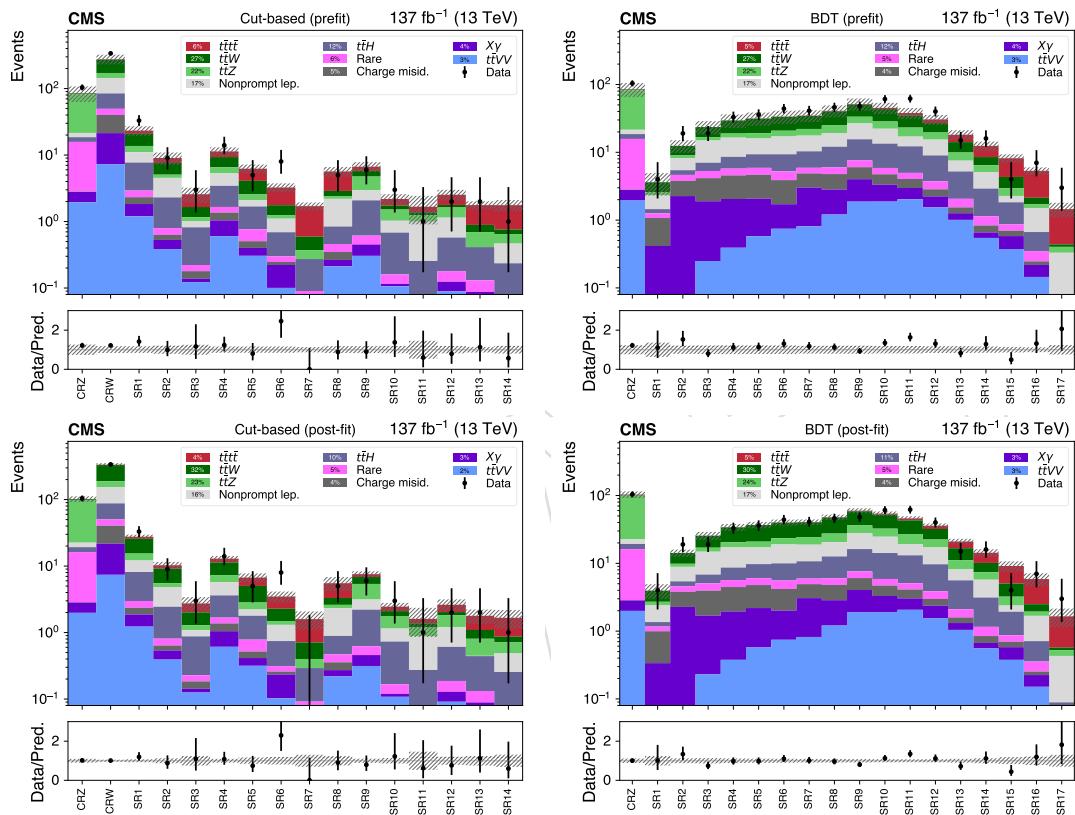


Figure 116: Run2: data yields compared to prefit (top) and postfit (bottom) SM predictions in CRZ, CRW, and the cut-based signal regions (left), and CRZ with the BDT signal regions (right).

Table 54: Prefit event yields in cut-based regions for 2018.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	0.87±0.33	29.83±11.02	1.33±0.36	0.85±0.12	0.71±0.41	6.28±1.30	0.01±0.00	1.72±0.60	41.61±11.49	51	0.48±0.03
CRW	48.99±18.52	13.77±4.92	16.49±4.22	3.45±0.44	5.75±1.14	5.18±1.11	8.79±1.75	33.41±9.77	135.83±23.00	163	2.21±0.08
SR1	3.21±1.37	1.23±0.47	2.29±0.64	0.57±0.09	0.30±0.09	0.22±0.05	0.23±0.05	1.32±0.62	9.37±1.96	15	1.16±0.07
SR2	1.27±0.53	0.25±0.16	0.70±0.24	0.19±0.04	0.08±0.04	0.10±0.03	0.06±0.01	1.27±0.79	3.92±1.09	4	0.64±0.04
SR3	0.18±0.17	0.04±0.08	0.26±0.10	0.05±0.01	0.01±0.02	0.02±0.00	0.03±0.01	0.05±0.04	0.62±0.31	2	0.39±0.07
SR4	1.48±0.62	0.71±0.31	0.86±0.27	0.32±0.05	0.20±0.06	0.18±0.05	0.15±0.03	0.89±0.42	4.79±1.01	4	0.96±0.05
SR5	0.76±0.32	0.28±0.18	0.46±0.15	0.16±0.03	0.05±0.03	0.14±0.05	0.07±0.01	0.01±0.05	1.94±0.54	3	1.04±0.05
SR6	0.29±0.14	0.09±0.04	0.18±0.07	0.05±0.01	0.04±0.01	0.02±0.00	0.02±0.00	0.36±0.24	1.06±0.33	3	0.70±0.08
SR7	0.10±0.05	0.05±0.03	0.09±0.04	0.02±0.00	0.00±0.00	0.02±0.01	0.00±0.00	0.00±0.04	0.28±0.11	0	0.52±0.07
SR8	0.25±0.16	0.14±0.09	0.22±0.09	0.11±0.02	0.02±0.01	0.06±0.01	0.04±0.01	0.90±0.75	1.75±0.83	5	1.37±0.13
SR9	0.46±0.23	0.83±0.36	0.67±0.19	0.13±0.02	0.09±0.02	0.11±0.09	0.00±0.00	0.56±0.34	2.85±0.71	3	0.43±0.04
SR10	0.15±0.09	0.16±0.08	0.23±0.07	0.05±0.01	0.00±0.01	0.02±0.01	0.00±0.00	0.00±0.04	0.62±0.17	0	0.18±0.04
SR11	0.07±0.05	0.10±0.06	0.09±0.03	0.02±0.00	0.01±0.01	0.01±0.01	0.00±0.00	0.69±0.75	0.98±0.75	1	0.14±0.02
SR12	0.09±0.05	0.23±0.09	0.22±0.06	0.05±0.01	0.03±0.01	0.03±0.01	0.00±0.00	0.14±0.11	0.79±0.20	0	0.35±0.03
SR13	0.10±0.06	0.16±0.08	0.14±0.05	0.04±0.01	0.00±0.00	0.02±0.01	0.00±0.00	0.00±0.05	0.47±0.16	2	0.42±0.04
SR14	0.04±0.02	0.09±0.06	0.09±0.04	0.02±0.01	0.00±0.00	0.02±0.00	0.00±0.00	0.23±0.25	0.48±0.27	0	0.45±0.06

Table 55: Postfit event yields in cut-based regions for 2018.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	1.11±0.28	36.09±6.39	1.40±0.36	0.87±0.11	0.75±0.54	6.53±1.39	0.01±0.00	1.89±0.67	48.65±6.64	51	0.45±0.31
CRW	62.04±14.89	16.59±2.94	17.26±4.06	3.52±0.40	5.90±1.46	5.36±1.19	8.87±1.59	36.98±12.30	156.51±12.23	163	2.02±1.45
SR1	4.15±1.27	1.49±0.29	2.42±0.62	0.60±0.09	0.31±0.10	0.23±0.05	0.24±0.04	1.52±0.84	10.96±1.46	15	1.06±0.83
SR2	1.62±0.57	0.33±0.16	0.75±0.24	0.20±0.04	0.09±0.07	0.10±0.04	0.06±0.01	1.25±0.88	4.39±1.10	4	0.59±0.40
SR3	0.25±0.28	0.06±0.11	0.28±0.11	0.05±0.02	0.01±0.02	0.02±0.00	0.03±0.00	0.05±0.04	0.75±0.48	2	0.37±0.30
SR4	1.88±0.60	0.85±0.23	0.91±0.26	0.33±0.05	0.21±0.10	0.19±0.05	0.15±0.03	0.93±0.38	5.44±0.87	4	0.88±0.61
SR5	0.97±0.30	0.35±0.16	0.49±0.15	0.17±0.03	0.06±0.06	0.15±0.05	0.07±0.01	0.02±0.05	2.28±0.48	3	0.95±0.67
SR6	0.39±0.16	0.11±0.04	0.20±0.08	0.05±0.01	0.04±0.02	0.03±0.01	0.02±0.00	0.43±0.30	1.27±0.36	3	0.66±0.47
SR7	0.13±0.06	0.06±0.03	0.10±0.04	0.02±0.00	0.00±0.00	0.02±0.01	0.00±0.00	0.00±0.05	0.32±0.12	0	0.49±0.36
SR8	0.34±0.19	0.17±0.10	0.24±0.10	0.12±0.02	0.03±0.02	0.07±0.02	0.04±0.01	1.18±0.68	2.18±0.75	5	1.28±0.92
SR9	0.58±0.29	1.00±0.24	0.71±0.18	0.14±0.02	0.09±0.02	0.12±0.15	0.00±0.00	0.60±0.37	3.24±0.64	3	0.39±0.26
SR10	0.18±0.14	0.20±0.11	0.24±0.07	0.06±0.01	0.00±0.02	0.02±0.01	0.00±0.00	0.00±0.04	0.71±0.22	0	0.17±0.15
SR11	0.09±0.08	0.13±0.07	0.10±0.03	0.02±0.01	0.01±0.01	0.01±0.01	0.00±0.00	0.68±0.69	1.03±0.70	1	0.13±0.08
SR12	0.12±0.07	0.28±0.10	0.23±0.07	0.05±0.01	0.03±0.01	0.03±0.01	0.00±0.00	0.14±0.11	0.88±0.22	0	0.32±0.21
SR13	0.14±0.08	0.20±0.10	0.15±0.05	0.05±0.01	0.00±0.00	0.02±0.01	0.00±0.00	0.01±0.05	0.56±0.16	2	0.39±0.27
SR14	0.05±0.03	0.11±0.06	0.10±0.04	0.02±0.01	0.00±0.00	0.02±0.00	0.00±0.00	0.19±0.27	0.48±0.29	0	0.42±0.30

Table 56: Prefit event yields in BDT regions for 2018.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	0.87±0.31	29.83±11.11	1.33±0.35	0.85±0.13	0.71±0.43	6.28±1.56	0.01±0.00	1.72±0.66	41.61±11.56	51	0.48±0.03
SR1	0.49±0.38	0.18±0.10	0.08±0.06	0.01±0.01	0.24±0.20	0.11±0.17	0.29±0.05	0.36±0.36	1.75±0.87	1	0.00±0.00
SR2	1.66±0.88	0.43±0.28	0.38±0.20	0.03±0.02	0.95±0.49	0.49±0.19	0.68±0.13	1.97±1.04	6.60±1.58	8	0.00±0.01
SR3	3.54±1.50	0.89±0.48	0.86±0.31	0.13±0.04	0.97±0.29	0.76±0.24	1.08±0.20	4.74±2.28	12.98±3.15	13	0.02±0.01
SR4	5.09±2.29	1.23±0.70	1.38±0.42	0.21±0.04	0.40±0.47	0.71±0.30	1.17±0.22	3.95±1.41	14.14±3.61	12	0.04±0.01
SR5	6.26±2.38	1.53±1.03	1.68±0.50	0.29±0.06	0.38±0.32	0.61±0.35	1.19±0.22	3.47±1.57	15.40±3.92	17	0.06±0.01
SR6	6.38±2.61	1.83±0.89	1.96±0.54	0.35±0.05	0.03±0.62	0.59±0.21	1.06±0.20	4.10±1.97	16.30±4.32	25	0.10±0.02
SR7	6.05±2.30	2.23±1.25	2.11±0.58	0.42±0.08	0.91±0.20	0.51±0.13	0.87±0.16	4.74±2.16	17.83±3.90	19	0.15±0.03
SR8	7.65±2.83	2.85±1.70	2.84±0.77	0.58±0.08	0.86±0.24	0.59±0.26	0.89±0.17	3.86±1.57	20.12±4.17	21	0.29±0.04
SR9	8.18±3.15	3.26±2.41	3.83±0.98	0.85±0.12	0.93±0.24	0.86±0.27	0.99±0.19	6.14±2.59	25.05±5.46	19	0.45±0.06
SR10	7.28±2.75	2.93±1.82	3.49±0.95	0.89±0.11	0.79±0.33	0.55±0.28	0.69±0.13	5.99±2.66	22.60±4.84	32	0.82±0.04
SR11	5.77±2.33	2.94±1.59	3.47±0.93	0.98±0.12	0.37±0.07	0.51±0.15	0.52±0.10	2.60±1.35	17.14±3.76	35	1.36±0.06
SR12	3.71±1.35	2.00±1.04	2.47±0.68	0.71±0.09	0.33±0.21	0.51±0.15	0.31±0.06	3.04±1.69	13.08±2.83	17	1.63±0.07
SR13	2.10±0.87	1.07±0.68	1.53±0.42	0.48±0.07	0.18±0.04	0.27±0.09	0.13±0.03	1.22±0.51	6.98±1.54	5	1.97±0.09
SR14	0.97±0.44	0.66±0.33	0.89±0.26	0.27±0.04	0.08±0.02	0.13±0.03	0.10±0.02	1.71±1.10	4.81±1.39	6	1.81±0.10
SR15	0.78±0.37	0.41±0.20	0.48±0.14	0.19±0.02	0.16±0.11	0.10±0.02	0.07±0.01	0.19±0.10	2.37±0.59	3	1.75±0.14
SR16	0.17±0.10	0.10±0.09	0.13±0.05	0.08±0.01	0.01±0.00	0.05±0.01	0.01±0.00	0.70±0.47	1.26±0.51	4	1.51±0.11
SR17	0.00±0.02	0.05±0.05	0.02±0.01	0.01±0.00	0.00±0.00	0.01±0.00	0.00±0.00	0.10±0.08	3	0.51±0.07	

Table 57: Postfit event yields in BDT regions for 2018.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	1.01±0.25	36.74±6.82	1.43±0.31	0.89±0.12	0.71±0.41	6.60±1.34	0.01±0.00	1.86±0.67	49.25±6.88	51	0.70±0.29
SR1	0.50±0.23	0.18±0.06	0.07±0.04	0.01±0.00	0.14±0.15	0.09±0.09	0.29±0.05	0.34±0.33	1.62±0.59	1	0.00±0.00
SR2	1.83±0.58	0.45±0.17	0.35±0.14	0.03±0.01	0.85±0.32	0.40±0.15	0.68±0.12	2.30±1.17	6.89±1.34	8	0.00±0.01
SR3	3.82±1.08	0.95±0.29	0.79±0.26	0.11±0.03	0.82±0.28	0.73±0.22	1.08±0.19	5.04±1.99	13.34±2.15	13	0.02±0.02
SR4	5.79±1.59	1.42±0.37	1.36±0.35	0.19±0.04	0.30±0.37	0.71±0.23	1.17±0.21	4.08±1.53	15.02±2.15	12	0.05±0.03
SR5	7.02±1.78	1.76±0.52	1.69±0.42	0.29±0.05	0.53±0.24	0.65±0.24	1.19±0.21	3.74±1.42	16.86±2.16	17	0.08±0.04
SR6	7.41±1.92	2.04±0.52	1.99±0.47	0.35±0.05	0.37±0.46	0.63±0.17	1.06±0.19	4.97±2.09	18.83±2.37	25	0.14±0.08
SR7	6.97±1.75	2.56±0.66	2.15±0.51	0.43±0.07	0.91±0.14	0.51±0.11	0.87±0.15	5.03±1.92	19.43±2.11	19	0.19±0.10
SR8	8.77±2.15	3.18±0.92	2.96±0.69	0.57±0.08							

Table 58: Prefit event yields in cut-based regions for Run2.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	1.91± 0.70	60.40±22.00	2.79± 0.73	1.93± 0.22	0.85± 0.41	12.86± 2.77	0.03± 0.00	2.94± 1.10	83.71±22.75	104	1.05± 0.06
CRW	98.82±35.54	27.53±10.10	34.45± 8.64	7.20± 0.79	13.76± 1.80	9.33± 2.08	18.72± 2.14	59.11±21.59	268.90±45.98	337	4.87± 0.20
SR1	7.00± 2.66	2.59± 0.96	4.72± 1.30	1.18± 0.15	0.63± 0.14	0.60± 0.14	0.48± 0.06	3.37± 1.41	20.57± 3.76	33	2.47± 0.14
SR2	2.49± 0.96	0.53± 0.28	1.49± 0.47	0.38± 0.06	0.14± 0.05	0.16± 0.04	0.10± 0.01	2.23± 1.24	7.53± 1.87	9	1.41± 0.07
SR3	0.49± 0.33	0.15± 0.11	0.59± 0.22	0.12± 0.02	0.01± 0.02	0.04± 0.01	0.04± 0.01	0.19± 0.18	1.64± 0.61	3	0.91± 0.12
SR4	2.74± 1.00	1.22± 0.53	1.76± 0.49	0.59± 0.07	0.43± 0.14	0.29± 0.07	0.32± 0.04	1.86± 0.87	9.21± 1.74	14	1.97± 0.10
SR5	1.51± 0.63	0.47± 0.21	0.91± 0.29	0.30± 0.04	0.10± 0.04	0.26± 0.07	0.10± 0.01	0.40± 0.28	4.05± 0.97	5	2.10± 0.10
SR6	0.52± 0.20	0.13± 0.10	0.39± 0.14	0.10± 0.02	0.12± 0.05	0.05± 0.01	0.02± 0.00	0.41± 0.27	1.75± 0.48	8	1.46± 0.09
SR7	0.22± 0.09	0.10± 0.04	0.18± 0.07	0.04± 0.01	0.01± 0.01	0.03± 0.01	0.01± 0.00	0.00± 0.08	0.59± 0.19	0	1.09± 0.10
SR8	0.47± 0.24	0.21± 0.11	0.38± 0.14	0.21± 0.03	0.05± 0.02	0.11± 0.02	0.08± 0.01	1.36± 0.81	2.86± 0.93	5	2.73± 0.17
SR9	1.03± 0.44	1.79± 0.66	1.46± 0.39	0.30± 0.04	0.15± 0.02	0.16± 0.09	0.00± 0.00	0.86± 0.45	5.75± 1.18	6	0.90± 0.08
SR10	0.23± 0.10	0.48± 0.17	0.52± 0.16	0.11± 0.02	0.01± 0.01	0.04± 0.01	0.00± 0.00	0.35± 0.24	1.73± 0.39	3	0.44± 0.06
SR11	0.13± 0.06	0.25± 0.11	0.19± 0.06	0.04± 0.01	0.01± 0.01	0.02± 0.01	0.00± 0.00	0.74± 0.74	1.38± 0.76	1	0.29± 0.03
SR12	0.15± 0.08	0.48± 0.17	0.39± 0.10	0.09± 0.01	0.04± 0.01	0.06± 0.02	0.00± 0.00	0.57± 0.37	1.77± 0.46	2	0.75± 0.04
SR13	0.20± 0.10	0.28± 0.11	0.28± 0.09	0.07± 0.01	0.02± 0.01	0.04± 0.01	0.00± 0.00	0.00± 0.09	0.89± 0.24	2	0.87± 0.05
SR14	0.11± 0.05	0.18± 0.09	0.17± 0.07	0.04± 0.01	0.00± 0.00	0.03± 0.01	0.00± 0.00	0.23± 0.26	0.76± 0.32	1	1.00± 0.09

Table 59: postfit event yields in cut-based regions for Run2.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	2.74± 0.57	75.94±10.52	3.00± 0.68	1.97± 0.24	0.86± 0.38	13.33± 2.24	0.03± 0.00	3.25± 1.03	101.14±10.08	104	0.84± 0.50
CRW	141.84±27.96	34.54± 4.77	37.01± 8.03	7.31± 0.83	14.04± 1.78	9.71± 1.73	18.78± 2.53	66.93±19.96	330.16±18.88	337	3.88± 2.28
SR1	10.20± 2.18	3.26± 0.47	5.12± 1.17	1.21± 0.15	0.64± 0.12	0.62± 0.12	0.49± 0.07	4.01± 1.61	25.55± 2.11	33	1.98± 1.18
SR2	3.60± 0.86	0.70± 0.22	1.62± 0.41	0.39± 0.06	0.15± 0.05	0.17± 0.03	0.10± 0.01	2.41± 1.05	9.13± 1.27	9	1.13± 0.65
SR3	0.72± 0.34	0.20± 0.10	0.64± 0.20	0.12± 0.02	0.02± 0.02	0.04± 0.01	0.04± 0.01	0.21± 0.17	2.00± 0.59	3	0.73± 0.42
SR4	4.01± 0.93	1.56± 0.34	1.93± 0.46	0.60± 0.07	0.43± 0.13	0.30± 0.06	0.32± 0.04	2.14± 0.85	11.30± 1.26	14	1.58± 0.90
SR5	2.20± 0.58	0.60± 0.13	1.00± 0.27	0.31± 0.05	0.10± 0.04	0.26± 0.05	0.10± 0.02	0.44± 0.27	5.01± 0.77	5	1.69± 0.95
SR6	0.80± 0.21	0.18± 0.08	0.43± 0.13	0.10± 0.02	0.13± 0.05	0.05± 0.01	0.02± 0.00	0.56± 0.29	2.29± 0.40	8	1.20± 0.67
SR7	0.31± 0.11	0.12± 0.04	0.20± 0.06	0.04± 0.01	0.01± 0.01	0.03± 0.01	0.01± 0.00	0.00± 0.08	0.71± 0.20	0	0.89± 0.48
SR8	0.70± 0.28	0.28± 0.12	0.42± 0.15	0.22± 0.03	0.05± 0.02	0.11± 0.02	0.08± 0.01	1.45± 0.82	3.30± 0.95	5	2.21± 1.27
SR9	1.47± 0.42	2.24± 0.34	1.58± 0.36	0.31± 0.05	0.14± 0.02	0.16± 0.09	0.00± 0.00	0.94± 0.46	6.85± 0.80	6	0.72± 0.39
SR10	0.33± 0.11	0.62± 0.14	0.56± 0.14	0.11± 0.02	0.01± 0.01	0.05± 0.01	0.00± 0.00	0.42± 0.26	2.10± 0.31	3	0.36± 0.22
SR11	0.19± 0.07	0.32± 0.07	0.20± 0.06	0.04± 0.01	0.01± 0.01	0.02± 0.01	0.00± 0.00	0.60± 0.72	1.38± 0.75	1	0.23± 0.14
SR12	0.22± 0.08	0.61± 0.12	0.42± 0.10	0.09± 0.01	0.04± 0.01	0.06± 0.01	0.00± 0.00	0.59± 0.40	2.04± 0.48	2	0.59± 0.34
SR13	0.29± 0.11	0.36± 0.12	0.31± 0.09	0.07± 0.01	0.02± 0.01	0.04± 0.01	0.00± 0.00	0.00± 0.11	1.10± 0.28	2	0.70± 0.40
SR14	0.16± 0.05	0.23± 0.07	0.19± 0.06	0.04± 0.01	0.00± 0.00	0.03± 0.01	0.00± 0.00	0.23± 0.27	0.87± 0.30	1	0.81± 0.45

Table 60: Prefit event yields in BDT regions for Run2.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	1.91± 0.70	60.40±19.18	2.79± 0.60	1.93± 0.25	0.85± 0.36	12.86± 2.43	0.03± 0.00	2.94± 0.97	83.71±19.68	104	1.05± 0.06
SR1	1.01± 0.65	0.27± 0.15	0.19± 0.13	0.02± 0.01	0.39± 0.27	0.18± 0.24	0.66± 0.07	0.86± 0.55	4	0.00± 0.00	
SR2	3.40± 1.73	0.81± 0.44	0.79± 0.34	0.07± 0.03	2.14± 0.61	0.84± 0.27	1.52± 0.17	2.72± 1.28	12.31± 2.61	19	0.01± 0.01
SR3	6.73± 3.26	1.79± 0.78	1.70± 0.57	0.25± 0.05	1.61± 0.43	1.12± 0.38	2.26± 0.26	7.91± 2.79	23.35± 5.16	19	0.04± 0.02
SR4	10.36± 4.68	2.44± 1.11	2.80± 0.80	0.38± 0.08	1.64± 0.63	1.17± 0.45	2.51± 0.28	7.87± 2.55	29.17± 6.61	33	0.08± 0.03
SR5	12.19± 5.18	3.01± 1.54	3.41± 0.95	0.57± 0.10	1.44± 0.61	1.21± 0.53	2.59± 0.29	6.65± 2.23	31.08± 7.00	35	0.15± 0.03
SR6	12.42± 5.43	3.66± 1.52	4.07± 1.02	0.74± 0.09	0.94± 0.85	1.34± 0.37	2.17± 0.25	7.74± 2.76	33.08± 7.38	44	0.22± 0.06
SR7	12.18± 5.10	4.43± 2.07	4.38± 1.10	0.80± 0.14	2.17± 0.35	1.01± 0.26	1.86± 0.21	7.27± 2.42	34.09± 6.84	41	0.31± 0.07
SR8	15.09± 6.14	5.59± 2.83	6.03± 1.44	1.19± 0.15	1.60± 0.37	1.19± 0.42	1.95± 0.22	7.17± 2.61	39.80± 8.25	46	0.67± 0.05
SR9	16.78± 7.08	6.74± 3.65	7.87± 1.83	1.85± 0.25	2.11± 0.34	1.54± 0.42	2.02± 0.23	11.11± 4.16	50.01± 10.62	48	1.11± 0.08
SR10	14.71± 6.01	6.14± 3.04	7.45± 1.70	1.85± 0.24	1.43± 0.44	1.12± 0.36	1.40± 0.16	8.86± 3.31	42.96± 8.47	61	1.78± 0.08
SR11	11.80± 5.11	6.03± 2.64	7.20± 1.66	2.00± 0.26	0.95± 0.19	0.91± 0.22	1.07± 0.12	4.80± 2.15	34.75± 7.12	62	2.80± 0.15
SR12	7.32± 3.04	4.14± 1.72	5.18± 1.22	1.51± 0.18	0.69± 0.22	0.88± 0.26	0.61± 0.07	6.41± 2.60	26.76± 5.08	40	3.50± 0.14
SR13	4.14± 1.85	2.38± 1.16	3.13± 0.77	1.00± 0.13	0.25± 0.06	0.51± 0.13	0.28± 0.03	2.31± 0.88	14.00± 2.99	15	3.97± 0.16
SR14	2.04± 1.00	1.30± 0.50	1.76± 0.47	0.54± 0.07	0.11± 0.08	0.30± 0.05	0.18± 0.02	2.21± 1.13	8.43± 1.93	16	3.88± 0.17
SR15	1.37± 0.73	0.67± 0.29	0.95± 0.27	0.37± 0.05	0.20± 0.10	0.17± 0.03	0.12± 0.01	0.44± 0.24	4.27± 1.12	4	3.79± 0.22
SR16	0.42± 0.24	0.20± 0.12	0.31± 0.13	0.14± 0.02	0.07± 0.04	0.10± 0.02	0.03± 0.00	0.84± 0.48	2.12± 0.66	7	3.15± 0.19
SR17	0.03± 0.03	0.07± 0.05	0.03± 0.02	0.02± 0.00	0.00± 0.00	0.02± 0.00	0.00± 0.00	0.25± 0.30	0.44± 0.33	3	1.00± 0.10

Table 61: postfit event yields in BDT regions for Run2.

	t̄W	t̄Z	t̄H	t̄VV	X+γ	Rares	Charge misid.	Nonprompt lep.	SM expected	Data	t̄t̄t̄
CRZ	2.51± 0.54	77.27±11.74	3.09± 0.65	1.98± 0.21	0.87± 0.36	13.29± 2.58	0.03± 0.00	3.25± 1.19	102.30±11.59	104	1.12± 0.43
SR1	1.24± 0.49	0.31± 0.11	0.19± 0.13	0.01± 0.01	0.32± 0.25	0.17± 0.19	0.66± 0.07	1.05± 0.65	3.95± 0.95	4	0.00± 0.00
SR2	4.33± 1.28	0.96± 0.32	0.82± 0.29	0.08± 0.02	2.16± 0.53	0.78± 0.23	1.53± 0.17	3.49± 1.46	14.15± 1.77	19	0.01± 0.01
SR3	8.48± 2.41	2.14± 0.50	1.74± 0.53	0.23± 0.05	1.43± 0.47	1.10± 0.38	2.26± 0.26	8.08± 2.79	25.45± 3.53	19	0.04± 0.03
SR4	13.42± 3.36	3.02± 0.68	2.98± 0.77	0.37± 0.08	1.56± 0.49	1.18± 0.40	2.51± 0.				

**H.3 Nuisances**

- Two sets of nuisance pull values, expected and observed for cut-based and BDT analyses, are tabulated in Table 62 (expected cut-based analysis), Table 63 (observed cut-based analysis), Table 64 (expected BDT analysis), and Table 65 (observed BDT analysis).
- The most constrained nuisances correspond to normalization parameters for ttW and ttZ ("TTWSF" and "TTZSF") due to high statistics in control regions and in the bulk (BDT). Their input normalization uncertainty is 40%.

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Table 62: Expected nuisance pulls for the cut-based analysis. The final column indicates the correlation between the nuisance and the signal strength

name	<i>b</i> -only fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	<i>s + b</i> fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	<b>+0.37, 0.99</b>	+0.00, 0.97	-0.12
TTVV	+0.05, 1.00	+0.00, 0.99	-0.02
TTWSF	<b>+0.09, 0.62</b>	<b>+0.00, 0.62</b>	-0.06
TTZSF	<b>-0.02, 0.45</b>	<b>+0.00, 0.45</b>	+0.02
XG	-0.01, 0.99	+0.00, 0.99	+0.00
alphas	+0.00, 0.99	+0.00, 0.99	+0.00
bb	<b>+0.72, 0.94</b>	+0.00, 0.98	-0.23
fakes	+0.25, 0.91	+0.00, 0.91	-0.09
pdf	+0.07, 0.99	-0.00, 0.99	-0.02
rares	-0.01, 0.99	+0.00, 0.99	+0.00
scale	-0.20, 0.97	-0.00, 0.99	+0.00
y2016_btaghf	+0.04, 0.99	+0.00, 0.99	-0.03
y2016_btagnf	+0.03, 0.99	+0.00, 0.99	-0.01
y2016_fakes_EWK	+0.04, 1.00	+0.00, 0.99	-0.01
y2016_flips	-0.02, 0.99	+0.00, 0.99	+0.00
y2016_fsrvar	+0.00, 0.99	+0.00, 0.99	+0.02
y2016_isr	-0.10, 0.99	-0.00, 0.99	+0.04
y2016_isrvar	+0.00, 0.99	-0.00, 0.99	-0.02
y2016_jer	<b>-0.13, 1.12</b>	+0.00, 0.99	+0.03
y2016_jes	+0.07, 0.97	+0.00, 0.99	-0.05
y2016_lep	+0.01, 0.99	+0.00, 0.99	-0.01
y2016_lumi	+0.02, 0.99	+0.00, 0.99	-0.02
y2016_prefire	-0.00, 0.99	-0.00, 0.99	+0.00
y2016_pu	+0.01, 0.99	-0.00, 0.99	+0.00
y2016_trig	+0.01, 0.99	+0.00, 0.99	-0.02
y2017_btaghf	+0.04, 0.99	+0.00, 0.99	-0.03
y2017_btagnf	+0.03, 0.99	+0.00, 0.99	-0.01
y2017_fakes_EWK	-0.00, 0.99	+0.00, 0.98	-0.00
y2017_flips	-0.02, 0.99	+0.00, 0.99	+0.01
y2017_fsrvar	+0.00, 0.99	-0.00, 0.99	+0.00
y2017_isr	+0.07, 0.99	-0.00, 0.99	-0.03
y2017_isrvar	+0.00, 0.99	-0.00, 0.99	-0.01
y2017_jer	<b>-0.12, 1.20</b>	+0.00, 0.99	+0.03
y2017_jes	+0.07, 1.00	+0.00, 0.99	-0.05
y2017_lep	+0.01, 0.99	+0.00, 0.99	-0.01
y2017_lumi	+0.01, 0.99	+0.00, 0.99	-0.02
y2017_prefire	-0.01, 0.99	-0.00, 0.99	+0.01
y2017_pu	-0.01, 0.99	-0.00, 0.99	+0.00
y2017_trig	+0.01, 0.99	+0.00, 0.99	-0.02
y2018_btaghf	+0.09, 0.99	+0.00, 0.99	-0.05
y2018_btagnf	+0.12, 0.99	+0.00, 0.99	-0.04
y2018_fakes_EWK	+0.08, 1.03	+0.00, 0.99	-0.02
y2018_flips	-0.03, 0.99	+0.00, 0.99	+0.01
y2018_fsrvar	+0.00, 0.99	+0.00, 0.99	-0.00
y2018_isr	+0.08, 0.99	-0.00, 0.99	-0.03
y2018_isrvar	+0.00, 0.99	+0.00, 0.99	-0.00
y2018_jer	<b>-0.32, 1.08</b>	+0.00, 0.99	+0.08
y2018_jes	+0.15, 0.98	+0.00, 0.99	-0.09
y2018_lep	+0.02, 0.99	+0.00, 0.99	-0.02
y2018_lumi	+0.06, 0.99	+0.00, 0.99	-0.07
y2018_pu	+0.00, 1.00	-0.00, 0.99	+0.00
y2018_trig	+0.02, 0.99	+0.00, 0.99	-0.03

Table 63: Observed nuisance pulls for the cut-based analysis. The final column indicates the correlation between the nuisance and the signal strength

name	$b$ -only fit $\Delta x / \sigma_{\text{in}}, \sigma_{\text{out}} / \sigma_{\text{in}}$	$s + b$ fit $\Delta x / \sigma_{\text{in}}, \sigma_{\text{out}} / \sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	<b>+0.49, 1.02</b>	+0.24, 1.00	-0.12
TTVV	+0.06, 1.00	+0.03, 1.00	-0.02
TTWSF	<b>+1.06, 0.58</b>	<b>+1.02, 0.56</b>	-0.05
TTZSF	<b>+0.57, 0.42</b>	<b>+0.61, 0.42</b>	+0.06
XG	+0.04, 1.00	+0.04, 1.00	+0.00
alphas	+0.00, 0.99	+0.00, 0.99	+0.00
bb	<b>+0.62, 0.92</b>	+0.10, 0.97	-0.27
fakes	<b>+0.55, 1.02</b>	<b>+0.34, 1.00</b>	-0.09
pdf	+0.15, 0.98	+0.06, 1.00	-0.04
rares	+0.06, 1.00	+0.07, 1.00	+0.01
scale	-0.09, 0.91	-0.02, 0.96	-0.00
y2016_btaghf	+0.07, 0.99	+0.04, 0.99	-0.03
y2016_btaglf	+0.03, 0.99	+0.00, 0.99	-0.01
y2016_fakes_EWK	+0.08, 1.00	+0.06, 0.99	-0.01
y2016_flips	+0.00, 0.99	+0.01, 1.00	+0.01
y2016_fsrvvar	+0.00, 0.99	-0.01, 1.04	+0.01
y2016_isr	-0.13, 0.99	-0.04, 0.99	+0.05
y2016_isrvvar	+0.00, 0.99	-0.04, 1.06	-0.03
y2016_jer	-0.07, 1.04	+0.01, 0.98	+0.04
y2016_jes	<b>+0.39, 1.07</b>	<b>+0.26, 1.12</b>	-0.08
y2016_lep	+0.04, 0.99	+0.03, 0.99	-0.01
y2016_lumi	+0.05, 1.00	+0.04, 1.00	-0.02
y2016_prefire	-0.01, 0.99	-0.00, 0.99	+0.00
y2016_pu	+0.03, 0.99	+0.01, 0.99	-0.01
y2016_trig	+0.04, 0.99	+0.04, 0.99	-0.01
y2017_btaghf	+0.10, 1.00	+0.06, 0.99	-0.03
y2017_btaglf	+0.03, 0.99	+0.00, 0.99	-0.01
y2017_fakes_EWK	+0.24, 0.98	+0.22, 0.97	-0.01
y2017_flips	+0.00, 0.99	+0.01, 1.00	+0.01
y2017_isrvvar	+0.00, 0.99	-0.01, 0.99	+0.00
y2017_isr	+0.08, 0.99	+0.02, 0.99	-0.04
y2017_isrvvar	+0.00, 0.99	-0.01, 0.99	-0.01
y2017_jer	<b>-0.07, 1.24</b>	+0.01, 0.96	+0.03
y2017_jes	+0.14, 0.93	+0.09, 0.93	-0.05
y2017_lep	+0.03, 0.99	+0.03, 0.99	-0.01
y2017_lumi	+0.05, 1.00	+0.04, 1.00	-0.01
y2017_prefire	-0.02, 0.99	-0.02, 0.99	+0.01
y2017_pu	+0.02, 1.00	+0.01, 1.00	-0.00
y2017_trig	+0.05, 0.99	+0.04, 0.99	-0.01
y2018_btaghf	+0.15, 0.99	+0.08, 0.99	-0.05
y2018_btaglf	+0.12, 0.99	+0.03, 0.99	-0.05
y2018_fakes_EWK	+0.15, 1.01	+0.09, 0.99	-0.03
y2018_flips	+0.01, 0.99	+0.02, 1.00	+0.01
y2018_isrvvar	+0.00, 0.99	+0.00, 0.99	+0.00
y2018_isr	+0.08, 0.99	+0.02, 0.99	-0.04
y2018_isrvvar	+0.00, 0.99	+0.00, 0.99	+0.00
y2018_jer	<b>-0.19, 1.16</b>	+0.06, 1.05	+0.09
y2018_jes	<b>+0.48, 1.03</b>	<b>+0.30, 1.08</b>	-0.13
y2018_lep	+0.06, 0.99	+0.05, 0.99	-0.02
y2018_lumi	+0.20, 0.99	+0.17, 0.99	-0.06
y2018_pu	+0.00, 1.00	-0.00, 0.99	+0.00
y2018_trig	+0.09, 0.99	+0.07, 0.99	-0.02

Table 64: Expected nuisance pulls for the BDT analysis. The final column indicates the correlation between the nuisance and the signal strength

name	$b$ -only fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$s + b$ fit $\Delta x/\sigma_{\text{in}}, \sigma_{\text{out}}/\sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	<b>+0.31, 0.98</b>	-0.00, 0.96	-0.11
TTVV	+0.06, 1.00	-0.00, 0.99	-0.02
TTWSF	<b>+0.07, 0.64</b>	<b>-0.00, 0.63</b>	-0.06
TTZSF	<b>-0.03, 0.50</b>	<b>+0.00, 0.48</b>	+0.03
XG	-0.01, 0.99	+0.00, 0.99	+0.00
alphas	+0.00, 0.99	+0.00, 0.99	+0.00
bb	<b>+0.89, 0.94</b>	-0.00, 0.99	-0.26
fakes	<b>+0.31, 0.90</b>	<b>+0.00, 0.90</b>	-0.09
pdf	+0.12, 0.97	-0.00, 0.99	-0.03
rares	-0.02, 0.99	+0.00, 0.99	+0.00
scale	-0.25, 0.98	+0.00, 0.99	+0.01
y2016_btaghf	+0.04, 0.99	+0.00, 0.99	-0.02
y2016_btaglf	+0.03, 0.99	-0.00, 0.99	-0.01
y2016_fakes_EWK	+0.03, 0.99	-0.00, 0.99	-0.01
y2016_flips	-0.02, 0.99	+0.00, 0.99	+0.01
y2016_fsrvvar	+0.00, 0.99	+0.00, 0.99	+0.01
y2016_isr	-0.13, 0.99	+0.00, 0.99	+0.04
y2016_isrvvar	+0.00, 0.99	-0.00, 0.99	-0.01
y2016_jer	-0.04, 1.03	+0.00, 0.99	+0.01
y2016_jes	+0.07, 0.90	+0.00, 0.99	-0.04
y2016_lep	+0.01, 0.99	+0.00, 0.99	-0.01
y2016_lumi	+0.01, 0.99	+0.00, 0.99	-0.02
y2016_prefire	-0.00, 0.99	-0.00, 0.99	+0.00
y2016_pu	+0.04, 0.99	-0.00, 0.99	-0.01
y2016_trig	+0.01, 0.99	+0.00, 0.99	-0.02
y2017_btaghf	+0.03, 0.99	+0.00, 0.99	-0.03
y2017_btaglf	+0.03, 0.99	+0.00, 0.99	-0.01
y2017_fakes_EWK	+0.03, 0.98	+0.00, 0.96	-0.01
y2017_flips	-0.02, 0.99	+0.00, 0.99	+0.00
y2017_fsrvvar	+0.00, 0.99	+0.00, 0.99	+0.00
y2017_isr	+0.07, 0.99	-0.00, 0.99	-0.02
y2017_isrvvar	+0.00, 0.99	-0.00, 0.99	-0.01
y2017_jer	<b>-0.00, 1.18</b>	+0.00, 0.99	+0.00
y2017_jes	+0.07, 0.90	+0.00, 0.99	-0.04
y2017_lep	+0.00, 0.99	+0.00, 0.99	-0.01
y2017_lumi	+0.00, 0.99	+0.00, 0.99	-0.02
y2017_prefire	-0.00, 0.99	-0.00, 0.99	+0.01
y2017_pu	+0.03, 1.00	-0.00, 0.99	-0.01
y2017_trig	+0.00, 0.99	+0.00, 0.99	-0.02
y2018_btaghf	+0.08, 0.99	+0.00, 0.99	-0.05
y2018_btaglf	+0.11, 0.99	-0.00, 0.99	-0.04
y2018_fakes_EWK	+0.06, 0.98	+0.00, 0.98	-0.01
y2018_flips	-0.03, 0.99	+0.00, 0.99	+0.01
y2018_fsrvvar	+0.00, 0.99	+0.00, 0.99	+0.00
y2018_isr	+0.04, 0.99	-0.00, 0.99	-0.02
y2018_isrvvar	+0.00, 0.99	+0.00, 0.99	+0.00
y2018_jer	<b>+0.01, 1.49</b>	+0.00, 0.95	-0.00
y2018_jes	<b>+0.12, 0.85</b>	+0.00, 0.97	-0.08
y2018_lep	+0.01, 0.99	+0.00, 0.99	-0.02
y2018_lumi	+0.05, 0.99	+0.00, 0.99	-0.07
y2018_pu	+0.02, 1.01	-0.00, 0.99	-0.01
y2018_trig	+0.02, 0.99	+0.00, 0.99	-0.03

Table 65: Observed nuisance pulls for the BDT analysis. The final column indicates the correlation between the nuisance and the signal strength

name	$b$ -only fit $\Delta x / \sigma_{\text{in}}, \sigma_{\text{out}} / \sigma_{\text{in}}$	$s + b$ fit $\Delta x / \sigma_{\text{in}}, \sigma_{\text{out}} / \sigma_{\text{in}}$	$\rho(\theta, \mu)$
TTH	<b>+0.59, 1.02</b>	<b>+0.34, 1.01</b>	-0.09
TTVV	+0.08, 1.00	+0.04, 1.00	-0.02
TTWSF	<b>+0.65, 0.67</b>	<b>+0.77, 0.60</b>	-0.00
TTZSF	<b>+0.62, 0.43</b>	<b>+0.66, 0.43</b>	+0.05
XG	+0.02, 0.99	+0.04, 1.00	+0.01
alphas	+0.00, 0.99	+0.00, 1.00	+0.00
bb	<b>+0.88, 0.95</b>	+0.07, 0.98	-0.27
fakes	<b>+0.95, 0.99</b>	<b>+0.34, 0.99</b>	-0.13
pdf	+0.04, 1.00	-0.04, 1.01	-0.02
rares	+0.03, 0.99	+0.04, 0.99	+0.00
scale	-0.25, 0.94	-0.05, 0.94	+0.00
y2016_btaghf	+0.08, 0.99	+0.05, 0.99	-0.02
y2016_btaglf	+0.04, 1.00	+0.01, 0.99	-0.01
y2016_fakes_EWK	+0.17, 0.99	+0.08, 0.99	-0.02
y2016_flips	-0.01, 0.99	+0.00, 1.00	+0.00
y2016_fsrvar	+0.00, 0.99	+0.02, 1.07	+0.02
y2016_isr	-0.19, 0.99	-0.07, 0.99	+0.04
y2016_isrvar	+0.00, 0.99	+0.01, 1.00	-0.01
y2016_jer	<b>-0.06, 0.89</b>	<b>-0.03, 0.87</b>	+0.01
y2016_jes	+0.04, 0.99	<b>-0.04, 1.11</b>	-0.03
y2016_lep	+0.04, 0.99	+0.03, 0.99	-0.01
y2016_lumi	+0.06, 1.00	+0.05, 1.00	-0.02
y2016_prefire	-0.01, 0.99	-0.00, 0.99	+0.00
y2016_pu	-0.01, 1.00	-0.05, 1.00	-0.01
y2016_trig	+0.05, 0.99	+0.04, 0.99	-0.02
y2017_btaghf	+0.09, 0.99	+0.08, 0.99	-0.03
y2017_btaglf	+0.05, 0.99	+0.02, 0.99	-0.01
y2017_fakes_EWK	<b>+0.44, 1.03</b>	+0.26, 0.98	-0.04
y2017_flips	-0.01, 0.99	+0.01, 1.00	+0.01
y2017_fsrvar	+0.00, 0.99	+0.02, 1.00	+0.01
y2017_isr	+0.08, 0.99	+0.02, 0.99	-0.02
y2017_isrvar	+0.00, 0.99	+0.00, 1.00	-0.01
y2017_jer	<b>+0.09, 0.88</b>	<b>+0.07, 0.84</b>	-0.01
y2017_jes	<b>+0.15, 0.90</b>	+0.09, 1.00	-0.04
y2017_lep	+0.05, 0.99	+0.04, 0.99	-0.01
y2017_lumi	+0.05, 1.00	+0.05, 1.00	-0.02
y2017_prefire	-0.01, 0.99	-0.01, 0.99	+0.01
y2017_pu	+0.05, 1.00	+0.03, 1.00	-0.01
y2017_trig	+0.04, 0.99	+0.05, 0.99	-0.01
y2018_btaghf	+0.14, 0.99	+0.09, 0.99	-0.05
y2018_btaglf	+0.06, 0.99	-0.03, 0.99	-0.04
y2018_fakes_EWK	+0.10, 0.98	-0.01, 0.99	-0.03
y2018_flips	-0.02, 0.99	+0.00, 0.99	+0.01
y2018_fsrvar	+0.00, 0.99	+0.00, 0.99	-0.00
y2018_isr	+0.04, 0.99	+0.01, 0.99	-0.02
y2018_isrvar	+0.00, 0.99	+0.00, 0.99	-0.00
y2018_jer	<b>+0.74, 0.77</b>	<b>+0.50, 0.97</b>	-0.14
y2018_jes	<b>+0.21, 0.67</b>	<b>+0.14, 0.81</b>	-0.07
y2018_lep	+0.06, 0.99	+0.05, 0.99	-0.02
y2018_lumi	+0.16, 0.99	+0.15, 0.99	-0.06
y2018_pu	+0.09, 1.01	+0.06, 0.99	-0.02
y2018_trig	+0.06, 0.99	+0.05, 0.99	-0.03

#### 1806 H.4 Impacts

1807 The leading 30 nuisance impacts for two sets of impacts, expected and observed for cut-based  
 1808 and BDT analyses, are shown in Figure 117 (expected cut-based analysis), Figure 118 (observed  
 1809 cut-based analysis), Figure 119 (expected BDT analysis), and Figure 120 (observed BDT anal-  
 1810 ysis). The leading expected nuisance in both cases corresponds to the  $\sigma(\text{ttbb})/\sigma(\text{ttjj})$  scaling.  
 1811 Note that the “prop binSS” nuisances for MC statistics include (and are dominated by) tight-  
 1812 loose sideband statistics.

1813 The obseved pulls show the most constrained/pulled nuisances correspond to normalization  
 1814 parameters for ttW and ttZ, as we would expect from the control regions. “TTWSF” is moved  
 1815 by approximately  $1\sigma$  ( $0.8\sigma$ ) with respect to the input nuisance sizes for the cut-based (BDT)  
 1816 analysis. “TTZSF” is moved up by approximately  $0.6\sigma$  ( $0.7\sigma$ ) for the cut-based (BDT) analysis.

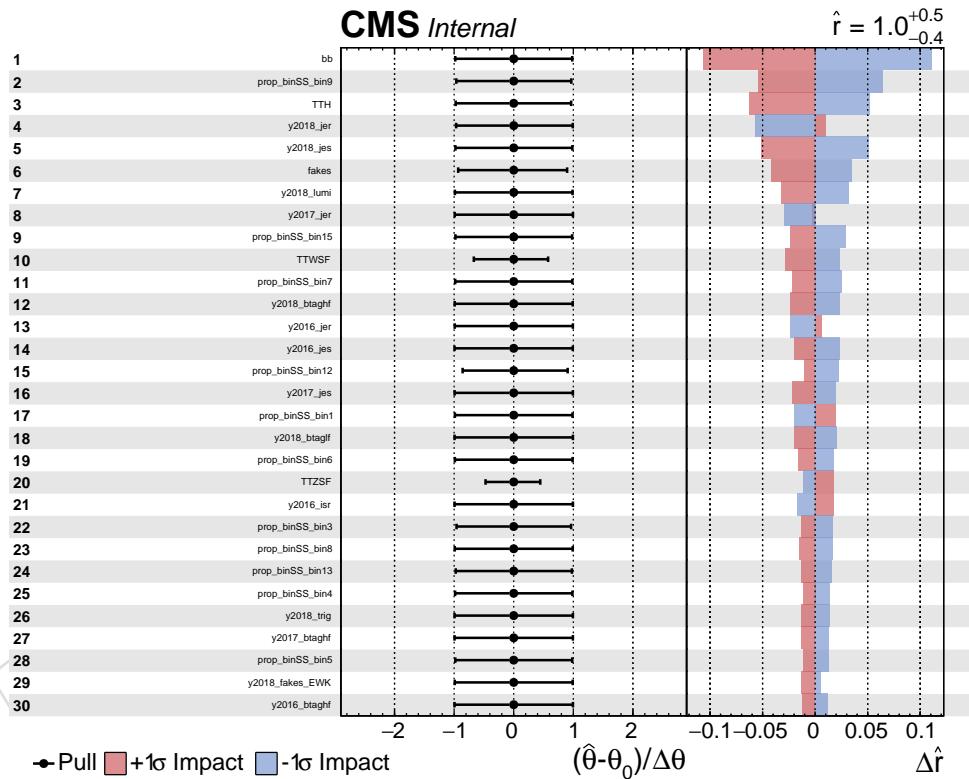


Figure 117: Expected nuisance impacts for the cut-based analysis.

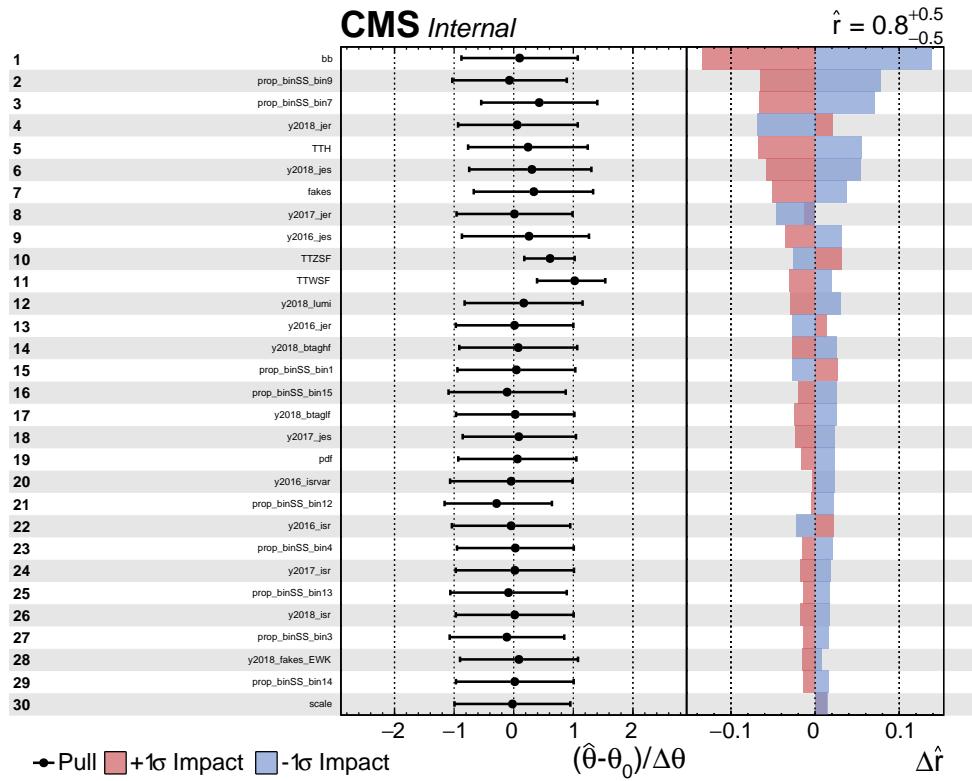


Figure 118: Observed nuisance impacts for the cut-based analysis.

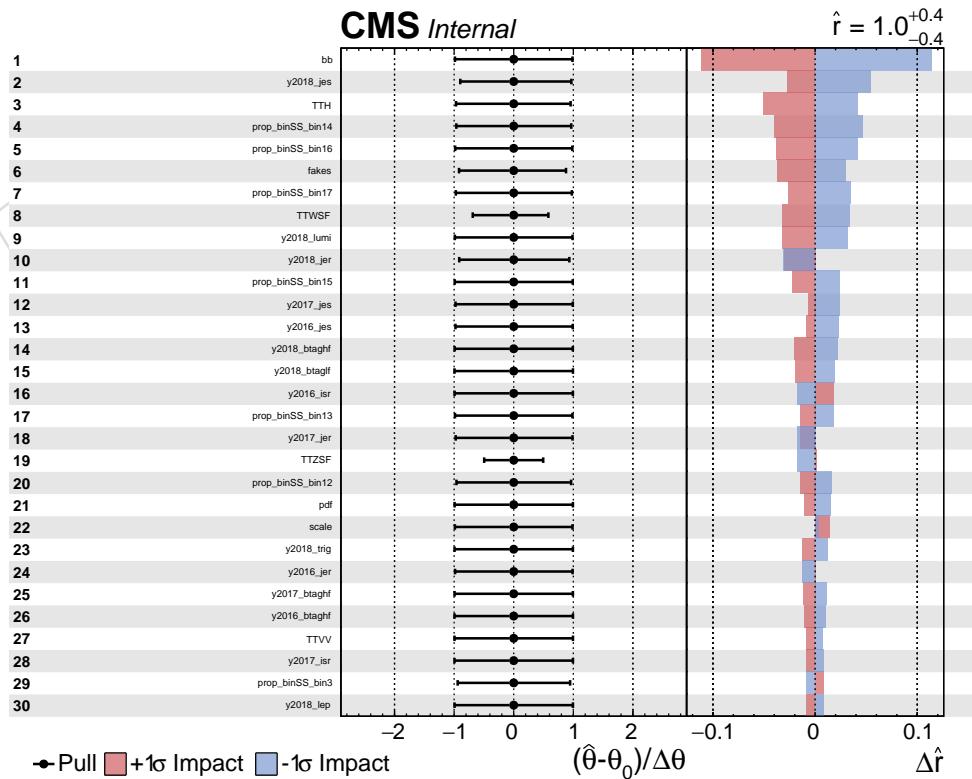


Figure 119: Expected nuisance impacts for the BDT-based analysis.

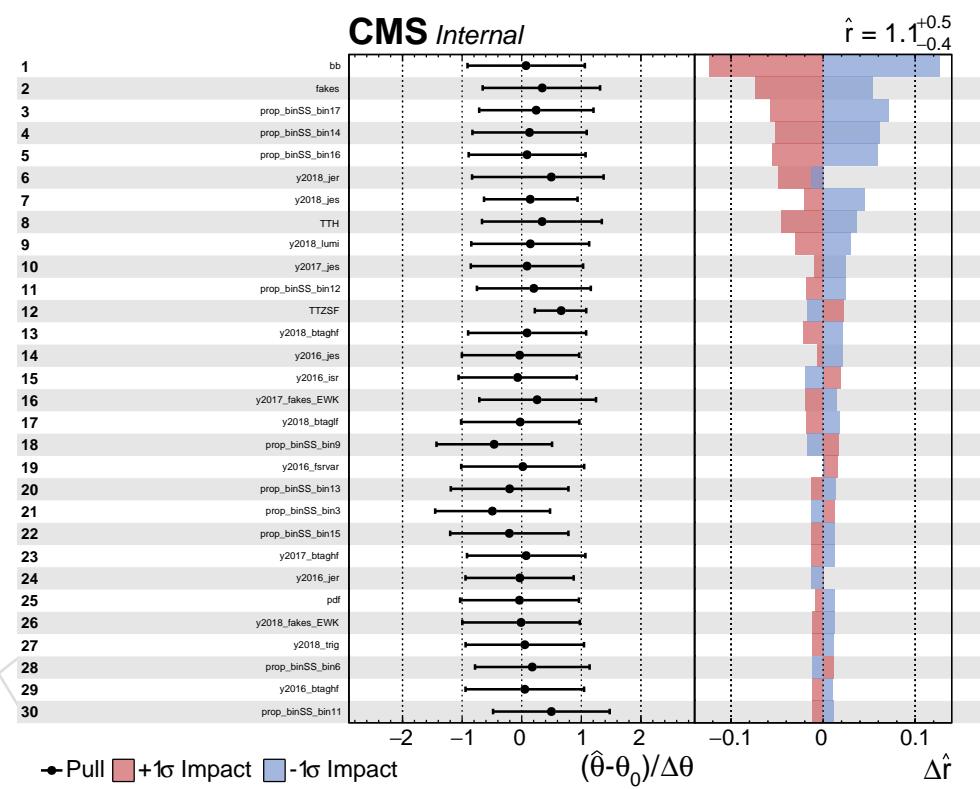


Figure 120: Observed nuisance impacts for the BDT-based analysis.

1817 **H.5 Goodness of fits**

1818 The goodness of fit distributions (using the saturated, Kolmogorov-Smirnov, and Anderson-  
 1819 Darling test statistics) with the signal+background fit to data for the cut-based and BDT anal-  
 1820 yses are shown in Figure 121. We note that the observation is generally within the bulk of the  
 expected distributions for both analyses and all three fit types.

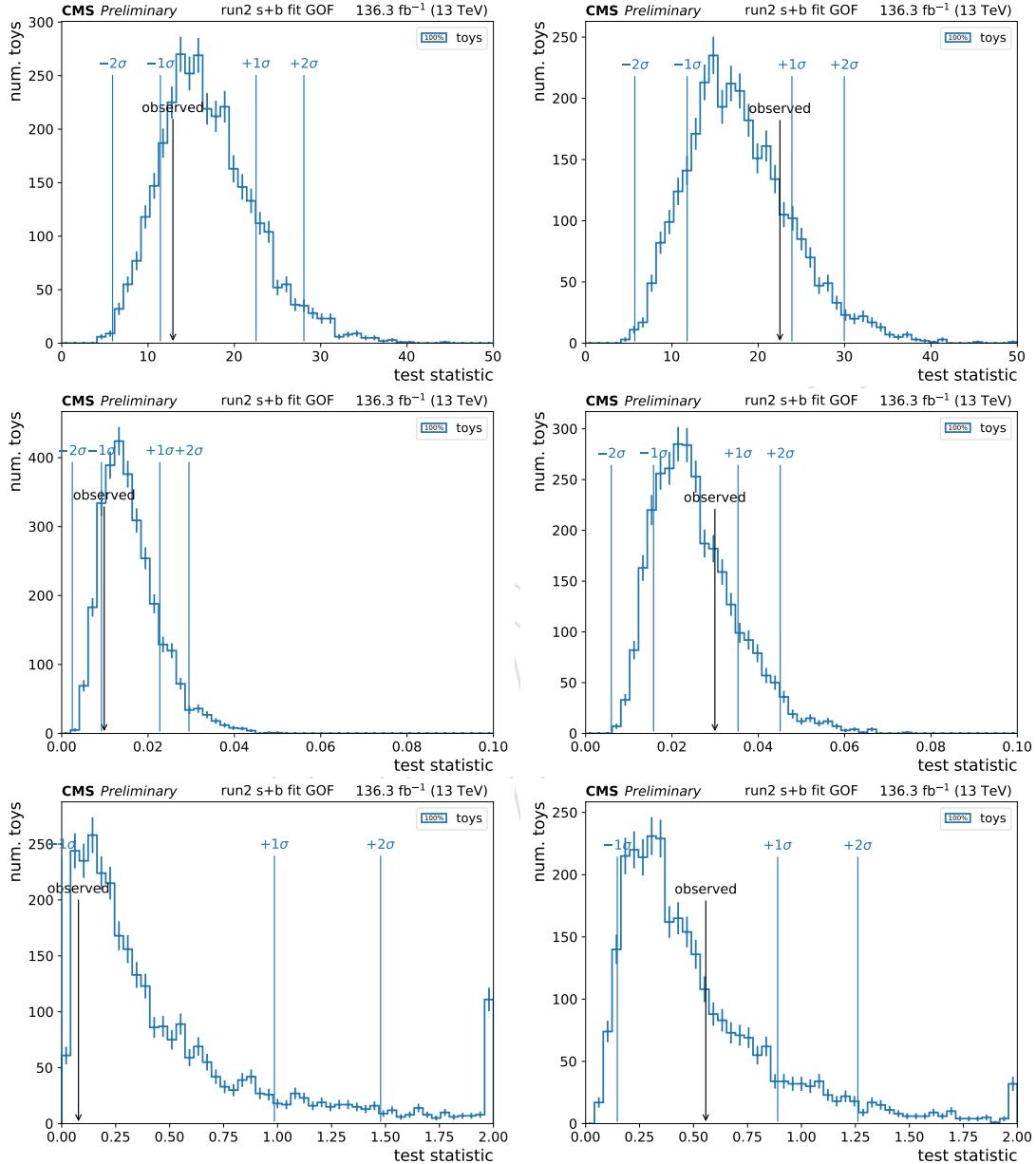


Figure 121: GOF test for the cut-based analysis (left) and BDT-based analysis (right) using the saturated (top), KS (middle), AD (bottom) test statistics.

Table 66: Event kinematics for tail BDT events

run:lumievent	SRnum	year	nleps	id1	id2	pT1	pT2	eta1	eta2	phi1	phi2	Nj	Nb	MET	HT
2767761569261539766	17	2016	2	11	13	131.313	71.6004	1.54234	0.574022	-0.434035	-1.93809	7	3	256.921	1458.27
2830505892826095	17	2016	2	-13	-11	48.3626	39.4692	-0.217067	-2.34033	2.24725	-0.568915	7	3	168.629	844.845
3053779441751515966	17	2017	2	13	13	146.593	112.307	-1.86365	0.71062	2.00505	0.010847	7	3	212.79	836.855
31564487101203005	17	2018	2	-13	-13	102.218	24.2224	-1.61099	-0.637295	1.11559	-1.64171	7	5	82.6309	739.944
3162409151292897676	17	2018	2	11	13	46.3985	30.3899	0.592867	0.175213	0.0199213	3.82597	7	4	56.2346	632.974
3174351373196950764	17	2018	3	-11	-11	62.5448	38.0043	0.172945	0.109995	2.11431	-0.932109	9	2	95.8601	1622.7
3234000000000000000	18	2018	2	13	13	57.4322	50.6672	1.42404	0.887908	2.50299	-1.42865	8	4	79.4541	558.675
3217745085657761	18	2018	2	-11	-11	112.112	107.589	1.17477	0.895216	1.18255	-2.27924	9	4	245.434	1466.6
32235679153159025	18	2018	2	-11	-13	54.3049	31.4315	-0.017497	-0.924084	1.47584	1.63024	8	4	174.04	650.043
323727421771405322	17	2018	2	11	11	55.2933	42.2259	-0.0188495	-1.35181	2.1227	-1.93872	7	3	126.791	592.848

## 1822 H.6 Tail BDT events

1823 We included event information about the 7+3 events in the two tail BDT bins in Table 66. As  
 1824 expected, the high number of jets and b-tag multiplicity contribute to a high BDT score. We  
 1825 don't see any localization in the detector (for example, due to HEM).

DRAFT

## 1826 I Checks from ARC review

### 1827 I.1 Additional uncertainty from tt+bb

1828 Figure 126 shows the fraction of events affected by the correction to account for the measured  
 1829 ratio of  $t\bar{t}bb/t\bar{t}jj$  cross sections ( $1.7 \pm 0.6$ ), as discussed in Section 2.  $t\bar{t}$  and  $t\bar{t}Z$  show agreement  
 1830 in this fraction (approximately 1%), while  $t\bar{t}W$  is low by 30%. To study possible effects due  
 1831 to differences in kinematics with respect to  $t\bar{t}$ , we tested a configuration where an additional  
 1832 30% uncertainty was included in quadrature, so that the ratio and corresponding systematic  
 1833 uncertainty become  $1.7 \pm 0.78$ . The effect on the analysis was to reduce the observed signal  
 1834 strength of the cut-based and BDT analyses by 0.5%, and the expected (observed) significance  
 1835 of both analyses by 1% (1.5%). This additional uncertainty was eventually not adopted in the  
 1836 analysis.

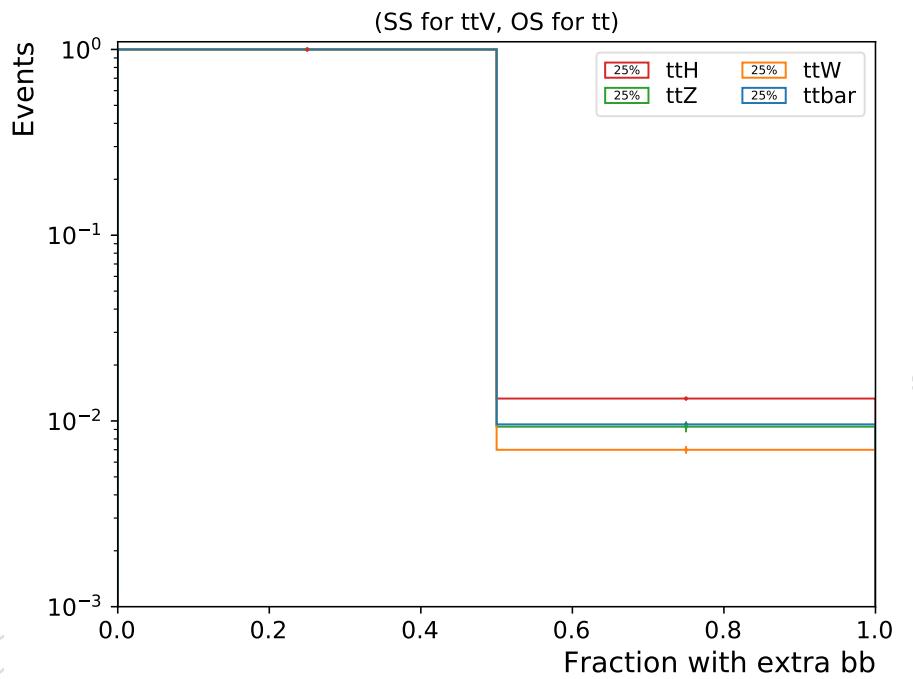


Figure 122: Fraction of events reweighted by measured ratio of  $t\bar{t}bb/t\bar{t}jj$  for  $t\bar{t}$ ,  $t\bar{t}W$ , and  $t\bar{t}Z$ , inclusively. Note that the same-sign requirement has been lifted from  $t\bar{t}$ , which requires opposite-sign dileptons.

## I.2 Distributions of BDT inputs for signal region events

Figure 123 shows distributions of the 19 BDT input variables for the same events that enter the BDT signal regions, except for CRZ. They are shown with “postfit yields”, meaning prefit shapes are used, but normalization of processes is adjusted to match postfit/prefit scale factors. Good agreement is observed.

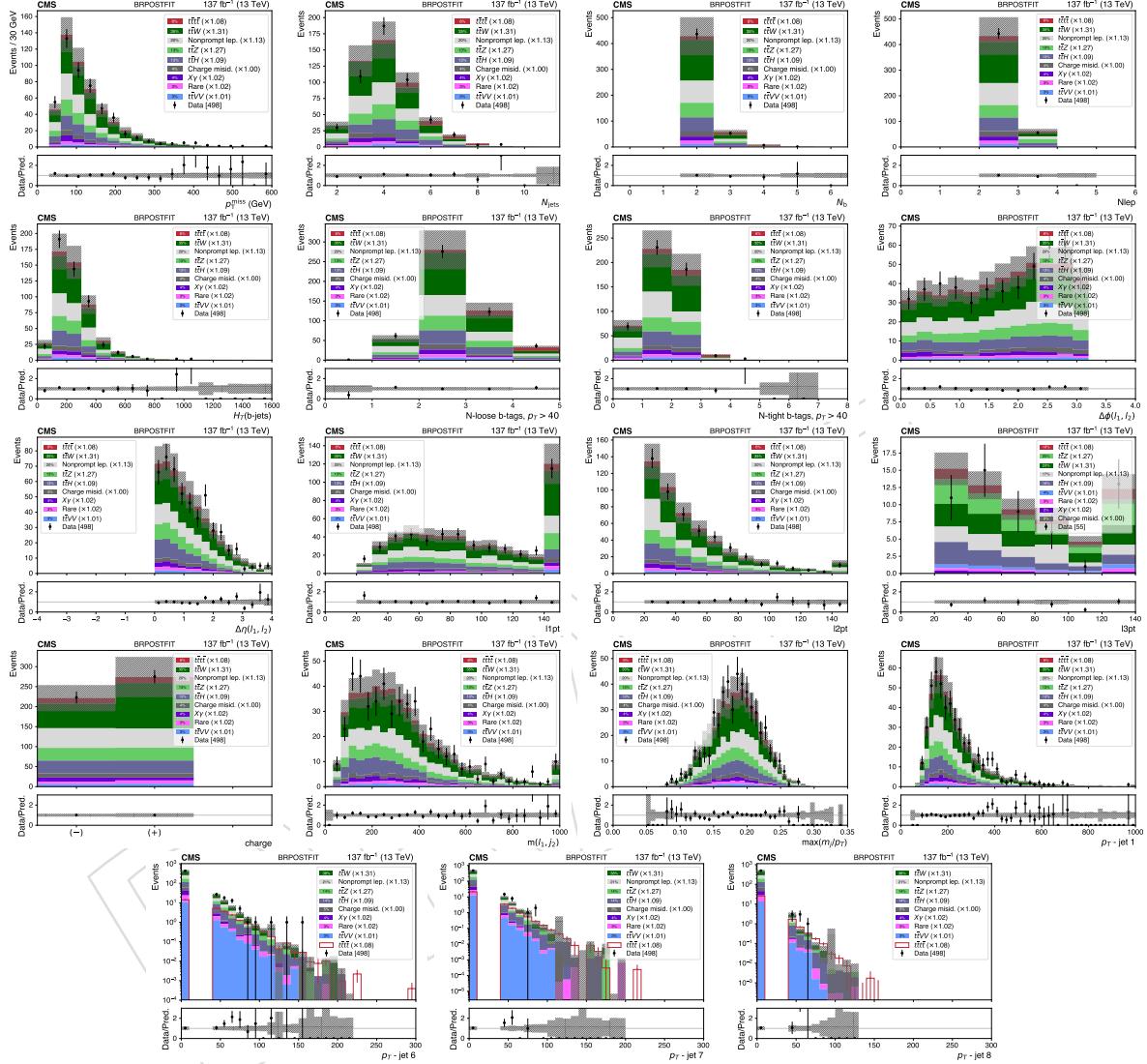


Figure 123: Run2 Data to prediction comparisons for the BDT input variables. From left top to right bottom,  $\not{E}_T$ ,  $N_{\text{jets}}$ ,  $N_b$  jets,  $H_T^b$ ,  $N_{\text{looseb}}$ ,  $N_{\text{tightb}}$ ,  $\Delta\phi(\ell_1, \ell_2)$ ,  $\Delta\eta(\ell_1, \ell_2)$ ,  $p_T(\ell_1)$ ,  $p_T(\ell_2)$ ,  $p_T(\ell_3)$ ,  $q_1$ ,  $m(\ell_1, j_1)$ ,  $\max(m(j)/p_T(j))$  and the  $p_T$  for jets 1,6,7, and 8, are shown for the BDT signal region events excluding CRZ, with scaled  $t\bar{t}t\bar{t}$  signal stacked when in solid red.

## I.3 More jets at higher $\eta$

The analysis has a jet requirement of  $|\eta| < 2.4$  to keep (b)jets almost completely within tracker acceptance. Additionally, b-tag scale factor uncertainties are provided up to  $|\eta|$  of 2.4 (2.5) for 2016 (2017, 2018) according to <https://twiki.cern.ch/twiki/bin/view/CMS/BtagRecommendation>, and reflected in the .csv files.

1848 While optimizing the Run 2 analysis, we considered relaxing both the  $p_T$  and  $|\eta|$  requirements  
 1849 for jets and b-jets. First, we studied the additional acceptance for signal jets when relaxing these  
 1850 cuts. The additional jet acceptance from relaxing the  $|\eta|$  requirement is small. Figure 124 shows  
 1851 the  $|\eta|$  distribution of  $p_T > 30\text{GeV}$  generator-level jets which are matched to top decay products  
 1852 (light-flavor quarks and b quarks) in  $t\bar{t}t\bar{t}$  events with 2 leptonically-decaying W bosons.  
 1853 The left plot shows an additional 3.8% of jets would be gained by completely relaxing the  $|\eta|$   
 1854 requirement up to approximately 3.6. The right plot shows up to 0.8% of jets would be gained  
 1855 by relaxing  $|\eta|$  to 2.5, in order to stay within an appropriate range of the b-tag scale factors and  
 1856 allow for well-behaved JEC factors. Figure 125 shows the significantly more impressive gain  
 1857 from relaxing  $p_T$  thresholds instead of  $\eta$ .

1858 As the gain from relaxing  $p_T$  thresholds was much larger, we put aside the  $\eta$  cut and focused on  
 1859 the  $p_T$  cut. This resulted in the studies shown in Appendix C, where the whole analysis was re-  
 1860 run with varying  $p_T$  cuts, signal and background counts were re-evaluated, but the significance  
 1861 was not found to improve. For these reasons, we decided not to change our thresholds with  
 1862 respect to the 2016 analysis.

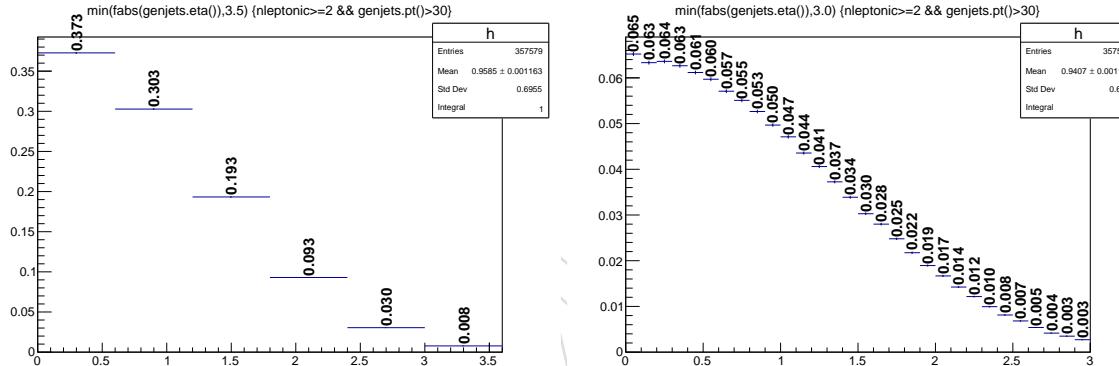


Figure 124:  $|\eta|$  distribution of generator-level jets which are matched to top decay products for 2 lepton  $t\bar{t}t\bar{t}$  events with coarse-binning (left) and fine-binning (right). The plots are normalized to unit area and bin values are superimposed.

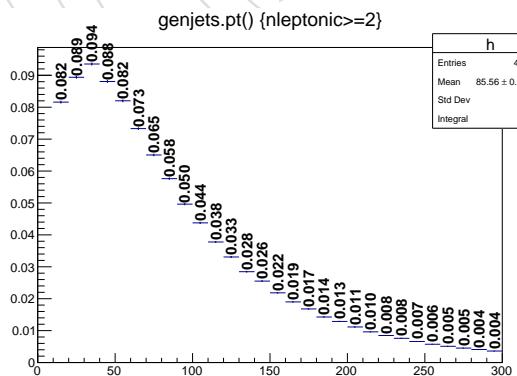


Figure 125:  $p_T$  distribution of generator-level jets which are matched to top decay products for 2 lepton  $t\bar{t}t\bar{t}$  events.

#### 1863 I.4 Effect of ttH and ttbb/ttjj measurement updates

1864 In this section, we quantify the effects of potential updates to the  $t\bar{t}H$  and  $ttbb$  measurements  
 1865 on this analysis.

1866 **I.4.1 ttH**

1867 As we do not scale the  $t\bar{t}H$  background, an updated measurement of  $t\bar{t}H$  would only change  
 1868 the 25% relative uncertainty we place on this process. The effect of modifying this uncertainty  
 1869 can be seen in the impact plots from Appendix A, which shows the  $t\bar{t}H$  normalization uncer-  
 1870 tainty nuisance is up to a 6% effect on the measured signal strength, approximately. Changes  
 1871 with respect to the 25% normalization uncertainty can then be extrapolated from this 6%.

1872 **I.4.2  $t\bar{t}bb/t\bar{t}jj$** 

1873 Conversely, an updated measurement of  $t\bar{t}bb/t\bar{t}jj$  can change the central value of background  
 1874 predictions. Currently, the  $t\bar{t}bb/t\bar{t}jj$  ratio is  $1.7 \pm 0.6$ . Based on TOP-18-002 (<http://cms.cern.ch/analysisadmin/cadilines?line=TOP-18-002>), the  $t\bar{t}bb/t\bar{t}jj$  ratio in  
 1875 the visible phase space is  $1.17 \pm 0.16$  for the lepton+jets channels, and  $1.38 \pm 0.24$  for the dilep-  
 1876 ton channels. While a numerical combination is not provided, a crude estimate using the av-  
 1877 erage would be around  $1.30 \pm 0.20$ . Thus, we use  $1.3 \pm 0.2$  and re-evaluate the results of this  
 1878 analysis. Numbers for both values/uncertainties of  $t\bar{t}bb/t\bar{t}jj$  are tabulated in Table 67. The  
 1879 smaller scale factor from  $t\bar{t}bb/t\bar{t}jj$  results in 10% higher expected and observed significance  
 1880 values on average.

BDT			
$t\bar{t}bb/t\bar{t}jj$	exp. $\sigma$	obs. $\sigma$	obs. $\mu$
$1.7 \pm 0.6$	2.699	2.561	$1.052^{+0.483}_{-0.438}$
$1.3 \pm 0.2$	2.904	2.841	$1.108^{+0.470}_{-0.420}$

Cut-based			
$t\bar{t}bb/t\bar{t}jj$	exp. $\sigma$	obs. $\sigma$	obs. $\mu$
$1.7 \pm 0.6$	2.478	1.713	$0.784^{+0.514}_{-0.469}$
$1.3 \pm 0.2$	2.629	1.968	$0.852^{+0.498}_{-0.450}$

Table 67

1882 **I.5 tt+bb correction applied to ttV samples**

1883 Figure 126 shows the fraction of events affected by the correction to account for the measured  
 1884 ratio of  $t\bar{t}bb/t\bar{t}jj$  cross sections ( $1.7 \pm 0.6$ ), as discussed in Section 2.  $t\bar{t}$  and  $t\bar{t}Z$  show agreement  
 1885 in this fraction (approximately 1%), while  $t\bar{t}W$  is low by 30%. To study possible effects due  
 1886 to differences in kinematics with respect to  $t\bar{t}$ , we tested a configuration where an additional  
 1887 30% uncertainty was included in quadrature, so that the ratio and corresponding systematic  
 1888 uncertainty become  $1.7 \pm 0.78$ . The effect on the analysis was to reduce the observed signal  
 1889 strength of the cut-based and BDT analyses by 0.5%, and the expected (observed) significance  
 1890 of both analyses by 1% (1.5%). This additional uncertainty was eventually not adopted in the  
 1891 analysis.

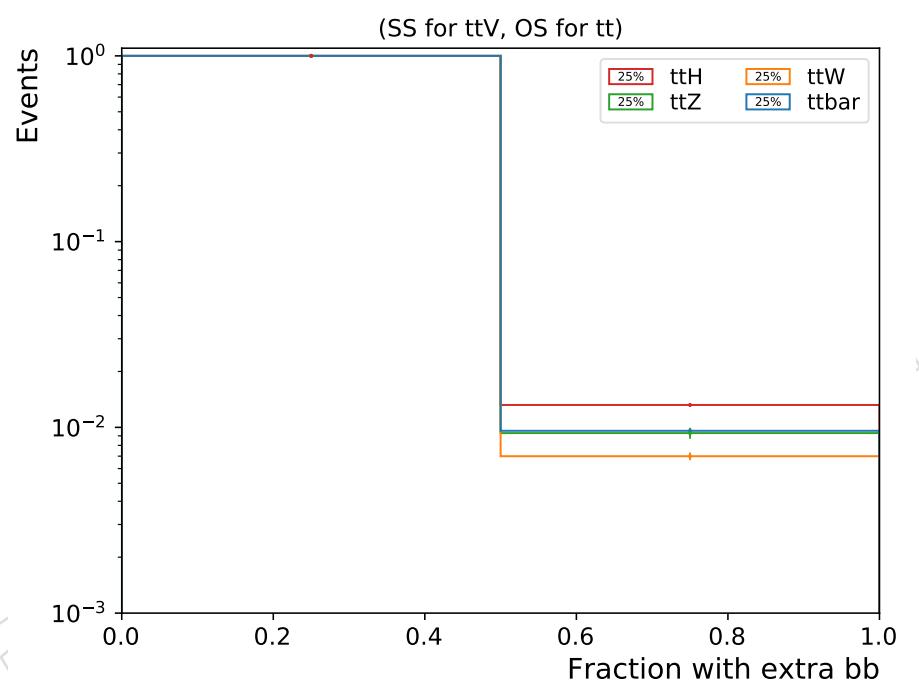


Figure 126: Fraction of events reweighted by measured ratio of  $t\bar{t}bb\bar{b}/t\bar{t}jj$  for  $t\bar{t}$ ,  $t\bar{t}W$ , and  $t\bar{t}Z$ , inclusively. Note that the same-sign requirement has been lifted from  $t\bar{t}$ , which requires opposite-sign dileptons.

1892 **I.6  $N_J^{\text{ISR/FSR}}$  systematic**

1893 As introduced in Section 2,  $t\bar{t}W$  and  $t\bar{t}Z$  MC events are reweighted based on  $N_J^{\text{ISR/FSR}}$ , and a  
 1894 systematic corresponding to half of the reweighting factor deviation from unity is used (Sec-  
 1895 tion 9).

1896 The size of the systematic comes from deviations observed in an orthogonal single-lepton chan-  
 1897 nel.

1898 The tunes for 2016 (TuneCUETP8M1) and 2017/2018 (TuneCP5) are different, so both need  
 1899 to be checked separately. For the case of 2016, the study documented in Ref. [8] assesses a  
 1900 systematic of half of the reweighting factor deviation from unity. Below, we check the 2018  
 1901 dataset explicitly to verify this is also the case for the new tune, TuneCP5.

1902 Following Ref. [8], we make a tight single-lepton selection with analysis jet, b-tagging, and  
 1903 lepton selections. Selected events have lepton  $p_T > 25 \text{ GeV}$ ,  $E_T^{\text{miss}} > 50 \text{ GeV}$ ,  $H_T > 300 \text{ GeV}$ ,  
 1904 and 2 b tagged jets. MC integral is normalized to data to account for trigger and lepton scale  
 1905 factors.

1906 Figure 127 shows the jet multiplicity distribution for 2018 data/MC before (left) and after (right)  
 1907 applying the 2018-derived reweighting factors. The systematic uncertainty of half of the de-  
 1908 viation from unity covers variations present in the reweighted single lepton jet multiplicity  
 1909 distributions for the updated TuneCP5.

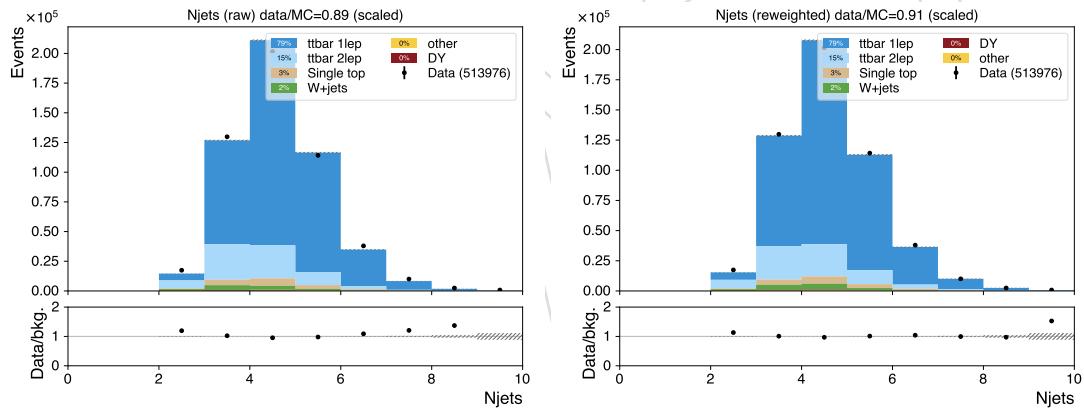


Figure 127: Jet multiplicity distribution for 2018 data/MC before (left) and after (right) ap-  
 plying the 2018-derived reweighting factors. The “other” category includes VV,  $t\bar{t}W$ , and  $t\bar{t}Z$   
 contributions.